# THE LINEAR SYSTEMS LIE ALGEBRA，THE SEGAL－SHALE－WEIL REPRESENTATION AND ALL KALMAN－BUCY FILTERS 

Michiel Hazewinkel<br>（Center for Math．and Comp．Sci．，The Netherlands）

## 1．Introduction

Let $l s_{n}$ be the Lie algebra of all differential operators in $n$ variables with polynomial coeffi－ cients of total degree in variables and derivatives $\leqslant 2$ ．Thus e．g．$l s_{1}$ is the Lie algebra with basis

$$
\begin{equation*}
x^{2}, x \frac{\partial}{\partial x}, \frac{\partial^{2}}{\partial x^{2}}, x, \frac{\partial}{\partial x}, 1 \tag{1.1}
\end{equation*}
$$

（The product is of course the commutator product）．The symbol $l s$ for this Lie algebra stands for＂linear systems＂．The reason for this appellation derives from the following．Consider a linear stochastic system

$$
\begin{equation*}
d x_{t}=A x_{t} d t+B d w_{t}, \quad d y_{t}=C x_{t}+d v_{t} \tag{1.2}
\end{equation*}
$$

Then an unnormalized version of the density of the conditional expectation of the state $x_{t}$ given the past observations $y_{s}, 0 \leqslant s \leqslant t$ ，satisfies a（stochastic）evolution equation

$$
\begin{equation*}
d \rho(x, t)=L \rho(x, t) d t+L_{1} \rho(x, t) d y_{1 t}+\cdots+L_{p} \rho(x, t) d y_{p t} \tag{1.3}
\end{equation*}
$$

with $L, L_{1}, \cdots, L_{p} \in l s_{n}$ ．And for varying systems（1．2）these operators generate all of $l s_{n}$ ．
The Kalman－Bucy filter for $\hat{x}_{t}=E\left[x_{t} \mid y_{s}, 0 \leqslant s \leqslant t\right]$ is a system of the form

$$
\begin{equation*}
d z=\alpha(z) d t+\beta_{1}(z) d y_{1 t}+\cdots+\beta_{p}(z) d y_{p t} \tag{1.4}
\end{equation*}
$$

where $z$ is short for $(p, \hat{x})$ and $\alpha, \beta_{1}, \cdots, \beta_{p}$ are vectorfields on（ $P, \hat{x}$ ）－space．Let $V\left(\underline{\underline{R}}^{N}\right)$ denote the Lie algebra of vectorfields on $\underline{\underline{R}}^{N}$ ．Then the first main point of this paper is that all Kalman－Bucy filters combine to define a＂universal Kalman－Bucy filter＂in the shape of an anti－homomorphism of Lie algebras

$$
\begin{equation*}
\kappa: l s_{n} \rightarrow V\left(\underline{\underline{R}}^{N}\right), N=\frac{1}{2} n(n+1)+n \tag{1.5}
\end{equation*}
$$

（and it is even possible to use this to propagate nongaussian initial densities）．Here＂anti＂ means that

$$
\kappa\left[D, D^{\prime}\right]=\left[\kappa\left(D^{\prime}\right), \kappa(D)\right] \text { rather than } \kappa\left[D, D^{\prime}\right]=\left[\kappa(D), \kappa\left(D^{\prime}\right)\right]
$$

This also establishes that the Kalman filter does indeed define an antihomomorphism of Lie algebras
from the Lie algebra generated by $L, L_{1}, \cdots, L_{p}$ in (1.3) (the so-called estimation Lie algebra) to a suitable Lie algebra of vectorfields, as it should according to a philosophy (almost a theorem now) first proposed by Brockett and Clark [1].

The structure of $l s_{n}$ is simple. It is an extension of the real symplectic Lie algebra $S p_{n}$ by the Heisenberg Lie algebra $h_{n}$. Let $S p_{n}$ be the symplectic Lie group. Then there is a famous and somewhat mysterious representation of $S p_{n}$ (or more precisely its 2 -fold covering $\tilde{S} p_{n}$ ) which turns up in many distinct areas of mathematics, e.g. number theory and quantum mechanics. It is called the Segal-Shale-Weil representation or sometimes the oscillator representationThe second main point of this paper is that this Segal-Shale-Weil representation and the "filter anti-representation ${ }^{3}$ (1.5) above are intimately related. This extends and strengthens the links between filtering theory and quantum mechanics which had been noted before [11], cf. also various contributions in [5].

