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T.H. KOORNWINDER

A NEW PROOF OF A PALEY-WIENER TYPE THEOREM FOR THE
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A new proof of a Paley-Wiener type theorem for the Jacobi transform *)

by

T.H. Koornwinder

ABSTRACT

A new integral representation for Jacobi functions is derived, containing the Mehler-Dirichlet formula for Legendre functions as a special case. As a result, the Fourier-Jacobi transform, which generalizes the Mehler-Fok transform, can be factorized as the product of two Weyl type fractional integral transforms and a Fourier-cosine transform. There follow new short proofs of a Paley-Wiener type theorem and the inversion formula for the Jacobi transform. By analytic continuation these results hold for all complex values of the parameters.

*) This paper is not for review; it is meant for publication in a journal.

1. INTRODUCTION

Jacobi functions $\phi_\lambda(t)$ of order (α, β) are the eigenfunctions of the differential operator $(\Delta(t))^{-1}(d/dt)(\Delta(t)d/dt)$, $\Delta(t) = (e^t - e^{-t})^{2\alpha+1}(e^t + e^{-t})^{2\beta+1}$, such that $\phi_\lambda(0) = 1$, $\phi'_\lambda(0) = 0$. The Jacobi transform

$$(1.1) \quad f^\wedge(\lambda) = (2^{\frac{1}{2}}/\Gamma(\alpha+1)) \int_0^\infty f(t)\phi_\lambda(t)\Delta(t)dt,$$

which generalizes the Mehler-Fok transform, was studied by TITCHMARSH [23, §4.17], OLEVSKIĬ [21], BRAAKSMA and MEULENBELD [2], FLENSTED-JENSEN [9], [11, §2 and §12] and FLENSTED-JENSEN and KOORNWINDER [12]. Some papers by CHÉBLI [3], [4], [5] deal with a larger class of integral transforms which includes the Jacobi transform. An even more general class was considered by BRAAKSMA and DE SNOO [24].

In the present paper short proofs will be given of a Paley-Wiener type theorem and the inversion formula for the Jacobi transform. The L^2 -theory, i.e. the Plancherel theorem, is then an easy consequence. These results were earlier obtained by FLENSTED-JENSEN [9], [11, §2] and by CHÉBLI [5]. However, to prove the Paley-Wiener theorem these two authors needed the L^2 -theory, which can be obtained as a corollary of the Weyl-Stone-Titchmarsh-Kodaira theorem about the spectral decomposition of a singular Sturm-Liouville operator (cf. for instance DUNFORD and SCHWARTZ [6, Chap.13, §5]). The proofs presented here exploit the properties of Jacobi functions as hypergeometric functions and no general theorem needs to be invoked. Furthermore, it turns out that the Paley-Wiener theorem, which was proved by FLENSTED-JENSEN [11, §2] for real α, β , $\alpha > -1$, holds for all complex values of α and β .

The key formula in this paper is a generalized Mehler formula

$$(1.2) \quad (\Gamma(\alpha+1))^{-1}\Delta(t)\phi_\lambda(t) = \pi^{-\frac{1}{2}} \int_0^t \cos \lambda s A(s, t) ds,$$

where for $\operatorname{Re} \alpha > \operatorname{Re} \beta > -\frac{1}{2}$ $A(s, t)$ is given as an integral of elementary functions. Substituting (1.2) in (1.1) we can write the Jacobi transform

f^\wedge as the Fourier-cosine transform of $F(f)$, where the mapping F consists of two successive Weyl type fractional integral transforms. Thus the Jacobi transform is factorized as the product of three integral transforms with elementary kernels and the Paley-Wiener theorem follows from the mapping properties of these elementary transforms.

For certain discrete values of α and β the mapping F has a geometric and group-theoretic interpretation as a Radon transform on rank one symmetric spaces (cf. HELGASON [16, Chap.1,2]). For integer or half integer values of α and β such that $\alpha \geq \beta \geq -\frac{1}{2}$ a similar interpretation was given by FLENSTED-JENSEN [10] on certain pseudo-Riemannian symmetric spaces. A large class of integral transforms for which the corresponding mapping F is positive was examined by CHÉBLI [5]. Finally, FLENSTED-JENSEN and RAGOZIN [13] wrote a note on the analogue of (1.2) for spherical functions on non-compact symmetric spaces of arbitrary rank.

In section 2 of this paper some properties and formulas for Jacobi functions are given. Section 3 contains the proof of the Paley-Wiener theorem for all complex α and β . Formula (1.2) is the only result on Jacobi functions which is needed there. In section 4 the inversion formula is derived by using the Paley-Wiener theorem, some estimates for Jacobi functions and a formula for Jacobi functions of the second kind which is dual to (1.2). The paper concludes with some remarks, in particular about the Plancherel theorem and about Paley-Wiener type theorems for the Hankel transform and for Jacobi series.

Notation. This is mainly similar to the notation used in [12]. For reasons of elegance and in order to avoid singularities if $\alpha = -1, -2, \dots$, some constant factors have been changed. If no confusion is possible the indices α, β denoting the order may be deleted.

2. JACOBI FUNCTIONS

Consider for $\alpha, \beta, \lambda \in \mathbb{C}$ and $0 < t < \infty$ the differential equation

$$(2.1) \quad (\Delta_{\alpha, \beta}(t))^{-1} \frac{d}{dt} (\Delta_{\alpha, \beta}(t) \frac{du(t)}{dt}) = -(\lambda^2 + \rho^2)u(t),$$

where $\rho = \alpha + \beta + 1$ and

$$(2.2) \quad \Delta_{\alpha, \beta}(t) = (e^t - e^{-t})^{2\alpha+1} (e^t + e^{-t})^{2\beta+1}.$$

By substituting $z = -(\sinh t)^2$ in (2.1) a hypergeometric differential equation is obtained (cf.[7,2.1(1)]) with parameters $\frac{1}{2}(\rho+i\lambda)$, $\frac{1}{2}(\rho-i\lambda)$, $\alpha + 1$. Hence, if $\alpha \neq -1, -2, -3, \dots$ then the function

$$(2.3) \quad \phi_{\lambda}^{(\alpha, \beta)}(t) = F(\frac{1}{2}(\rho+i\lambda), \frac{1}{2}(\rho-i\lambda); \alpha + 1; -(\sinh t)^2)$$

is the solution of (2.1) which satisfies $\phi_{\lambda}(0) = 1$, $\phi'_{\lambda}(0) = 0$. Here the hypergeometric function $F(a, b; c; z)$ denotes the unique analytic continuation for $z \notin [1, \infty)$ of the power series

$$\sum_{n=0}^{\infty} \frac{(a)_n (b)_n}{(c)_n n!} z^n, \quad |z| < 1.$$

Note that $(\Gamma(\alpha+1))^{-1} \phi_{\lambda}^{(\alpha, \beta)}(t)$ is an entire function of α, β and λ (also for $\alpha = -1, -2, \dots$).

