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# On the Nonexistence of a Strong Solution in the Boundary Problem for a Sticky Brownian Motion 

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#### Abstract

We prove that a sticky Brownian motion is not measurable with respect to a driving Wiener process thereby verifying Skorokhod's conjecture. Note: Work on this paper was completed while the author was visiting the Centre for Mathematics and Computer Science at Amsterdam, during the period september - december 1988. The paper will appear in Probability Theory and Related Fields. Keywords \& Phrases: Brownian motion, boundary problem, strong solution. 1980 Mathematics Subject Classification: 60H10.


We consider the stochastic differential equation

$$
\begin{equation*}
d \xi_{t}=a I_{\left[\xi_{t}=0\right]} d t+I_{\left[\xi_{t}>0\right]} d W_{t} \quad \xi_{0}=X_{0} \geqslant 0 \tag{1}
\end{equation*}
$$

where $W$ is a Wiener process and $0 \leqslant a \leqslant \infty$.
It is convenient to represent (1) as a boundary problem with the departure rate of a general form (giving by the way the exact meaning to the condition $a=\infty$ ).

For the nonnegative measurable function $a_{t}, t \geqslant 0$, the process $\xi$ is a solution of the boundary problem if
a) $\xi_{0}=X_{0}$,
b) $\xi_{t} \geqslant 0, t \geqslant 0$,
c) $I_{\left[\xi_{t}>0\right]} d \xi_{t}=I_{\left[\xi_{t}>0\right]} d W_{t}$,
d) $d \zeta_{t}=I_{\left[\xi_{t}=0\right]} d \xi_{t} \geqslant 0$,
e) $\left.d \zeta_{t}=a_{t} I_{[ } a_{t}<\infty\right] d \wedge_{t}$,
f) $I_{\left[a_{t}=\infty\right]} d \wedge_{t}=0$,
where $\wedge_{t}=\int_{0}^{t} I_{\left[\xi_{j}=0\right]} d s, t \geqslant 0$.
It is well known that the equation (1) has a unique weak solution ([1]; for more general boundary problems see [2]).

For the extremal cases $a=0$ (absorbtion) and $a=\infty$ (instantaneous reflection) the problem (2) admits the strong solutions expressed explicitly by

$$
\begin{equation*}
X_{t}^{0}=X_{0}+W_{\tau \wedge t}, \tau=\min \left\{s \geqslant 0, X_{0}+W_{s}=0\right\} \tag{3}
\end{equation*}
$$

and

$$
\begin{equation*}
X_{t}^{*}=\max \left(X_{0}+W_{t}, \max _{0 \leqslant s \leqslant t}\left(W_{t}-W_{s}\right)\right) \tag{4}
\end{equation*}
$$

In the intermediate case with the finite positive rate of departure from the boundary, however, the sticky (or slowly reflecting) Brownian motion $\xi$ is not representable as a functional of the (driving)

Theorem. The equation (1) does not admit a strong solution for $0<a<\infty$.
We prove this statement by constructing a sequence $\xi^{n}, n \geqslant 1$, of solutions of (2) adopted to $W$, which converges to the solution of (1) in the stable topology ([3]), but diverges in the strong sense (pathwise). It follows from this that the solution $\xi$ of (1) is nonmeasurable with respect to $W$. To this end we proceed in several stages.

1. The uniqueness of the solution measure of (1).

Definition. A probability $\mu$ on the Borel $\sigma$-algebra $B\left(C^{2}\right)$ of the product space $C^{2}=C \times C$ of two dimentional continuous functions $(x, y)=\left(x_{t}, y_{t}\right), t \geqslant 0$, is called the solution measure of (2) if

$$
\begin{align*}
& \mu\left\{(x, y): x_{t}=x_{0}+a \int_{0}^{t} I_{\left[x_{\mathrm{s}}=0\right]} d s+\int_{0}^{t} I_{\left[x_{s}>0\right]} d y_{s}\right\}=1,  \tag{5}\\
& \mu\{(x, y): y \in B\}=P^{W}(B), B \in B(C), \tag{6}
\end{align*}
$$

where $P^{W}$ is a Wiener measure.
Denote by $\mu^{1}, \mu^{2}, \mu^{1}(\cdot \cdot \cdot), \mu^{2}(\cdot 1 \cdot)$ the marginal and conditional distributions corresponding to $\mu$ :

$$
\begin{aligned}
& \mu^{\prime}(B)=\mu(B \times C), \mu^{2}(B)=\mu(C \times B) ; B \in B(C) \\
& \mu^{\prime}(B \mid y)=\mu\{x \in B \mid y\}, \mu^{2}(B \mid x)=\mu(y \in B \mid x) ; x, y \in C .
\end{aligned}
$$

The following construction of the solution measure is presented in [1]. Let $P^{x}$ be the (unique) distribution of the weak solution of (1), and define $\bar{\mu}$ as an image measure corresponding to the mapping $\varphi$ from the probability space ( $C^{2}, B\left(C^{2}\right), P^{X} \times P^{W}$ ) into the space ( $C^{2}, B\left(C^{2}\right)$ ) expressed by

$$
\varphi(x, y)=\left(x_{t}, x_{t}-x_{0}-a \int_{0}^{t} I_{\left[X_{s}=0\right]} d s+\int_{0}^{t} I_{\left[X_{s}=0\right]} d y_{s}\right), t \geqslant 0 .
$$