It seems likely that the fact that all Kalman-Bucy filters fit together nicely will be useful both for theory and applications. In fact it is definitely of importance in a class of nonlinear filtering problems coming from identification and tracking [4, 10] where the estimation Lie algebra is a lways a subalgebra of a current algebra $l s_{n} \otimes R$ where $R$ is a ring of polynomials. Further applications of the "universal filter" (1.5) and/or its relations with the" Segal-ShaleWeil representation seem likely.

## 2. The Linear Systems Lie Aigebra $l s_{n}$

2.1. Definition of $l s_{n}$. Let $n \in N$. If $\alpha$ is a multi-index $\alpha=\left(\alpha_{1}, \cdots, \alpha_{n}\right), \alpha_{i} \in \underline{\underline{N} U}$ $\{0\}$, then $|\alpha|$ denotes $\alpha_{1}+\cdots+\alpha_{n}$ and we write

$$
x^{\alpha}=x_{1}^{\alpha_{1}} \cdots x_{n^{\prime}}^{\alpha}, \quad \partial_{\beta}=\frac{\partial^{\beta}}{\partial x^{\beta}}=\frac{\partial^{\beta_{1}}}{\partial x_{1}^{\beta}} \cdots \frac{\partial^{\beta_{n}}}{\partial x_{n}^{\beta_{n}^{n}}} .
$$

With these notations $l s_{n}$ is by definition the Lie algebra of all differential operators of total degree $\leqslant 2$, i. e. all differential operators $\sum c_{\alpha, \beta} x^{\alpha} \partial_{\beta}$ with $c_{\alpha, \beta}=0$ unless $|\alpha|+|\beta| \leqslant 2$. These operators are considered to act on some suitable space of (real or complex valued) smooth functions on $\mathbb{R}^{n}$, say the Schwartz space $S\left(\mathbb{R}^{n}\right)$ of rapidly decreasing smooth functions on $\mathbb{R}^{n}$. The product (Lie bracket) of $D_{1}, D_{2}$ is then of course given by the commutator

$$
\left[D_{1}, D_{2}\right](\phi)=D_{1}\left(D_{2} \phi\right)-D_{2}\left(D_{1} \phi\right), \quad \phi \in S\left(\mathbb{R}^{n}\right)
$$

It is an elementary observation that $l s_{n}$ is closed under this commutator product.
I shall call $l s_{n}$ the linear systems Lie algebra. The reason for this name will become clear later (in section 4 below).
2.2. The Heisenberg Lie algebra $h_{n}$. Let $h_{n}$ be the subspace of $l s_{n}$ spanned by the operators of total degree $\leqslant 1$, i.e. the operators $x_{1}, \cdots, x_{n} ; \partial_{1}, \cdots, \partial_{n} ; 1$ (with an obvious notation). The products in $h_{n}$ are of course the Heisenberg commutation relations

$$
\begin{equation*}
\left[\partial_{i}, x_{j}\right]=\delta_{i j},\left[x_{i}, x_{j}\right]=\left[\partial_{i}, \partial_{j}\right]=\left[x_{i}, 1\right]=\left[\partial_{i}, 1\right]=0 \tag{2,3}
\end{equation*}
$$

where $\delta_{i j}$ is the Kronecker $\delta$. The Lie algebra $h_{n}$ is called the Heisenberg Lie algebra.
2.4. The symplectic Lie algebra $s p_{n}$. Let $J$ be the $2 n \times 2 n$ matrix

$$
J=\left(\begin{array}{rr}
0 & I \\
-I & 0
\end{array}\right)
$$

where $I$ stands for the $n \times n$ unit matrix. The Lie algebra $s p_{n}$ consists of all $2 n \times 2 n$ matrices $M$ which satisfy $M J+J M^{T}=0$ (where $M^{T}$ is the transpose of $M$ ). The product on $s p_{n}$ is the commutator matrix product $\left[M, M^{\prime}\right]=M M^{\prime}-M^{\prime} M$.
2.5. Structure of $l s_{n}$. It is an easy observation that $h_{n} \subset l s_{n}$ is an ideal, i.e. $\left[D, D^{\prime}\right] \in$ $h_{n}$ for all $D \in l s_{n}, D^{\prime} \in h_{n}$. The quotient Lie algebra $l s_{n} / h_{n}$ is isomorphic to $s p_{n}$. This can e. g. be seen as follows. Let $E_{i, j}$ denote the matrix with a 1 at spot ( $i, j$ ) and 0 everywhere else. Then the homomorphism of vectorspaces defined by