For $\lambda \neq -i, -2i, -3i, \dots$ another solution of (2.1) (cf.[7,2.9(9)]) is given by the function

$$(2.4) \quad \phi_{\lambda}^{(\alpha, \beta)}(t) = (e^t - e^{-t})^{i\lambda - \rho} \cdot F(\frac{1}{2}(-\alpha + \beta + 1 - i\lambda), \frac{1}{2}(\alpha + \beta + 1 - i\lambda); 1 - i\lambda; -(\sinh t)^{-2}).$$

This solution is characterized by the property that $\phi_{\lambda}(t) = e^{(i\lambda - \rho)t} (1 + o(1))$ for $t \rightarrow \infty$. The functions $\phi_{\lambda}(t)$ and $\Phi_{\lambda}(t)$ are called Jacobi functions of the first and second kind, respectively.

Using [7,2.10(2) and 2.10(5)] we obtain for non-integer λ the identity

$$(2.5) \quad \pi^{\frac{1}{2}}(\Gamma(\alpha+1))^{-1} \phi_{\lambda}(t) = \frac{1}{2}c(\lambda)\phi_{\lambda}(t) + \frac{1}{2}c(-\lambda)\phi_{-\lambda}(t),$$

where

$$(2.6) \quad c_{\alpha,\beta}(\lambda) = \frac{2^{\alpha+\beta+1} \Gamma(\frac{1}{2}i\lambda) \Gamma(\frac{1}{2}(1+i\lambda))}{\Gamma(\frac{1}{2}(\alpha+\beta+1+i\lambda)) \Gamma(\frac{1}{2}(\alpha-\beta+1+i\lambda))}.$$

Note that for real λ, α, β $\overline{c(\lambda)} = c(-\lambda)$.

It follows easily from (2.1) and the definitions of $\phi_{\lambda}(t)$, $\Phi_{\lambda}(t)$, $\Delta(t)$ and $c(\lambda)$ that

$$(2.7) \quad \left\{ \begin{array}{ll} \phi_{\lambda}^{(-\frac{1}{2}, -\frac{1}{2})}(t) = \cos \lambda t & , \quad \Phi_{\lambda}^{(-\frac{1}{2}, -\frac{1}{2})}(t) = e^{i\lambda t} , \\ \Delta_{-\frac{1}{2}, -\frac{1}{2}}(t) = 1 & , \quad c_{-\frac{1}{2}, -\frac{1}{2}}(\lambda) = 1 , \end{array} \right.$$

and

$$(2.8) \quad \left\{ \begin{array}{ll} \phi_{2\lambda}^{(\alpha, \alpha)}(t) = \phi_{\lambda}^{(\alpha, -\frac{1}{2})}(2t) & , \quad \Phi_{2\lambda}^{(\alpha, \alpha)}(t) = \Phi_{\lambda}^{(\alpha, -\frac{1}{2})}(2t) , \\ \Delta_{\alpha, \alpha}(t) = \Delta_{\alpha, -\frac{1}{2}}(2t) & , \quad c_{\alpha, \alpha}(2\lambda) = c_{\alpha, -\frac{1}{2}}(\lambda) . \end{array} \right.$$

The first two formulas of (2.8) can also be interpreted as quadratic transformations for hypergeometric functions, cf. [7,2.11(2) and 2.11(26)].

Application of [7,2.8(20) and 2.8(27)] gives the differentiation formulas

$$(2.9) \quad (\Gamma(\alpha+1))^{-1} \frac{d\phi_{\lambda}^{(\alpha, \beta)}(t)}{dt} =$$

$$= -\frac{1}{4}((\alpha+\beta+1)^2 + \lambda^2)(\Gamma(\alpha+2))^{-1} \sinh 2t \phi_{\lambda}^{(\alpha+1, \beta+1)}(t)$$

and

$$\begin{aligned}
(2.10) \quad & (\Gamma(\alpha+2))^{-1} \frac{d}{dt} \left[(\sinh 2t)^{-1} \Delta_{\alpha+1, \beta+1}(t) \phi_{\lambda}^{(\alpha+1, \beta+1)}(t) \right] = \\
& = 16(\Gamma(\alpha+1))^{-1} \Delta_{\alpha, \beta}(t) \phi_{\lambda}^{(\alpha, \beta)}(t).
\end{aligned}$$

Next we derive some useful integration formulas for Jacobi functions. It follows from Bateman's integral [7, 2.4(2)] and the identity

$$(2.11) \quad F(a, b; c; z) = (1-z)^{c-a-b} F(c-a, c-b; c; z)$$

(cf. [7, 2.1(23)]) that for $y > 0$, $\operatorname{Re} \mu > 0$, $\operatorname{Re} c > 0$

$$\begin{aligned}
(2.12) \quad & (\Gamma(c+\mu))^{-1} y^{c+\mu-1} (1+y)^{a+b-c+\mu} F(a+\mu, b+\mu; c+\mu; -y) = \\
& = \frac{1}{\Gamma(c)\Gamma(\mu)} \int_0^y x^{c-1} (1+x)^{a+b-c} F(a, b; c; -x) (y-x)^{\mu-1} dx.
\end{aligned}$$

It follows from ASKEY and FITCH [1, (2.10)] that for $x > 0$, $\operatorname{Re} \mu > 0$, $\operatorname{Re} b > 0$

$$\begin{aligned}
(2.13) \quad & \Gamma(b) x^{-b} F(a, b; c; -x^{-1}) = \\
& = \frac{\Gamma(b+\mu)}{\Gamma(\mu)} \int_x^\infty y^{-b-\mu} F(a, b+\mu; c; -y^{-1}) (y-x)^{\mu-1} dy.
\end{aligned}$$

Translating (2.12) and (2.13) in terms of Jacobi functions we obtain

$$\begin{aligned}
(2.14) \quad & (\Gamma(\alpha+\mu+1))^{-1} \Delta_{\alpha+\mu, \beta+\mu}(t) \phi_{\lambda}^{(\alpha+\mu, \beta+\mu)}(t) = \\
& = \frac{2^{3\mu+1} \sinh 2t}{\Gamma(\alpha+1)\Gamma(\mu)} \int_0^t \Delta_{\alpha, \beta}(s) \phi_{\lambda}^{(\alpha, \beta)}(s) (\cosh 2t - \cosh 2s)^{\mu-1} ds,
\end{aligned}$$

where $t > 0$, $\operatorname{Re} \mu > 0$, $\operatorname{Re} \alpha > -1$, and

$$\begin{aligned}
(2.15) \quad & (c_{\alpha, \beta}(-\lambda))^{-1} \phi_{\lambda}^{(\alpha, \beta)}(s) = \frac{2^{3\mu+1}}{c_{\alpha+\mu, \beta+\mu}(-\lambda)\Gamma(\mu)} \\
& \cdot \int_s^\infty \phi_{\lambda}^{(\alpha+\mu, \beta+\mu)}(t) (\cosh 2t - \cosh 2s)^{\mu-1} \sinh 2t dt,
\end{aligned}$$

where $s > 0$, $\operatorname{Re} \mu > 0$, $\operatorname{Im} \lambda > -\operatorname{Re}(\alpha+\beta+1)$.