It is easily seen that $\bar{\mu}$ is a solution measure of (1) characterized by the property

$$
\bar{\mu}\left\{\int_{0}^{1} I_{\left[X_{s}=0\right]} d y_{s} \in B \mid x\right\}=Q_{x}(B), B \in B(C), \text { a.s; }
$$

where, for each $x \in C, Q_{x}$ is the distribution of a Gaussian process with independent increments expressable as

$$
\int_{0}^{1} I_{\left[X_{1}=0\right]} d W_{s}
$$

with a Wiener process $W$.
The property $(7)$ is equivalent to the relation

$$
\begin{equation*}
\left.E^{\bar{\mu}} f(x) \exp \left(i \int_{0}^{1} C_{s} I_{[X,}=0\right] d y_{s}\right)=E^{\bar{\mu}} f(x) \exp \left(-\frac{1}{2} \int_{0}^{1} C_{s}^{2} I_{[X,=0]} d s\right), \tag{8}
\end{equation*}
$$

satisfied for any bounded measurable functions

$$
\begin{aligned}
& f:(C, B(C)) \rightarrow\left(R^{(1)}, B\left(R^{(1)}\right)\right) \\
& C:\left(R_{+}^{(1)}, B\left(R_{+}^{(1)}\right) \rightarrow\left(R^{(1)}, B\left(R^{(1)}\right)\right)\right.
\end{aligned}
$$

and $t \geqslant 0$.
Lemma. The equation (1) admits an unique solution measure.
Proof. It is sufficient to show only that for any solution measure of (1) the relation (8) takes place. By the uniqueness of $\mu^{1}=P^{X}$ as a solution of the Martingale problem corresponding to (1), it
follows (see [4]) that every (bounded) measurable function

$$
\phi:(C, B(C)) \rightarrow\left(R^{(1)}, B\left(R^{(1)}\right)\right)
$$

is expressable as a stochastic integral

$$
\phi(x)=E^{\mu} \phi(x)+\int_{0}^{\infty} g_{s}\left(d x_{s}-a I_{\left[X_{s}=0\right]} d s\right), P^{X} \text { a.s. }
$$

with some nonanticipative functional $g_{t}=g_{t}(X)$.
Taking

$$
\phi(x)=f(x) \exp \left(-\frac{1}{2} \int_{0}^{1} C_{s}^{2} I_{[X,=0]} d s\right)
$$

and denoting

$$
\rho_{t}=\exp \left(i \int_{0}^{t} C_{s} I_{\left[X_{s}=0\right]} d y_{s}+\frac{1}{2} \int_{0}^{t} C_{s}^{2} I_{\left[X_{s}=0\right]} d_{s}\right)
$$

we have

$$
\begin{aligned}
& E^{\mu} f(x) \exp \left(i \int_{0}^{t} C_{s} I_{\left[X_{s}=0\right]} d y_{s}\right)=E^{\mu} \phi(x) \rho_{t}=E^{\mu} \phi(x) E^{\mu} \rho_{t}+ \\
+ & E^{\mu} \int_{0}^{t} g_{s}\left(d x_{s}-a I_{\left[x_{s}=0\right]} d s\right) \rho_{t}=E^{\mu} \phi(x)+E^{\mu} \int_{0}^{t} g_{s}\left(d x_{s}-a I_{\left[x_{s}=0\right]} d s\right)+ \\
+ & E^{\mu} \int_{0}^{t} i g_{s} C_{s} I_{\left[x_{s}>0\right]} I_{\left[x_{s}=0\right]} d s=E^{\mu} \phi(x) \square
\end{aligned}
$$

2. The compactness of the class of measure solutions of (2)

Let $M$ be the class of all solution measures corresponding to the (uniquely solvable) boundary problems (2) with arbitrary departure rates $\left(a_{t}\right), t \geqslant 0$. We shall consider the convergence

$$
u_{n} \rightarrow \mu, n \rightarrow \infty
$$

of measures $\mu_{n}, \mu$ on $B\left(C^{2}\right)$ determined by the convergence of integrals

$$
E^{\mu} f_{n}(x, y) \rightarrow E^{\mu} f(x, y), n \rightarrow \infty
$$

for each $t$ and each bounded measurable in $y$ and continuous in $x$ function on $f$ depending on first $t$ coordinates $\left(f(x, y)=f\left(x_{s}, y_{s} ; s \leqslant t\right)\right.$ ).
Denote by $M^{s}$ the subclass of $M$ corresponding to the boundary problems (2) having unique strong solutions. If $\mu \in M^{s}$, then there exists some mapping

$$
\varphi^{\mu}:(C, B(C)) \rightarrow(C, B(C))
$$

such that for each $A, B \in B(C)$

$$
\mu(A \times B)=\mu^{2}\left(B \cap\left(\varphi^{\mu}\right)^{-1}(A)\right)=\int_{\{y \in B, \varphi(y) \in A\}} \mu^{2}(d y)
$$

Evidently $\mu^{2}=P^{W}$, and

$$
\begin{equation*}
\xi_{t}=\varphi_{t}^{\mu}(y) P^{W} \text { a.s. } 0 \leqslant t \tag{9}
\end{equation*}
$$

represents the solution of (2) as a functional of the driving process.