$$
\begin{aligned}
x_{i} x_{i} & \rightarrow E_{i, n+j}+E_{j, n+i}, \quad i, j=1, \cdots, n \\
x_{i} \frac{\partial}{\partial x_{j}} & \rightarrow E_{i, i}-E_{n+j, n+i}, \quad i, j=1, \cdots, n \\
\frac{\partial^{2}}{\partial x_{i} \partial x_{j}} & \rightarrow E_{n+i, j}-E_{n+j, i}, \quad i, j=1, \cdots, n \\
h_{n} & \rightarrow 0,
\end{aligned}
$$

is a surjective homomorphism of Lie algebras as is easily checked and induces an isomorphism $l s_{n} / h_{n} \simeq s p_{n}$. Thus we have an exact sequence

$$
\begin{equation*}
0 \rightarrow h_{n} \xrightarrow{l} l s_{n} \xrightarrow{\pi} s p_{n} \rightarrow 0 \tag{2.6}
\end{equation*}
$$

A lift of $\pi$ (i. e. a homomorphism of Lie algebras $\sigma: s p_{n} \rightarrow l s_{n}$ such that $\pi \circ \sigma=$ id) is given by $\sigma\left(E_{i, n+j}+E_{i, n+i}\right)=x_{i} x_{j}, \quad \sigma\left(E_{n+i, i}-E_{n+i, i}\right)=-\frac{\partial^{2}}{\partial x_{i} \partial x_{j}}, \sigma\left(E_{i, j}-E_{n+j, n+i}\right)=$ $x_{i} \frac{\partial}{\partial x_{j}}+\frac{1}{2} \delta_{i j}$. This defines an action of $s p_{n}$ on $h_{n}$ and also on $h_{n} / Z \simeq \mathbb{R}^{2 n}$ (as an abelian Lie algebra) where $Z$ is the one dimensional centre of $h_{n}$ and $l s_{n}$. Identifying $\mathbb{R}^{2 n}$ with $h_{n} / Z$ by means of $e_{i} \rightarrow x_{i}, e_{n+i} \rightarrow-\partial_{i}, i=1, \cdots, n$, this action becomes the usual action of $s p_{n}$ as a Lie algebra of $2 n \times 2 n$ matrices on $\mathbb{R}^{2 n}$.

## 3. The Filter Anti-Representation of $l s_{n}$

3.1. Description of the anti-representation. If $M$ is a smooth manifold, $F(M)$ denotes the smooth functions on $M$ and $V(M)$ denotes the Lie algebra of vectorfields on $M$ (considered as the Lie algebra of derivations $F(M) \rightarrow F(M)$ ). If $M=\mathbb{R}^{n}$ then in the coordinates $\left(x_{1}\right.$, $\cdots, x_{n}$ ) every vectorfield on $\mathbb{R}^{n}$ can be written as $\sum_{i} f_{i}(x) \frac{\partial}{\partial x_{i}}$, where the $f_{i}(x)$ are smooth functions.

Now consider $\mathbb{R}^{N}$ with $N=\frac{1}{2} n(n+1)+n+1$ with coordinates $P_{i j}=P_{i j}, i, j=1$, $\cdots, n ; m_{i}, i=1, \cdots, n$. Consider the hemcmorphism of real vectorspaces
defined by the formulas

$$
\begin{gather*}
\kappa: l s_{n} \rightarrow V\left(\underline{\underline{R}}^{N}\right)  \tag{3.2}\\
1 \rightarrow \frac{\partial}{\partial c}  \tag{3.3}\\
x_{i} \rightarrow m_{i} \frac{\partial}{\partial c}+\sum_{t=1}^{n} P_{i t} \frac{\partial}{\partial m_{t}} \tag{3.4}
\end{gather*}
$$