The integrals (2.14) and (2.15) connect Jacobi functions of order (α, β) with functions of order $(\alpha-\beta-\frac{1}{2}, -\frac{1}{2})$ and Jacobi functions of order $(\alpha-\beta-\frac{1}{2}, \alpha-\beta-\frac{1}{2})$ with functions of order $(-\frac{1}{2}, -\frac{1}{2})$. Hence, by (2.7), (2.8), (2.14) and (2.15) we conclude that for $\operatorname{Re} \alpha > \operatorname{Re} \beta > -\frac{1}{2}$

$$(2.16) \quad (\Gamma(\alpha+1))^{-1} \Delta(t) \phi_\lambda(t) = \pi^{-\frac{1}{2}} \int_0^t \cos \lambda s A(s, t) ds$$

and

$$(2.17) \quad e^{i\lambda s} = (c(-\lambda))^{-1} \int_s^\infty \Phi_\lambda(t) A(s, t) dt, \quad \operatorname{Im} \lambda > 0,$$

where the kernel is given by

$$(2.18) \quad A_{\alpha, \beta}(s, t) = \frac{2^{3\alpha+5/2} \sinh 2t}{\Gamma(\alpha-\beta)\Gamma(\beta+\frac{1}{2})} \cdot \int_s^t (\cosh 2t - \cosh 2w)^{\beta-\frac{1}{2}} (\cosh w - \cosh s)^{\alpha-\beta-1} \sinh w dw.$$

By substituting $\tau = (\cosh t - \cosh w)/(\cosh t - \cosh s)$ in (2.18) and using Euler's integral [7, 2.1(10)] we obtain

$$(2.19) \quad A_{\alpha, \beta}(s, t) = 2^{3\alpha+2\beta+3/2} (\Gamma(\alpha+\frac{1}{2}))^{-1} \sinh 2t (\cosh t)^{\beta-\frac{1}{2}} \cdot (\cosh t - \cosh s)^{\alpha-\frac{1}{2}} F(\frac{1}{2}+\beta, \frac{1}{2}-\beta; \alpha+\frac{1}{2}; \frac{\cosh t - \cosh s}{2\cosh t}).$$

Combination of (2.19) and (2.11) gives

$$(2.20) \quad A_{\alpha, \beta}(s, t) = 2^{\alpha+2\beta+5/2} (\Gamma(\alpha+\frac{1}{2}))^{-1} \sinh 2t (\cosh t)^{\beta-\alpha} \cdot (\cosh 2t - \cosh 2s)^{\alpha-\frac{1}{2}} F(\alpha+\beta, \alpha-\beta; \alpha+\frac{1}{2}; \frac{\cosh t - \cosh s}{2\cosh t}).$$

Note that for $0 \leq s < t$ the argument of the hypergeometric functions in (2.19) and (2.20) has its value between 0 and $\frac{1}{2}$. Hence these hypergeometric functions are bounded functions in s and t . By analytic continuation with

respect to α and β and by using the expressions (2.19) or (2.20) for the kernel it follows that formula (2.16) is valid for $\operatorname{Re} \alpha > -\frac{1}{2}$ and formula (2.17) holds if $\operatorname{Re} \alpha > -\frac{1}{2}$, $\operatorname{Im} \lambda > 0$. It is clear from (2.19) and (2.20) that $A_{\alpha,\beta}(s,t) > 0$ if $0 \leq s < t$, $\alpha > -\frac{1}{2}$ and $|\beta| \leq \max(\frac{1}{2}, \alpha)$.

From (2.16) and (2.20) we have the integral representation

$$(2.21) \quad \phi_{\lambda}^{(\alpha,\beta)}(t) = 2^{-\alpha+3/2} \frac{\Gamma(\alpha+1)}{\Gamma(\alpha+\frac{1}{2})\Gamma(\frac{1}{2})} \frac{1}{(\sinh t)^{2\alpha}(\cosh t)^{\alpha+\beta}} \\ \cdot \int_0^t \cos \lambda s (\cosh 2t - \cosh 2s)^{\alpha-\frac{1}{2}} F(\alpha+\beta, \alpha-\beta; \alpha+\frac{1}{2}; \frac{\cosh t - \cosh s}{2\cosh t}) ds,$$

valid for $\operatorname{Re} \alpha > -\frac{1}{2}$. In view of (2.9) formula (2.21) in the case of order $(\alpha+1, \beta+1)$ gives an integral representation for $d\phi_{\lambda}^{(\alpha,\beta)}(t)/dt$. This last integral can be rewritten by using integration by parts and by [7, 2.8(27)]. Thus we obtain the integral representation

$$(2.22) \quad \frac{d\phi_{\lambda}^{(\alpha,\beta)}(t)}{dt} = -2^{-\alpha+3/2} \frac{\Gamma(\alpha+1)}{\Gamma(\alpha+\frac{1}{2})\Gamma(\frac{1}{2})} \frac{(\alpha+\beta+1)^2 + \lambda^2}{\lambda} \\ \cdot \frac{1}{(\sinh t)^{2\alpha+1}(\cosh t)^{\alpha+\beta}} \int_0^t \sin \lambda s \sinh s (\cosh 2t - \cosh 2s)^{\alpha-\frac{1}{2}} \\ \cdot F(\alpha+\beta+1, \alpha-\beta-1; \alpha+\frac{1}{2}; \frac{\cosh t - \cosh s}{2\cosh t}) ds,$$

which is also valid for $\operatorname{Re} \alpha > -\frac{1}{2}$.

We shall need some estimates which are essentially due to FLENSTED-JENSEN [9, Theorem 2], [11, §2], but which will be stated here for all $\alpha, \beta \in \mathbb{C}$. The proof of lemma 2.3 below is different from the proof given in [9].

LEMMA 2.1. *For each $\alpha, \beta \in \mathbb{C}$ and $\delta > 0$ there exists a positive constant K such that for all $t \geq \delta$ and all $\lambda \in \mathbb{C}$ with $\operatorname{Im} \lambda \geq 0$*

$$|\phi_{\lambda}^{(\alpha,\beta)}(t)| \leq K e^{-(\operatorname{Im} \lambda + \operatorname{Re} \rho)t}.$$

LEMMA 2.2. *For each $\alpha, \beta \in \mathbb{C}$ there exists a positive constant K such that if $\lambda \in \mathbb{C}$, $\operatorname{Im} \lambda \geq 0$ and λ is outside arbitrary small neighborhoods of the poles of $(c_{\alpha,\beta}(-\lambda))^{-1}$ then*

$$|c_{\alpha,\beta}(-\lambda)|^{-1} \leq K(1+|\lambda|)^{\alpha+\frac{1}{2}}.$$

Lemma 2.1 follows by extending the proof of [9, Lemma 7] to the case of complex α and β . Lemma 2.2 follows from (2.6) and Stirling's formula.