LEmma 2. ([3]) a). The weak compactness of the class $\left(\mu^{1}: \mu \in M\right.$ ) implies the (stable) compactness (corresponding to the convergence $\rightarrow$ ) of $M$.

$$
\text { b) If } \mu_{n} \in M^{s}, \mu_{n} \rightarrow \mu, n \rightarrow \infty, \mu \in M^{S}
$$

then

$$
\sup _{0 \leqslant s \leqslant t}\left|\varphi_{s}^{\mu_{n}}-\varphi_{s}^{\mu}\right| \xrightarrow{P^{*}} 0, n \rightarrow \infty, t \geqslant 0
$$

Note that the class $M^{s}$ consists of the solution measures corresponding to the class $a^{+}$of functions $\left(a_{t}\right), t \rightarrow 0$ which are stepwise constant taking on only two values 0 and $\infty$.

Lemma 3. The class $M$ is (relatively) compact.
Proof. Let $\mu \in M$ with some $\left(a_{t}\right), t \geqslant 0$. Consider the time change $\tau_{t}=\inf \left(s: s-\wedge_{s}=t\right)$ and the process

$$
X_{t}^{\mu \mu}=X_{\tau_{i}}
$$

It is clear from

$$
\begin{aligned}
\int_{0}^{t} I_{\left[X_{s}^{\mu}=0\right]} d s & =\int_{0}^{t} I_{\left[X_{\mathrm{rs}}=0\right]} d s=\int_{0}^{t} I_{\left[X_{\mathrm{r}}=0\right]} d\left(\tau_{s}-\wedge_{\tau_{s}}\right) \\
& =\int_{0}^{t} I_{\left[X_{\mathrm{s}}=0\right]}\left(1-I_{\left[X_{\mathrm{s}}=0\right]}\right) d \tau_{s}=0
\end{aligned}
$$

that the distribution of $X^{* \mu}$ coincides with the distribution of an instantaneous reflecting Brownian motion $X^{*}$ stopped at some random time.

We have for each $0 \leqslant h \leqslant 1$ and $t \geqslant 0$, using $\left|\tau_{t}-\tau_{s}\right| \leqslant|t-s|$,

$$
\begin{aligned}
& \Delta_{\epsilon}^{\mu}(h)=\mu^{1}\left\{X: \sup _{0 \leqslant s<u<s+h \leqslant t}\left|X_{s}-X_{u}\right|>\epsilon\right\} \leqslant \\
& \leqslant P^{X}\left\{X: \sup _{0 \leqslant s<u<s+h \leqslant t}\left|X_{s}-X_{u}\right|>\epsilon\right\}=\Delta_{\epsilon}(h)
\end{aligned}
$$

Thus the condition of weak (relative) compactness of the class $\left\{\mu^{1}: \mu \in M\right\}$, ([5]), follows from the continuity of the process $X^{*}$.

$$
\lim _{h \rightarrow 0} \sup _{\mu \in M} \Delta_{\epsilon}^{\mu}(h) \leqslant \lim _{n \rightarrow 0} \Delta_{\epsilon}(h)=0
$$

Hence the assertion is true by Lemma 2.a)

## 3. The conditions of the convergence to the measure solution of (1)

Lemma 4. Let $\mu$ be the (unique) measure solution of (1) and let $\mu_{n} \in M$, with $\mu_{n}^{1}$ having strong Markovian property, $n \geqslant 1$. Then, if for each $t>0, \lambda>0$

$$
\begin{align*}
& L_{t}^{n}(\lambda)=E^{\mu_{n}}\left(\int_{t}^{\infty} \exp (-\lambda(s-t)) d \zeta_{s} \mid x_{t}=0\right) \rightarrow \frac{a}{\lambda+\sqrt{2 \lambda} a}, n \rightarrow \infty  \tag{10}\\
& M_{t}^{n}(\lambda)=E^{\mu_{n}}\left(\int_{t}^{\infty} \exp (-\lambda(s-t)) I_{\left[X_{s}=0\right]} d s \mid x_{t}=0\right) \rightarrow \frac{1}{\lambda+\sqrt{2 \lambda} a}, n \rightarrow \infty
\end{align*}
$$

then $\mu_{n} \rightarrow \mu$
Proof. If $\mu_{n}^{1} \rightarrow \mu^{1}$, then for each limit point $\tilde{\mu}$ of the sequence $\mu_{n}$ we have for each $\Gamma>0$

$$
\begin{equation*}
\left.\left.\eta=\eta(x, y)=\sup _{0 \leqslant s<t \leqslant T S \leqslant u \leqslant t} \inf _{u} X_{u}\right)\left(\sup _{S \leqslant v \leqslant t} \mid\left(X_{\sigma}-X_{s}\right)-y_{v}-y_{s}\right) \mid\right)=0 \tilde{\mu} \text { a.s. } \tag{11}
\end{equation*}
$$