$$
\begin{gather*}
\frac{\partial}{\partial x_{i}} \rightarrow-\frac{\partial}{\partial m_{i}},  \tag{3.5}\\
x_{i} x_{j} \rightarrow\left(m_{i} m_{i}+P_{i j}\right) \frac{\partial}{\partial c}+\sum_{t}\left(m_{i} P_{j t}+m_{j} P_{i t}\right) \frac{\partial}{\partial m_{t}} \\
+\sum_{s, t} P_{i s} P_{i t} \frac{\partial}{\partial P_{s t}}+\sum_{t} P_{i t} P_{i t} \frac{\partial}{\partial P_{t t}},  \tag{3.6}\\
x_{i} \frac{\partial}{\partial x_{j}} \rightarrow-m_{i} \frac{\partial}{\partial m_{j}}-\delta_{i j} \frac{\partial}{\partial c}-P_{i j} \frac{\partial}{\partial P_{j j}}-\sum_{t} P_{i t} \frac{\partial}{\partial P_{j t}},  \tag{3.7}\\
\frac{\partial^{2}}{\partial x_{i} \partial x_{j}} \rightarrow \frac{\partial}{\partial P_{i j}} \text { if } i \neq j, \frac{\partial^{2}}{\partial x_{i}^{2}} \rightarrow 2 \frac{\partial}{\partial P_{i i}} \tag{3.8}
\end{gather*}
$$

3.9. Theorem. The vectorspace homomorphism $\kappa: l s_{n} \rightarrow V\left(\mathbb{R}^{N}\right)$ defined by the formulae (3.3)-(3.8) is an injective anti-homomorphism of Lie algebras (i.e. it satisfies $\kappa\left[D, D^{\prime}\right]=$ $\left[\kappa\left(D^{\prime}\right), \kappa(D)\right]$ for all $\left.D, D^{\prime} \in l s_{n}\right)$.

The proof of this theorem is a straightforward but perhaps somewhat tedious calculation. As an example we have $\left[\partial_{i}, x_{i}\right]=1$ and

$$
\left[m_{i} \frac{\partial}{\partial c}+\sum_{t} P_{i t} \frac{\partial}{\partial m_{t}},-\frac{\partial}{\partial m_{i}}\right]=\frac{\partial}{\partial c}
$$

which fits. As another example if $i \neq j$ we have

$$
\left[\frac{\partial^{2}}{\partial x_{i} \partial x_{j}}, x_{i} x_{j}\right]=x_{i} \frac{\partial}{\partial x_{i}}+x_{j} \frac{\partial}{\partial x_{j}}+1
$$

Now

$$
\begin{aligned}
& {\left[\frac{\partial}{\partial P_{i j}},\left(m_{i} m_{j}+P_{i j}\right) \frac{\partial}{\partial c}\right]=\frac{\partial}{\partial c}} \\
& {\left[\frac{\partial}{\partial P_{i j}}, \sum_{t}\left(m_{i} P_{j t}+m_{j} P_{i t}\right) \frac{\partial}{\partial m_{t}}\right]=m_{i} \frac{\partial}{\partial m_{i}}+m_{j} \frac{\partial}{\partial m_{j}}} \\
& {\left[\frac{\partial}{\partial P_{i j}}, \sum_{s, t} P_{i s} P_{i t} \frac{\partial}{\partial P_{s, t}}\right]=\sum_{i} P_{i t} \frac{\partial}{\partial P_{j t}}+\sum_{s} P_{i s} \frac{\partial}{\partial P_{s i}},} \\
& {\left[\frac{\partial}{\partial P_{i j}}, \sum_{t} P_{i t} P_{i t} \frac{\partial}{\partial P_{t t}}\right]=P_{j j} \frac{\partial}{\partial P_{j i}}+P_{i i} \frac{\partial}{\partial P_{i i}}}
\end{aligned}
$$

So indeed

$$
\begin{aligned}
{\left[\kappa\left(x_{i} x_{j}\right), \kappa\left(\frac{\partial^{2}}{\partial x_{i} \partial x_{i}}\right)\right]=} & -\frac{\partial}{\partial c}-m_{i} \frac{\partial}{\partial m_{i}}-m_{j} \frac{\partial}{\partial m_{j}}-\sum_{t} P_{i t} \frac{\partial}{\partial P_{j t}} \\
& -\sum_{s} P_{i s} \frac{\partial}{\partial P_{i s}}-P_{i j} \frac{\partial}{\partial P_{j j}}-P_{i i} \frac{\partial}{\partial P_{i i}} \\
& =\kappa\left(x_{i} \frac{\partial}{\partial x_{i}}\right)+\kappa\left(x_{i} \frac{\partial}{\partial x_{i}}\right) \kappa(1)
\end{aligned}
$$