LEMMA 2.3. *For each $\alpha, \beta \in \mathbb{C}$ and for each non-negative integer n there exists a positive constant K such that for all $t \geq 0$ and all $\lambda \in \mathbb{C}$*

$$|(\Gamma(\alpha+1))^{-1} \frac{d^n}{dt^n} \phi_{\lambda}^{(\alpha,\beta)}(t)| \leq K(1+|\lambda|)^{n+k} (1+t) e^{(|\operatorname{Im} \lambda| - \operatorname{Re} \rho)t},$$

where $k = 0$ if $\operatorname{Re} \alpha > -\frac{1}{2}$ and $k = [\frac{1}{2} - \operatorname{Re} \alpha]$ if $\operatorname{Re} \alpha \leq -\frac{1}{2}$.

Proof. Consider first the case that $n = 0$ and $\operatorname{Re} \alpha > -\frac{1}{2}$. It follows from (2.21) that

$$\begin{aligned} |\phi_{\lambda}^{(\alpha,\beta)}(t)| &\leq \text{const.} \cdot e^{(|\operatorname{Im} \lambda| + \operatorname{Re}(\alpha-\beta))t} \\ &\quad \cdot (\sinh t \cosh t)^{-2\operatorname{Re} \alpha} \int_0^t (\cosh 2t - \cosh 2s)^{\operatorname{Re} \alpha - \frac{1}{2}} ds \\ &= \text{const.} \cdot e^{(|\operatorname{Im} \lambda| + \operatorname{Re}(\alpha-\beta))t} \phi_0^{(\operatorname{Re} \alpha, \operatorname{Re} \alpha)}(t). \end{aligned}$$

Applying [7, 2.10(7)] we have the estimate

$$\phi_0^{(\operatorname{Re} \alpha, \operatorname{Re} \alpha)}(t) \leq \text{const.} (1+t) e^{-(2\operatorname{Re} \alpha + 1)t}.$$

By combining the last two equalities the lemma is proved for $n = 0$. The estimate in the case that $n = 1$, $\operatorname{Re} \alpha > -\frac{1}{2}$ and $|\lambda| < 1$ follows from (2.9) and the estimate for $\phi_{\lambda}^{(\alpha+1, \beta+1)}(t)$. If $n = 1$, $\operatorname{Re} \alpha > -\frac{1}{2}$ and $|\lambda| \geq 1$ then we conclude from (2.22) that

$$\begin{aligned} \left| \frac{d}{dt} \phi_{\lambda}^{(\alpha,\beta)}(t) \right| &\leq \text{const.} (1+|\lambda|) e^{(|\operatorname{Im} \lambda| + \operatorname{Re}(\alpha-\beta))t} \phi_0^{(\operatorname{Re} \alpha, \operatorname{Re} \alpha)}(t) \leq \\ &\leq \text{const.} (1+|\lambda|)(1+t) e^{(|\operatorname{Im} \lambda| - \operatorname{Re} \rho)t}. \end{aligned}$$

The case that $n = 0, 1$ and $\operatorname{Re} \alpha \leq -\frac{1}{2}$ follows by complete induction with respect to $k = [\frac{1}{2} - \operatorname{Re} \alpha]$, where formulas (2.9) and (2.10) are used. Finally we prove the case that $n=2, 3, \dots$ by complete induction with respect to n using the formula

$$\begin{aligned}
(\Gamma(\alpha+1))^{-1} \frac{d^n}{dt^n} \phi_\lambda^{(\alpha, \beta)}(t) &= -(\rho^2 + \lambda^2)(\Gamma(\alpha+1))^{-1} \frac{d^{n-2}}{dt^{n-2}} \phi_\lambda^{(\alpha, \beta)}(t) + \\
&+ \frac{1}{2}(\rho^2 + \lambda^2)(\Gamma(\alpha+2))^{-1} \frac{d^{n-2}}{dt^{n-2}} \left[(\rho \cosh 2t + \alpha - \beta) \phi_\lambda^{(\alpha+1, \beta+1)}(t) \right].
\end{aligned}$$

This formula follows by differentiating the formula

$$\frac{d^2}{dt^2} \phi_\lambda^{(\alpha, \beta)}(t) = (\rho^2 + \lambda^2) \left[\frac{\rho \cosh 2t + \alpha - \beta}{2(\alpha+1)} \phi_\lambda^{(\alpha+1, \beta+1)}(t) - \phi_\lambda^{(\alpha, \beta)}(t) \right],$$

which is a consequence of (2.1) and (2.9). \square

3. A PALEY-WIENER TYPE THEOREM

Let C_0^∞ be the class of all even infinitely differentiable functions on \mathbb{R} with compact support. Let H be the class of even, entire, rapidly decreasing functions of exponential type, i.e., $g \in H$ if and only if g is an even and entire analytic function on \mathbb{C} and there exist positive constants A and K_n ($n=0,1,2,\dots$) such that for all $\lambda \in \mathbb{C}$ and for all $n=0,1,2,\dots$

$$(3.1) \quad |g(\lambda)| \leq K_n (1+|\lambda|)^{-n} e^{A|\operatorname{Im} \lambda|}.$$

Let for $f \in C_0^\infty$ and $\operatorname{Re} \alpha > -1$ the Fourier-Jacobi transform $f \rightarrow f_{\alpha,\beta}^\wedge$ be defined by

$$(3.2) \quad f_{\alpha,\beta}^\wedge(\lambda) = (2^{\frac{1}{2}}/\Gamma(\alpha+1)) \int_0^\infty f(t) \phi_\lambda^{(\alpha,\beta)}(t) \Delta_{\alpha,\beta}(t) dt.$$

Clearly $f_{\alpha,\beta}^\wedge(\lambda)$ is analytic in $\alpha, \beta, \lambda \in \mathbb{C}$ with $\operatorname{Re} \alpha > -1$. Substitution of (2.10) in (3.2) and repeated integration by parts gives

$$(3.3) \quad f_{\alpha,\beta}^\wedge(\lambda) = \frac{(-1)^n}{2^{4n} \Gamma(\alpha+n+1)} \int_0^\infty \left(\left(\frac{1}{\sinh 2t} \frac{d}{dt} \right)^n \bar{f}(t) \right) \cdot \phi_\lambda^{(\alpha+n,\beta+n)}(t) \Delta_{\alpha+n,\beta+n}(t) dt, \quad n=0,1,2,\dots$$

This formula defines the analytic continuation of $f_{\alpha,\beta}^\wedge(\lambda)$ for $\operatorname{Re} \alpha > -n-1$. Hence $f_{\alpha,\beta}^\wedge(\lambda)$ is an entire function of α, β, λ .