In fact $\eta=0$ is equivalent to the property that for each $0 \leqslant \leqslant s<t \leqslant T$, with $X_{u}>0, S \leqslant u \leqslant t$ ) the equality $Y_{u}-Y_{s}=X_{u}-X_{s}, s \leqslant u \leqslant t$, takes place. From c) in (2) it follows that $\eta=0 \mu_{n}$ a.s. $n \geqslant 1$, and, hence, to obtain (11) it is sufficient to take limit in $E^{\mu_{n}} \min (\eta \wedge C)$ for each $C>0$, using the continuity of $\eta$ on $C^{2}$.

Thus, as the conditions a), b), d), e), f) in (2) are expressed only in terms of the marginal distribution $\mu^{1}$, we have $\tilde{\mu}=\mu$.

Hence it remains to prove that the convergence (10) implies the convergence of finite dimensional distributions corresponding to $\mu_{n}^{1}, n \geqslant 1$.

It is not difficult to calculate the Laplace transforms

$$
\begin{aligned}
& L_{t}(\lambda)=L_{0}(\lambda)=E^{\mu}\left(\int_{t}^{\infty} \exp (-\lambda(s-t)) d \zeta_{s} \mid x_{t}=0\right) \\
& M_{t}(\lambda)=M_{0}(\lambda)=E^{\mu}\left(\int_{t}^{\infty} \exp (-\lambda(s-t)) I_{\left[X_{s}=0\right]} d s \mid x_{t}=0\right)
\end{aligned}
$$

by solving the boundary problem

$$
\begin{align*}
& \left.\lambda U(x, \Lambda)=\frac{1}{2} U_{x x}^{(x,} \lambda\right), x>0  \tag{12}\\
& \lambda U(x, \lambda)=a U_{x x}(x, \lambda)+1, x=0
\end{align*}
$$

for

$$
U(x, \lambda)=E^{\mu}\left(\int_{0}^{\infty} \exp (-\lambda s) I_{\left[X_{s}=0\right]} d s \mid x_{0}=x\right)
$$

From (12) we obtain

$$
L_{t}(\lambda)=a M_{t}(\lambda)=a M_{0}(\lambda), M_{0}(\lambda)=\frac{1}{\lambda+\sqrt{2 \lambda} a} .
$$

Further, for each $n \rightarrow 1$, the conditional Laplace transform

$$
\phi_{n}(s, t, \lambda, x)=E^{\mu_{n}}\left(\exp \left(-\lambda x_{t}\right) \mid x_{s}=x\right)
$$

satisfies the equation (for each $t>s, x \geqslant 0$ )

$$
\begin{gathered}
\phi_{n}(s, t, \lambda, x)=e^{-\lambda x}-\lambda E^{\mu_{n}}\left(\zeta_{t}-\zeta_{s} \mid x_{s}=x\right)+\frac{\lambda^{2}}{2} \int_{s}^{t} \phi_{n}(s, u, \lambda, x) d u \\
-\frac{\lambda^{2}}{2} E^{\mu_{n}}\left(\wedge_{t}-\wedge_{s} \mid x_{s}=x\right),
\end{gathered}
$$

which gives

$$
\phi_{n}(s, t, \lambda, x)=e^{-\lambda x+\lambda^{2}}(t-s)-E^{\mu_{x}}\left(\left.\int_{s}^{t} \exp \left(\frac{\lambda^{2}}{2}(+-u)\right)\left(\lambda d \zeta_{u}+\frac{\lambda^{2}}{2} d \wedge_{u}\right) \right\rvert\, x_{s}=x\right)
$$

Besides, by the strong Markovian property, we have

$$
\begin{aligned}
& E^{\mu_{n}}\left(\zeta_{t}-\zeta_{s} \mid x_{s}=x\right)=\int_{s}^{t} e_{u}^{n}(t) k(x, u-s) d u \\
& E^{\mu_{s}}\left(\wedge_{t}-\wedge_{s} \mid x_{s}=x\right)=\int_{s}^{t} m_{u}^{n}(t) k(x, u-s) d u
\end{aligned}
$$

where, for $t>u$,

$$
l_{u}^{n}(t)=E^{\mu_{n}}\left(\zeta_{t}-\zeta_{u} \mid x_{u}=0\right), m_{u}^{n}(t)=E^{\mu_{n}}\left(\wedge_{t}-\wedge \mid x_{u}=0\right),
$$

and $k(x, t)$ is the distribution density function of the random moment

$$
\tau=\min \left(t \geqslant s: x_{t}=0\right)
$$

with the condition $x_{s}=x$. Evidently $k$ does not depend on $n \geqslant 1$ and $s$, and, in terms of the Wiener process $W$,

$$
k(x, t)=\frac{d}{d t} P\left(x+\inf _{0 \leqslant s \leqslant t} W_{s}<0\right)=\frac{d}{d t} P\left(\left|W_{t}\right|>x\right) .
$$

Thus the convergence of $l_{t}^{n}(s)$ and $m_{t}^{n}(s)$ (or the convergence (10) of their Laplace transforms $L_{t}^{n}(\lambda)$ and $\left.M_{t}^{n}(\lambda)\right)$ is sufficient for the convergence of the conditional Laplace transforms $\phi_{n}$
4. The necessary condition for the strong convergence

Let $\mu_{n} \in M^{s}$, and let

$$
X_{t}^{n}=\varphi_{t}^{\mu_{n}}(y)
$$

denote the strong solutions of (2).
Lemma 5. If for each $t>0, \lambda>0$

$$
P^{W}\left(\sup _{0 \leqslant s \leqslant t}\left|X_{s}^{n}-X_{s}^{m}\right|>\epsilon\right) \rightarrow 0, n, m \rightarrow \infty
$$

then

$$
E^{W} \int_{0}^{\infty} \exp (-\lambda s)\left[I_{\left[x_{s}^{n}=0\right]}-I_{\left[x_{s}^{m}=0\right]} \mid d s \rightarrow 0, n, m \rightarrow \infty\right.
$$

Proof. From (2) we have

$$
E^{W}\left(X_{t}^{n}-X_{t}^{m}\right)^{2}=2 E^{W} \int_{0}^{t}\left(X_{s}^{n}-X_{s}^{m}\right) d\left(\zeta_{s}^{n}-\zeta_{s}^{m}\right)+E^{W} \int_{0}^{t}\left(I_{\left[X_{s}^{n}>0\right]}-I_{\left[X_{s}^{\prime \prime}>0\right]}\right)^{2} d s
$$

Thus

$$
\begin{align*}
& E^{W} \int_{0}^{t}\left|I_{\left[X_{s}^{n}=0\right]}-I_{\left[X_{s}^{m}=0\right]}\right| d s \leqslant E^{W}\left(\sup _{0 \leqslant s \leqslant t}\left|X_{s}^{n}-X_{s}^{m}\right|\right)^{2}+ \\
& +2\left[E^{W}\left(\sup _{0 \leqslant s \leqslant t}\left|X_{s}^{n}-X_{s}^{m}\right|\right)^{2} E^{W}\left(\zeta_{t}^{n}+\zeta_{t}^{m}\right)^{2}\right]^{\frac{1}{2}} \tag{13}
\end{align*}
$$

Now it is sufficient to apply the fact that the instantaneous reflecting process $X^{*}$ is maximal in the class of all strong solutions ([6]) and, so $X_{t}^{n} \leqslant X_{t}^{*}$ a.s. Thus using the estimations

$$
\sup _{0 \leqslant s \leqslant t}\left|X_{s}^{n}-X_{s}^{m}\right| \leqslant 2 \sup _{0 \leqslant s \leqslant t} X_{s}^{*}, E^{W}\left(\zeta_{t}^{n}\right)^{2} \leqslant 2 E^{W}\left[\left(X_{t}^{n}\right)^{2}+t\right] \leqslant 2 E^{W}\left[\left(X_{t}^{*}\right)^{2}+t\right]
$$

we obtain that the right-hand side of (13) converges to zero

## 5. The description of the approximations and the weak convergence

Consider for each $\Delta>0,0 \leqslant \alpha \leqslant 1,0<\delta<\Delta$, the unique strong solution (expressed as a functional of a Wiener process $W$ )

$$
\xi_{t}(\Delta, \alpha, \delta)=\varphi_{t}^{\mu}(W)
$$

with the departure rate $\left(a_{t}\right), t \geqslant 0$ of the form

$$
\begin{aligned}
& a_{t}=0, \text { for } t \in[k \Delta, k \Delta+\alpha \Delta] \cup[k \Delta+\alpha \Delta+\delta,(k+1) \Delta[ \\
& \left.a_{t}=\infty, \text { for } t \in\right] k \Delta+\alpha \Delta, k \Delta+\alpha \Delta+\delta[, k \geqslant 0
\end{aligned}
$$

Thus the process $\xi(\Delta, \alpha, \delta)$ is everywhere absorbing except subintervals of length $\delta$ disposed at the one and the same positions inside the intervals $[k \Delta,(k+1) \Delta[, k=0,1, \ldots$.

Lemma 6. For each $0 \leqslant \alpha<1$ and $0<c<\infty$ the sequence $\xi^{n}=\xi\left(\frac{1}{n}, \alpha, \frac{c^{2}}{n^{2}}\right)$ converges weakly to the solution of (1) with $a=\sqrt{\frac{\pi}{2}} c$.