If $\alpha = \beta = -\frac{1}{2}$ then (3.2) reduces to the Fourier-cosine transform

$$(3.4) \quad f_{-\frac{1}{2}, -\frac{1}{2}}^\wedge(\lambda) = (2/\pi)^{\frac{1}{2}} \int_0^\infty f(t) \cos \lambda t dt.$$

THEOREM 3.1. (Paley and Wiener). *The Fourier-cosine transform is a bijection from C_0^∞ onto H .*

For a proof see for instance HÖRMANDER [17, Theorem 1.7.7]. In this section we shall generalize theorem 3.1 to general complex values of α and β .

Let for $f \in C_0^\infty$ and $\operatorname{Re} \alpha > -\frac{1}{2}$ the mapping $f \rightarrow F_{\alpha,\beta}(f)$ be defined by

$$(3.5) \quad (F_{\alpha,\beta}(f))(s) = \int_s^\infty f(t) A_{\alpha,\beta}(s,t) dt, \quad s > 0.$$

Note that $(F_{\alpha,\beta}(f))(s)$ is analytic in α and β . In particular, if $\operatorname{Re} \alpha > \operatorname{Re} \beta > -\frac{1}{2}$ then by (2.18) we have

$$(3.6) \quad (F_{\alpha,\beta}(f))(s) = \frac{2^{3\alpha+3/2}}{\Gamma(\alpha-\beta)} \int_s^\infty \left[\frac{1}{\Gamma(\beta+\frac{1}{2})} \int_w^\infty f(t) (\cosh 2t - \cosh 2w)^{\beta-\frac{1}{2}} d(\cosh 2t) \right] \\ \cdot (\cosh w - \cosh s)^{\alpha-\beta-1} d(\cosh w).$$

Combining (2.16), (3.2) and (3.5) we obtain that for $f \in C_0^\infty$ and $\operatorname{Re} \alpha > -\frac{1}{2}$

$$(3.7) \quad f_{\alpha,\beta}^\wedge(\lambda) = (2/\pi)^{\frac{1}{2}} \int_0^\infty (F_{\alpha,\beta}(f))(s) \cos \lambda s ds.$$

This means that the Jacobi transform of order (α,β) of f is the cosine transform of $F_{\alpha,\beta}(f)$.

To analyze the transform $F_{\alpha,\beta}$ consider the Weyl fractional integral transform \mathcal{W}_μ which is for $a \in \mathbb{R}$, $g \in C_0^\infty([a,\infty))$ and $\operatorname{Re} \mu > 0$ defined by

$$(3.8) \quad (\mathcal{W}_\mu(g))(y) = (\Gamma(\mu))^{-1} \int_y^\infty g(x) (x-y)^{\mu-1} dx$$

(cf. [8, Chap. 13]). Here $C_0^\infty([a,\infty))$ denotes the class of infinitely differentiable functions on the interval $[a,\infty)$ (right differentiable in a) with compact support. Repeated integration by parts in (3.8) gives

$$(3.9) \quad (\mathcal{W}_\mu(g))(y) = \frac{(-1)^n}{\Gamma(\mu+n)} \int_y^\infty \left(\frac{d^n}{dx^n} g(x) \right) (x-y)^{\mu+n-1} dx, \quad n=0,1,2,\dots$$

By (3.8) and (3.9) $(\mathcal{W}_\mu(g))(y)$ is defined as an entire function in μ , continuous in $(\mu, y) \in \mathbb{C} \times [a,\infty)$. Clearly, the function $\mathcal{W}_\mu(g)$ has also compact support and, since $(\mathcal{W}_\mu(g))' = \mathcal{W}_\mu(g')$ we conclude that $\mathcal{W}_\mu(g) \in C_0^\infty([a,\infty))$. It is an easy exercise to prove that $\mathcal{W}_0 = \operatorname{id}$, $\mathcal{W}_{-1}(g) = -g'$, $\mathcal{W}_\mu \circ \mathcal{W}_\nu = \mathcal{W}_{\mu+\nu}$. In particular, $\mathcal{W}_\mu \circ \mathcal{W}_{-\mu} = \operatorname{id} = \mathcal{W}_{-\mu} \circ \mathcal{W}_\mu$. This proves the following theorem.

THEOREM 3.2. For all $a \in \mathbb{R}$ and $\mu \in \mathbb{C}$ the mapping W_μ , defined by (3.9), is bijective from $C_0^\infty([a, \infty))$ onto itself.

Let us next define for $f \in C_0^\infty$, $\operatorname{Re} \mu > 0$, $\sigma > 0$, $s \geq 0$

$$(3.10) \quad (W_\mu^\sigma(f))(s) = (\Gamma(\mu))^{-1} \int_s^\infty f(t) (\cosh \sigma t - \cosh \sigma s)^{\mu-1} d(\cosh \sigma t).$$

Again we can extend $(W_\mu^\sigma(f))(s)$ to an entire function of μ by

$$(3.11) \quad (W_\mu^\sigma(f))(s) = \frac{(-1)^n}{(\mu+n)} \int_s^\infty \left(\frac{d^n}{d(\cosh \sigma t)^n} f(t) \right) \cdot (\cosh \sigma t - \cosh \sigma s)^{\mu+n-1} d(\cosh \sigma t),$$

$$n=0,1,2,\dots, \quad \operatorname{Re} \mu > -n.$$

Let $f(t) = g(\cosh \sigma t)$. Then $f \in C_0^\infty$ if and only if $g \in C_0^\infty([1, \infty))$. Hence it follows from theorem 3.2 that for each $\mu \in \mathbb{C}$ the mapping W_μ^σ is bijective from C_0^∞ onto itself. The inverse mapping of W_μ^σ is $W_{-\mu}^\sigma$. Applying this result to (3.6) we obtain

COROLLARY 3.3. If $f \in C_0^\infty$ then $(F_{\alpha,\beta}(f))(s)$ has an analytic continuation to an entire function in α and β which is given by

$$(3.12) \quad F_{\alpha,\beta}(f) = 2^{3\alpha+3/2} W_{\alpha-\beta}^1 \circ W_{\beta+\frac{1}{2}}^2(f).$$

For all $\alpha, \beta \in \mathbb{C}$ the mapping $F_{\alpha,\beta}$ is bijective from C_0^∞ onto itself. The inverse mapping is given by

$$(3.13) \quad f = 2^{-3\alpha-3/2} W_{-\beta-\frac{1}{2}}^2 \circ W_{\beta-\alpha}^1 \circ F_{\alpha,\beta}(f).$$

Combination of Theorem 3.1, corollary 3.3 and formula (3.7) gives the Paley-Wiener type theorem for the Jacobi-transform.