Proof. It is sufficient to verify the conditions (10).
Consider first the case $\alpha=0$. It is easy to notice that the functions $L_{i}^{\Delta, \delta}(\lambda)$ and $M_{i}^{\Delta \delta}(\lambda)$ corresponding to the solution $\xi(\Delta, 0, \delta)$ satisfy the relations

$$
\begin{align*}
& L_{t}^{\Delta \delta}(\lambda)=L_{0}^{\Delta . \delta}(\lambda), k \Delta+\delta \leqslant t \leqslant(k+1) \Delta, k \geqslant 0,  \tag{14}\\
& M_{t}^{\Delta, \delta}(\lambda)=M_{0}^{\delta}(\lambda) \exp \left(\left(\left[\frac{t}{\Delta}\right]+1\right) \Delta-t\right)+\frac{1}{\lambda}\left(\exp \left(\left[\left[\frac{t}{\Delta}\right]+1\right) \Delta-t\right)-1\right),
\end{align*}
$$

where $\left[\frac{t}{\Delta}\right]$ is a largest integer of $\frac{t}{\Delta}$.
To derive the recurrent equations for $M_{0}^{\Delta}, \delta(\lambda)$ and $L_{0}^{\Delta, \delta}(\lambda)$ introduce the moment

$$
\tau=\min \left(s \geqslant \delta, \xi_{s}(\Delta, 0, \delta)=0\right)
$$

Suppose $\xi_{0}(\Delta, 0, \delta)=0$. Then, by definition, the representations (3) and (4) give

$$
\begin{aligned}
& \xi_{t}(\Delta, 0, \delta)=\sup _{0 \leqslant s \leqslant t}\left(W_{t}-W_{s}\right), 0 \leqslant t \leqslant \sigma \\
& \xi_{t}(\Delta, 0, \Delta)-\xi_{\delta}(\Delta, 0, \delta)=W_{t}-W_{\delta}, \quad \delta \leqslant t \leqslant \tau
\end{aligned}
$$

Thus (taking into consideration that the random variables $\sup _{0 \leqslant s \leqslant t} W_{s}$ and $\left|W_{t}\right|$ have the same distribution) we obtain

$$
\begin{align*}
F^{\Delta, \delta}(t) & =P^{W}\left\{\tau<t \mid \zeta_{0}(\Delta, 0, \delta)=0\right\}=P^{W}\left\{\inf _{\delta \leqslant s \leqslant t+\delta}\left(W_{s}-W_{\delta}\right) \leqslant-\sup _{0 \leqslant s \leqslant \delta}\left(W_{\delta}-W_{s}\right)\right\}= \\
& =P\left\{\left|W_{t+\delta}-W_{\delta}\right| \geqslant\left|W_{\delta}\right|\right\}=\frac{2}{\pi t} \int_{0}^{\infty} \int_{y}^{\infty} \exp \left(-\frac{s^{2}}{2 t}-\frac{y^{2}}{2 \delta}\right) d x d y=\frac{2}{\pi} \operatorname{arctg}\left(\sqrt{\frac{t}{\delta}}\right) \tag{15}
\end{align*}
$$

Consider, for convenience,

$$
\tilde{M}_{t}^{\Delta, \delta}(\lambda)=\frac{1}{\lambda}-M_{t}^{\Delta, \delta}(\lambda)=E^{\mu}\left(\int_{t}^{\infty} \exp (-\lambda(s-t)) I_{\left[x_{s}>0\right]} d s \mid x_{t}=0\right)
$$

We have

$$
\begin{aligned}
\tilde{M}_{0}^{\Delta, \delta}(\lambda) & =E \int_{0}^{\tau} \exp (-\lambda t) d t+E \exp (-\lambda \tau) \tilde{M}_{\tau}^{\Delta, \delta}(\lambda)=\int_{0}^{\infty} \frac{1}{\lambda}(1-\exp (-\lambda t)) F^{\Delta, \delta}(d t)+ \\
& +\tilde{M}_{0}^{\Delta, \delta}(\lambda) E \exp \left(-\lambda\left(\left[\frac{\tau}{\Delta}\right]+1\right) \Delta\right) I_{\left[\tau-\left[\frac{\tau}{\Delta}\right] \Delta \geqslant \delta\right]}+E \exp (-\tau \lambda) I_{\left[\tau-\left[\frac{\tau}{\Delta}\right] \Delta \leqslant \delta\right]} \tilde{M}_{\tau}^{\Delta, \delta}(\lambda)
\end{aligned}
$$

Thus

$$
\begin{equation*}
\tilde{M}_{0}^{\Delta, \delta}(\lambda)=\left[\phi_{1}(\Delta, \delta, \lambda)+\phi_{2}(\Delta, \delta, \lambda)\right] \angle\left(1-\phi_{3}(\Delta, \delta, \lambda)\right), \tag{16}
\end{equation*}
$$

where

$$
\begin{aligned}
& \phi_{1}(\Delta, \delta, \lambda)=\int_{0}^{\infty} \frac{1}{\lambda}(1-\exp (-\lambda t)) F^{\Delta, \delta}(d t) \\
& \phi_{2}(\Delta, \delta, \lambda)=E\left(\exp (-\lambda \tau) I_{\left[\tau-\left[\frac{\tau}{\Delta}\right] \Delta \leqslant \delta\right]} \tilde{M}_{\tau}^{\Delta, \delta}(\lambda)\right. \\
& \phi_{3}(\Delta, \delta, \lambda)=E \exp \left(-\lambda \Delta\left(\left[\frac{\tau}{\Delta}\right]+1\right)\right) I_{\left[\tau-\left[\frac{\tau}{\Delta}\right] \Delta \geqslant \delta\right]}
\end{aligned}
$$

For large $X$ we shall use the estimation

$$
\begin{equation*}
\left|\operatorname{arctg}(x)-\frac{\pi}{2}+\frac{1}{x}\right| \leqslant 2\left(\frac{1}{x}\right)^{3} . \tag{17}
\end{equation*}
$$

Using (15) and (17), since obviously $\tilde{M}_{t}^{\Delta, \delta}(\lambda) \leqslant \frac{1}{\lambda}$, we have for $\delta=c^{2} \Delta^{2}$

$$
\begin{gather*}
\phi_{2}\left(\Delta, c^{2} \Delta^{2}, \lambda\right) \leqslant \sum_{k=1}^{\infty}\left(F^{\Delta, \delta}\left(k \Delta+c^{2} \Delta^{2}\right)-F^{\Delta \delta}(k \Delta)\right) \exp (-k \lambda \Delta)= \\
=\frac{2}{\pi \lambda} \sum_{k=1}^{\infty} \exp (-k \Delta \lambda)\left[\operatorname{arctg}\left(\frac{\sqrt{k+c^{2} \Delta}}{c^{2} \Delta}\right)-\operatorname{arctg}\left(\sqrt{\frac{k}{c^{2}} \Delta}\right) \leqslant\right. \\
\left.\leqslant \frac{2 c \sqrt{\Delta}}{\pi \lambda} \sum_{k=1}^{\infty} \exp (-k \lambda \Delta) \angle \sqrt{\frac{1}{k+c^{2} \Delta}}-\sqrt{\frac{1}{k}} \right\rvert\,+0\left(\Delta^{3 / 2}\right) \leqslant \frac{2 c \sqrt{\Delta}}{\pi \lambda} \sum_{k=1}^{\infty} \frac{c^{2} \Delta}{2 k^{3 / 2}}  \tag{18}\\
+0\left(\Delta^{3 / 2}\right)=0\left(\Delta^{3 / 2}\right)
\end{gather*}
$$

Further, applying

$$
\lim _{\Delta \rightarrow 0} \sqrt{\Delta} \sum_{k=1}^{\infty} \exp (-\lambda k \Delta) \frac{1}{\sqrt{k}}=\sqrt{\frac{\pi}{\lambda}}
$$

we obtain

$$
\begin{align*}
& \phi_{3}\left(\Delta, c^{2} \Delta^{2}, \lambda\right)=\sum_{k=1}^{\infty}(1-\exp (-\lambda k \Delta))\left(F^{\Delta, \delta}(k \Delta)-F^{\Delta \delta}((k-1) \Delta)=\right. \\
& =(1-\exp (-\lambda \Delta)) \sum_{k=1}^{\infty} \exp (-\lambda k \Delta) F^{\Delta, \delta}(k \Delta)=\exp (-\lambda \Delta)- \\
& -(1-\exp (-\lambda \Delta))\left[\sum_{k=1}^{\infty} \exp (-\lambda k \Delta) \sqrt{\frac{\Delta}{k}} \frac{2 c}{\pi}+0(\Delta)\right]=\exp (-\lambda \Delta)- \\
& -(1-\exp (-\lambda \Delta))\left[2 \sqrt{\frac{1}{\pi \lambda}} c+0(\Delta)\right]=1-\lambda \Delta-2 c \sqrt{\frac{\lambda}{\pi}}+o(\Delta) \tag{19}
\end{align*}
$$

As for the expression $\phi_{1}$, it is easily calculated that

$$
\begin{align*}
& \phi_{1}\left(\Delta, C^{2} \Delta^{2}, \lambda\right)=\frac{\Delta c}{\lambda \pi} \int_{0}^{\infty}(1-\exp (-\lambda t)) \frac{d t}{+} \frac{1}{2}\left(t+C^{2} \Delta^{2}\right)= \\
& =\frac{\Delta c}{\lambda \pi}\left[\int _ { 0 } ^ { \infty } \left(1-\exp \left(-\lambda \frac{t)) d t}{t^{3 / 2}}+0(\Delta)\right]=\frac{2 \Delta c}{\sqrt{\pi \lambda}}+0(\Delta)\right.\right. \tag{20}
\end{align*}
$$