THEOREM 3.4. For all $\alpha, \beta \in \mathbb{C}$ the mapping $f \rightarrow f_{\alpha,\beta}^\wedge$ is bijective from C_0^∞ onto H .

4. THE INVERSION FORMULA

It is well-known that the inversion formula for the cosine transform is given by

$$(4.1) \quad f(t) = (2/\pi)^{\frac{1}{2}} \int_0^{\infty} f_{-\frac{1}{2}, -\frac{1}{2}}^{\wedge}(\lambda) \cos \lambda t \, d\lambda,$$

where $f \in C_0^{\infty}$ and $f_{-\frac{1}{2}, -\frac{1}{2}}^{\wedge}(\lambda)$ is defined by (3.4). Substituting $\cos \lambda t = \frac{1}{2}e^{i\lambda t} + \frac{1}{2}e^{-i\lambda t}$ and changing the path of integration in (4.1) we also have

$$(4.2) \quad f(t) = (2\pi)^{-\frac{1}{2}} \int_{i\eta-\infty}^{i\eta+\infty} f_{-\frac{1}{2}, -\frac{1}{2}}^{\wedge}(\lambda) e^{i\lambda t} d\lambda,$$

where η is an arbitrary real number. In this section we shall generalize (4.1) and (4.2) to inversion formulas for the Jacobi transform.

Let for $g \in H$, $t > 0$ and $\alpha, \beta \in \mathbb{C}$

$$(4.3) \quad g_{\alpha, \beta}^{\vee}(t) = (2\pi)^{-\frac{1}{2}} \int_{i\eta-\infty}^{i\eta+\infty} g(\lambda) \phi_{\lambda}^{(\alpha, \beta)}(t) (c_{\alpha, \beta}(-\lambda))^{-1} d\lambda,$$

where $\eta \geq 0$, $\eta > -\operatorname{Re}(\alpha + \beta + 1)$, $\eta > -\operatorname{Re}(\alpha - \beta + 1)$, i.e., $(c_{\alpha, \beta}(-\lambda))^{-1}$ is a regular function of λ for $\operatorname{Im} \lambda \geq \eta$. Let for $g \in H$ A be a positive constant such that the estimates (3.1) hold and choose $\delta > 0$. Then by lemmas 2.1 and 2.2 there is a positive constant K such that for all $t \geq \delta$ and all $\lambda \in \mathbb{C}$ with $\operatorname{Im} \lambda \geq 0$ and λ outside arbitrary small neighborhoods of the poles of $(c_{\alpha, \beta}(-\lambda))^{-1}$ we have

$$(4.4) \quad |g(\lambda) \phi_{\lambda}^{(\alpha, \beta)}(t) (c_{\alpha, \beta}(-\lambda))^{-1}| \leq K e^{-\rho t} (1 + |\lambda|)^{-2} e^{(A-t)\operatorname{Im} \lambda}.$$

It follows that the integral in (4.3) absolutely converges and that its value does not depend on the choice of η . In particular, if $|\operatorname{Re} \beta| < \operatorname{Re}(\alpha + 1)$ then we can put $\eta = 0$ in (4.3) and by (2.5) we obtain

$$(4.5) \quad g_{\alpha, \beta}^{\vee}(t) = \frac{\sqrt{2}}{\Gamma(\alpha + 1)} \int_0^{\infty} \frac{g(\lambda) \phi_{\lambda}^{(\alpha, \beta)}(t)}{c_{\alpha, \beta}(\lambda) c_{\alpha, \beta}(-\lambda)} d\lambda.$$

LEMMA 4.1. Let $\operatorname{Re} \alpha > -\frac{1}{2}$ and $|\operatorname{Re} \beta| < \operatorname{Re}(\alpha+1)$. If $g \in H$ then $g_{\alpha,\beta}^v \in C_0^\infty$ and $(g_{\alpha,\beta}^v)_{\alpha,\beta}^\wedge = g$.

Proof. It follows from (4.3) and (4.4) by letting $\eta \rightarrow \infty$ that $g_{\alpha,\beta}^v(t) = 0$ if $t > A$. It is clear from (4.5) that $g_{\alpha,\beta}^v$ is even. The estimates from lemmas 2.2 and 2.3 and formula (3.1) show that

$$\begin{aligned} |g(\lambda) \left(\frac{d^n}{dt^n} \phi_\lambda^{(\alpha,\beta)}(t) \right) (c_{\alpha,\beta}(\lambda) c_{\alpha,\beta}(-\lambda))^{-1}| &\leq \\ &\leq \text{const.} (1+t) e^{-(\operatorname{Re} \rho)t} (1+\lambda)^{-2}, \end{aligned}$$

uniformly if $\lambda, t \geq 0$. Hence, by (4.5), $g_{\alpha,\beta}^v \in C_0^\infty$. To prove the second part of the theorem observe that for $\eta > 0$ and $s > 0$

$$\begin{aligned} (F_{\alpha,\beta}(g_{\alpha,\beta}^v))(s) &= \\ &= (2\pi)^{-\frac{1}{2}} \int_s^\infty A_{\alpha,\beta}(s,t) dt \int_{i\eta-\infty}^{i\eta+\infty} g(\lambda) \phi_\lambda^{(\alpha,\beta)}(t) (c_{\alpha,\beta}(-\lambda))^{-1} d\lambda = \\ &= (2\pi)^{-\frac{1}{2}} \int_{i\eta-\infty}^{i\eta+\infty} \left[\int_s^\infty \phi_\lambda^{(\alpha,\beta)}(t) (c_{\alpha,\beta}(-\lambda))^{-1} A_{\alpha,\beta}(s,t) dt \right] g(\lambda) d\lambda, \end{aligned}$$

where the interchanging of integrals is allowed by Fubini's theorem, in view of (4.4) and the estimate

$$|A_{\alpha,\beta}(s,t)| \leq \text{const.} e^{\rho t} (t-s)^{\alpha-\frac{1}{2}}, \quad t > s > 0,$$

which is evident from (2.20). Inserting (2.17) we find that

$$(F_{\alpha,\beta}(g_{\alpha,\beta}^v))(s) = (2\pi)^{-\frac{1}{2}} \int_{i\eta-\infty}^{i\eta+\infty} g(\lambda) e^{i\lambda s} d\lambda.$$

By inverting this formula it follows that

$$(g_{\alpha,\beta}^v)_{\alpha,\beta}^\wedge(\lambda) = (2/\pi)^{\frac{1}{2}} \int_0^\infty (F_{\alpha,\beta}(g_{\alpha,\beta}^v))(s) \cos \lambda s ds = g(\lambda). \quad \square$$

THEOREM 4.2. Let $\alpha, \beta \in \mathbb{C}$. Then $f \in C_0^\infty$ and $g = f_{\alpha, \beta}^\wedge$ if and only if $g \in H$ and $f = g_{\alpha, \beta}^\vee$.