Thus using (18), (19) and (20) we obtain

$$
\lim _{\Delta \geqslant 0} \tilde{M}_{0}^{\Delta, \delta}(\lambda)=\sqrt{\frac{2}{\pi}} \frac{c \sqrt{2 \lambda}}{\lambda\left(\lambda+\sqrt{2 \lambda} c \sqrt{\frac{2}{\pi}}\right)}
$$

and hence

$$
\lim _{\Delta \rightarrow 0} M_{0}^{\Delta, \delta}=\frac{1}{\lambda+\sqrt{2 \lambda a}}
$$

with

$$
a=c \sqrt{\frac{2}{\pi}} .
$$

Analogeously, from the decomposition

$$
L_{0}^{\Delta, \delta}(\lambda)=E \int_{0}^{\tau} \exp (-\lambda t) d \zeta_{t}+E \exp (-\lambda \tau) L_{\tau}^{\Delta, \delta}(\lambda)
$$

we can derive, using the relation (14), and applying the same arguments as before, that

$$
L_{0}^{\Delta, \delta}(\lambda)=E \int_{0}^{\tau} e^{-\lambda t} d \zeta_{t}\left(\Delta\left(\lambda+\sqrt{2 \lambda} \sqrt{\frac{2}{\pi}} c\right)\right)^{-1}+o(1)
$$

Besides

$$
\begin{aligned}
& E \int_{0}^{\tau} \exp (-\lambda t) d \zeta_{t}=E \int_{0}^{\delta} \exp (-\lambda t) d \zeta_{t}=E \zeta_{\delta}+o(\Delta)= \\
& =E \sup _{0 \leqslant s \leqslant \delta} W_{s}+0(\Delta)=\Delta \sqrt{\frac{2}{\pi}} c+o(\Delta)
\end{aligned}
$$

Thus

$$
\lim _{\Delta \rightarrow 0} L_{0}^{\Delta} \delta=\frac{a}{\lambda+a \sqrt{2 \lambda}}
$$

with

$$
a=c \sqrt{\frac{2}{\pi}} .
$$

From (14) it is evident that (10) is true for all $t \geqslant 0$.
Finally, it is easyly seen that

$$
L_{t+\alpha \Delta}^{\Delta, \alpha, \delta}(\lambda)=L_{t}^{\Delta, 0, \delta}(\lambda), M_{t+\alpha \Delta}^{\Delta, \alpha, \delta}(\lambda)=M_{t}^{\Delta, 0, \delta}(\lambda)
$$

and, thus (10) is true for each $0 \leqslant \alpha<1, t \geqslant 0$.
6. The strong nonconvergence and the proof of the theorem Consider now the sequence $\xi^{n}$ defined for $m \geqslant 1$ as follows:

$$
\begin{aligned}
\xi^{n} & =\xi\left(\frac{1}{2 m}, 0, \frac{c^{2}}{(2 m)^{2}}\right), \text { as } n=2 m \\
\xi^{n} & =\xi\left(\frac{1}{2 m}, \frac{1}{2}, \frac{c^{2}}{(2 m)^{2}}\right), \text { as } n=2 m+1
\end{aligned}
$$

## Lemma 7. For the sequence $\xi^{n}$ defined above

$$
\left.\varlimsup_{n \rightarrow \infty} E \int_{0}^{\infty} \exp (-\lambda s) \mid I_{\left.\xi_{j}=0\right]}-I_{\left[\xi^{n}\right.}=0\right]
$$

Proof. Obviously

$$
\begin{align*}
& \left.\left.E \int_{0}^{\infty} \exp (-\lambda s)\right|_{\left[\xi_{s}^{\prime}=0\right]}-I_{\left[\xi_{2}^{n}\right.}=0\right] \\
& -E \int_{0}^{\infty} \exp (-\lambda s) I\left[\max \left(\xi_{s}^{n}, \xi_{1}^{n+1}\right)=0\right] d s . \tag{20}
\end{align*}
$$

Let us consider the process $\bar{\xi}$ which is the strong solution of (2) with

$$
\begin{aligned}
& a_{t}=\infty, \text { for } t \in[k \Delta, k+\Delta+\delta] \cup\left[k \Delta+\frac{\Delta}{2}, k \Delta+\frac{\Delta}{2}+\delta\right], \\
& a_{t}=0, \text { otherwise. }
\end{aligned}
$$

It is easyly seen that (with $\delta<\frac{\Delta}{2}$ )

$$
\bar{\xi}_{t}=\max \left(\xi_{t}(\Delta, 0, \delta), \xi_{t}\left(\Delta, \frac{1}{2}, \delta\right)\right)=\xi_{t}\left(\frac{\Delta}{2}, 0, \delta\right)
$$

Thus, for $\delta=c^{2} \Delta^{2}, \Delta=(2 m)^{-1}$

$$
\begin{aligned}
& E \int_{0}^{\infty} \exp (-\lambda s)\left|I_{\left[\xi_{0}^{n}=0\right]}-I_{\left[\xi_{0}^{\prime \prime}=0\right]}\right| d s \geqslant M_{0}^{\Delta \cdot c^{2} \Delta^{2}} \rightarrow \\
& \rightarrow \frac{1}{\lambda+a \sqrt{2 \lambda}}-\frac{1}{\lambda+a^{1} \sqrt{2 \lambda}}, \Delta \rightarrow 0,
\end{aligned}
$$

where

$$
a=c \sqrt{\frac{2}{\pi}}, a^{1}=2 c \sqrt{\frac{2}{\pi}}=2 a .
$$

Combining now the statements of lemmas and taking into consideration the necessary condition for the strong convergence in assertion b) of Lemma 2, we obtain the proof of the theorem.

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