Proof. In view of theorem 3.4 it is sufficient to prove that $(f_{\alpha, \beta}^\wedge)_{\alpha, \beta}^\vee(t) = f(t)$ if $f \in C_0^\infty$, $t > 0$ and $\alpha, \beta \in \mathbb{C}$. By theorem 4.1 this is true for $\operatorname{Re} \alpha > -\frac{1}{2}$, $|\operatorname{Re} \beta| < \operatorname{Re}(\alpha+1)$. By (3.3) and (4.3) $(f_{\alpha, \beta}^\wedge)_{\alpha, \beta}^\vee(t)$ is an entire function of α and β . Hence the theorem follows by analytic continuation. \square

5. SOME REMARKS

Remark 1. Suppose that $(c_{\alpha,\beta}(-\lambda))^{-1}$ has N poles $\lambda_1, \lambda_2, \dots, \lambda_N$ such that $\text{Im } \lambda_n > 0$. Then a formula similar to (4.5) can be derived with additional terms of the type $c_n^{(\alpha,\beta)} g(\lambda_n) \phi_{\lambda_n}^{(\alpha,\beta)}(t)$, $n=1,2,\dots,N$ (cf. FLENSTED-JENSEN [11,§2]). Complications arise if some pole of $(c_{\alpha,\beta}(-\lambda))^{-1}$ is not simple or lies on the real axis or coincides with a pole of $(c_{\alpha,\beta}(\lambda))^{-1}$.

Remark 2. Let $f \in C_0^\infty$ and $g \in H$. Suppose for convenience that $(c_{\alpha,\beta}(-\lambda))^{-1}$ has no poles for $\text{Im } \lambda \geq 0$, i.e., $|\text{Re } \beta| < \text{Re}(\alpha+1)$. Then it is clear from (3.2) and (4.5) that

$$\int_0^\infty f(t) g^\vee(t) \Delta(t) dt = \int_0^\infty f^\wedge(\lambda) g(\lambda) (c(\lambda) c(-\lambda))^{-1} d\lambda.$$

Here Fubini's theorem is used together with the estimates of lemmas 2.2 and 2.3 and formula (3.1). It follows by theorem 4.2 that for $f_1, f_2 \in C_0^\infty$

$$(5.1) \quad \int_0^\infty f_1(t) f_2(t) \Delta(t) dt = \int_0^\infty f_1^\wedge(\lambda) f_2^\wedge(\lambda) (c(\lambda) c(-\lambda))^{-1} d\lambda.$$

Remark 3. For real α and β , $|\beta| < \alpha+1$, formula (5.1) implies Parseval's formula

$$(5.2) \quad \int_0^\infty f_1(t) \overline{f_2(t)} \Delta(t) dt = \int_0^\infty f_1^\wedge(\lambda) \overline{f_2^\wedge(\lambda)} |c(\lambda)|^{-2} d\lambda,$$

where $f_1, f_2 \in C_0^\infty$. Hence, since C_0^∞ is dense in $L^2(\Delta)$ and H is dense in $L^2(|c(\lambda)|^{-2})$, the Jacobi transform can be extended to an isometric mapping from $L^2(\Delta)$ onto $L^2(|c(\lambda)|^{-2})$. This gives an alternative proof for the Plancherel theorem obtained by FLENSTED-JENSEN [9, Prop. 3]. If $(c(-\lambda))^{-1}$ has poles for $\text{Im } \lambda > 0$ then a discrete spectrum must be added (cf. [11, §2]).

Remark 4. A Paley-Wiener type theorem for the Hankel transform can be proved by similar methods as in section 3. Let $J_\alpha(t)$ be a solution of the differential equation $u''(t) + (2\alpha+1)t^{-1}u'(t) + u(t) = 0$, $\alpha \neq -1, -2, \dots$, such that $J_\alpha(0) = 1, J'_\alpha(0) = 0$. Then $J_\alpha(t) = 2^{\alpha\Gamma(\alpha+1)} t^{-\alpha} J_\alpha(t)$, where $J_\alpha(t)$ is a Bessel function. If $\text{Re } \alpha > -\frac{1}{2}$ then it follows from the Poisson integral representation

$$J_{\alpha}(t) = \frac{\Gamma(\alpha+1)}{\pi^{\frac{1}{2}}\Gamma(\alpha+\frac{1}{2})} \int_0^{\pi} e^{it \cos \phi} (\sin \phi)^{2\alpha} d\phi$$

that

$$(5.3) \quad t^{2\alpha} J_{\alpha}(\lambda t) = \frac{2\Gamma(\alpha+1)}{\pi^{\frac{1}{2}}\Gamma(\alpha+\frac{1}{2})} \int_0^t \cos \lambda s (t^2 - s^2)^{\alpha-\frac{1}{2}} ds.$$

Define for $f \in C_0^{\infty}$ and $\operatorname{Re} \alpha > -1$ the Hankel transform by

$$(5.4) \quad f^{\wedge}(\lambda) = \frac{1}{2^{\alpha}\Gamma(\alpha+1)} \int_0^{\infty} f(t) J_{\alpha}(\lambda t) t^{2\alpha+1} dt.$$

Then

$$(5.5) \quad f^{\wedge}(\lambda) = (2/\pi)^{\frac{1}{2}} \int_0^{\infty} \cos \lambda s ds \frac{1}{2^{\alpha+\frac{1}{2}}\Gamma(\alpha+\frac{1}{2})} \\ \cdot \int_s^{\infty} f(t) (t^2 - s^2)^{\alpha-\frac{1}{2}} d(t^2), \quad \operatorname{Re} \alpha > -\frac{1}{2}.$$

Formula (5.5) is analogous to (3.6) and (3.7) and it can be used in a similar way.

Remark 5. For certain discrete values of α and β Jacobi functions are the spherical functions on non-compact symmetric spaces of rank one. In this context many formulas and results of [9] and the present paper were earlier obtained. Formula (3.7) corresponds to HELGASON [15,(9)]. The function $e^{-\rho s}(F_{\alpha,\beta}(f))(s)$ has a geometric interpretation as a Radon transform, where f is a radial function on the symmetric space (cf. HELGASON [16,Chap.1,2]). The Paley-Wiener theorem for the spherical Fourier transform on non-compact symmetric spaces of rank one was first proved by HELGASON [15].

REMARK 6. Formulas (2.16) and (2.18) generalize the classical Mehler-Dirichlet formula (cf. MEHLER [20])

$$P_{\nu}(\cos \theta) = \frac{2^{\frac{1}{2}}}{\pi} \int_0^{\theta} \frac{\cos(\nu+\frac{1}{2})\phi}{(\cos \phi - \cos \theta)^{\frac{1}{2}}} d\phi,$$

where $P_{\nu}(x)$ is a Legendre function. These formulas can also be obtained

from the Laplace type integral representation

$$(5.6) \quad \phi_{\lambda}^{(\alpha, \beta)}(t) = \frac{2\Gamma(\alpha+1)}{\pi^{\frac{1}{2}}\Gamma(\alpha-\beta)\Gamma(\beta+\frac{1}{2})} \int_0^1 \int_0^{\pi} |\cosh t + \sinh t \operatorname{re}^{i\psi}|^{i\lambda-\rho} \\ \cdot (1-r^2)^{\alpha-\beta-1} r^{2\beta+1} (\sin \psi)^{2\beta} dr d\psi, \quad t > 0, \operatorname{Re} \alpha > \operatorname{Re} \beta > -\frac{1}{2}$$

(cf. [18,(4)], [9,(3.5)]) by substituting first $\cosh t + \sinh t \operatorname{re}^{i\psi} = e^s e^{i\chi}$ and next $\cosh w = \cos \chi \cosh t$. A general method of transforming integrals of type (5.6) into integrals of type (2.16) is discussed in [19,§5].

Remark 7. Let $R_n^{(\alpha, \beta)}(x) = P_n^{(\alpha, \beta)}(x) / P_n^{(\alpha, \beta)}(1)$, where $P_n^{(\alpha, \beta)}(x)$ is a Jacobi polynomial. Then

$$R_n^{(\alpha, \beta)}(\cos \theta) = F(-n, n+\alpha+\beta+1; \alpha+1; \sin^2 \frac{1}{2}\theta) = \phi_{(2n+\alpha+\beta+1)1}^{(\alpha, \beta)}(\frac{1}{2}i\theta).$$

Analogous to (2.16), (2.18) and (2.19) we obtain

$$(5.7) \quad R_n^{(\alpha, \beta)}(\cos \theta) = \frac{2^{\alpha-2\beta-\frac{1}{2}}\Gamma(\alpha+1)}{\pi^{\frac{1}{2}}\Gamma(\alpha-\beta)\Gamma(\beta+\frac{1}{2})} (\sin \frac{1}{2}\theta)^{-2\alpha} (\cos \frac{1}{2}\theta)^{-2\beta} \\ \cdot \int_0^{\theta} (\cos \psi - \cos \theta)^{\beta-\frac{1}{2}} \sin \frac{1}{2}\psi d\psi \int_0^{\phi} \cos(n+\frac{1}{2}(\alpha+\beta+1))\phi \\ \cdot (\cos \frac{1}{2}\phi - \cos \frac{1}{2}\psi)^{\alpha-\beta-1} d\phi, \quad \operatorname{Re} \alpha > \operatorname{Re} \beta > -\frac{1}{2}, 0 < \theta < \pi,$$

$$(5.8) \quad R_n^{(\alpha, \beta)}(\cos \theta) = \frac{2^{\alpha-\frac{1}{2}}\Gamma(\alpha+1)}{\pi^{\frac{1}{2}}\Gamma(\alpha+\frac{1}{2})} (\sin \frac{1}{2}\theta)^{-2\alpha} (\cos \frac{1}{2}\theta)^{-\beta-\frac{1}{2}} \\ \cdot \int_0^{\theta} \cos(n+\frac{1}{2}(\alpha+\beta+1))\phi F(\frac{1}{2}+\beta, \frac{1}{2}-\beta; \alpha+\frac{1}{2}; \frac{\cos \frac{1}{2}\theta - \cos \frac{1}{2}\phi}{2\cos \frac{1}{2}\theta}) d\phi, \\ \operatorname{Re} \alpha > -\frac{1}{2}, 0 < \theta < \pi.$$

Quadratic transformation of the hypergeometric function in (5.8) by means of [7,2.11(22)] gives another integral representation for $R_n^{(\alpha, \beta)}(\cos \theta)$, which was independently obtained by GASPER [14] in a quite different way.

Remark 8. Suppose that f is an even C^∞ -function on $(-\pi, \pi)$ with compact

support. If f is expanded in a Fourier-Jacobi series with respect to $R_n^{(\alpha, \beta)}(\cos \theta)$ ($\alpha > \beta > -\frac{1}{2}$) then the Fourier coefficients are given by

$$(5.9) \quad f^\wedge(n) = (\Gamma(\alpha+1))^{-1} \int_0^\pi f(\theta) R_n^{(\alpha, \beta)}(\cos \theta) (\sin \frac{1}{2}\theta)^{2\alpha+1} (\cos \frac{1}{2}\theta)^{2\beta+1} d\theta, \\ n=0, 1, 2, \dots$$

Substitution of (5.7) in (5.9) gives

$$(5.10) \quad f^\wedge(n) = \frac{2^{\alpha-2\beta-3/2}}{\pi^{1/2} \Gamma(\alpha-\beta) \Gamma(\beta+\frac{1}{2})} \int_0^\pi \cos(n+\frac{1}{2}(\alpha+\beta+1))\phi \, d\phi \\ \cdot \int_{\psi=\phi}^\pi (\cos \frac{1}{2}\phi - \cos \frac{1}{2}\psi)^{\alpha-\beta-1} d(\cos \frac{1}{2}\psi) \int_{\theta=\psi}^\pi f(\theta) (\cos \psi - \cos \theta)^{\beta-\frac{1}{2}} d(\cos \theta).$$

In the same way as in section 3 we can write

$$(5.11) \quad f^\wedge(n) = \int_0^\pi \cos(n+\frac{1}{2}(\alpha+\beta+1))\phi \, (F(f))(\phi) d\phi,$$

where the mapping F is a bijection from the class of even C^∞ -functions on $(-\pi, \pi)$ with compact support onto itself. Then the function f^\wedge is well-defined and analytic for all complex values of its argument. Now the classical Paley-Wiener theorem implies a Paley-Wiener type theorem for Jacobi series.

THEOREM 5.1. *Let $\alpha > \beta > -\frac{1}{2}$. The function f^\wedge is the Fourier-Jacobi transform of an even C^∞ -function on $(-\pi, \pi)$ with compact support if and only if there is a function $g \in H$ such that $A < \pi$ in (3.1) and $f^\wedge(n) = g(n+\frac{1}{2}(\alpha+\beta+1))$, $n=0, 1, 2, \dots$.*

Since g is of exponential type less than π an application of Carlson's theorem (cf. TITCHMARSH [23, §5.81]) shows that g is uniquely determined by $f_{\alpha, \beta}^\wedge(n)$, $n=0, 1, 2, \dots$. Just as in section 3 theorem 5.1 remains valid for all $\alpha, \beta \in \mathbb{C}$. R. ASKEY informed me that in the case $\alpha = \beta = 0$ this theorem is due to BEURLING (unpublished).

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