The remaining identities are checked similarly.
3.10. Remark.

$$
\frac{\partial}{\partial x_{i}} \rightarrow-\frac{\partial}{\partial x_{i}}, x_{i} \frac{\partial}{\partial x_{j}} \rightarrow-x_{i} \frac{\partial}{\partial x_{i}}, x_{i} x_{j} \rightarrow x_{i} x_{j}, \frac{\partial^{2}}{\partial x_{i} \partial x_{j}} \rightarrow \frac{\partial^{2}}{\partial x_{i} \partial x_{i}}
$$

$x_{i} \rightarrow x_{i}, 1 \rightarrow 1$ defines an anti-automorphism of $l s_{n}$. Thus changing the sign in formulas (3.5) and (3.7) defines a representation of $\delta s_{m}$ in $V\left(\mathbb{R}^{N}\right)$.

## 4. DMZ Equations and Kalman Filters

4.1. The Duncan-Mortenson-Zakai equation and the estimation Lie algebra. Consider a general nonlinear stochastic system (in Ito form)

$$
\begin{equation*}
d x_{t}=f\left(x_{t}\right) d t+G\left(x_{i}\right) d w_{t}, \quad d y_{t}=h\left(x_{t}\right) d t+d v_{t}, \quad x_{t} \in R^{n}, \tag{4.2}
\end{equation*}
$$

where $f, G, h$ are suitable vector and matrix valued functions and $w_{t}$ and $v_{t}$ are independent unit covariance Wiener processes also independent of the initial random vector $x_{0}$. Given sufficiently nice $f, G, h$, an unnormalized version $\rho(x, t)$ of the probability density $p(x, t)$ of the state $x_{t}$ given the past observations $y_{s}, 0 \leqslant s \leqslant t$, satisfies the (forced) diffusion equation (FiskStratonovič form)

$$
\begin{equation*}
d \rho=L \rho d t+\sum_{j=1}^{D} h_{i} \rho d y_{t} \tag{4.3}
\end{equation*}
$$

where $h_{i}$ is the $j$-th component of $h$ and $L$ is the second order differential operator

$$
\begin{equation*}
L \phi=\frac{1}{2} \sum_{i, i} \frac{\partial^{2}}{\partial x_{i} \partial x_{j}}\left(\left(G G^{T}\right)_{i i} \phi\right)-\sum_{i} \frac{\partial}{\partial x_{i}}\left(f_{i} \phi\right)-\frac{1}{2} \sum_{i} h_{i}^{2} \phi . \tag{4.4}
\end{equation*}
$$

Here $f_{i}$ is the $i$-th component of $f$ and $\left(G G^{T}\right)_{i j}$ the $(i, j)$-entry of the matrix product $G G^{T}$. Equation (4.3) is called the Duncan-Mortensen-Zakai equation. Cf. e.g. [3] for a derivation. The Lie algebra of differential operators (on $S\left(\mathbb{R}^{n}\right)$ say) generated by $L$ and $h_{1}, h_{2}, \cdots, h_{p}$ is called the estimation Lie algebra.
4.5. Exact filters and Liealgebra anti-homomorphisms. Now let

$$
\begin{equation*}
d \xi_{t}=\alpha\left(\xi_{t}\right) d t+\beta_{1}\left(\xi_{t}\right) d y_{1 t}+\cdots+\beta_{p}\left(\xi_{t}\right) d y_{p t}, \hat{x}_{t}=\gamma\left(x_{t}\right) \tag{4.6}
\end{equation*}
$$

be a stochastic system (in Fisk-Stratonovič form) driven by $y_{z}$ which calculates the conditional expectation

$$
\begin{equation*}
\hat{x}_{t}=E\left[x_{t} \mid y_{s}, 0 \leqslant s \leqslant t\right] \tag{4.7}
\end{equation*}
$$

of the state given the past observations. I. e. (4.6) is a filter for $\hat{x}_{t}$. Then as Brockett and Clark observed [1] we have two ways of calculating $\hat{x}_{t}$, one via (4.3) and one via (4.6). Minimal realization theory then suggests that there will be a corresponding homomorphism of Lie algebras from the estimation Lie algebra $L$ of the system to the Lie algebra of vectorfields generated by the vectorfields $\alpha, \beta_{1}, \cdots, \beta_{p}$ in (4.6) given by $A \rightarrow \alpha, h_{i} \rightarrow \beta_{i}, i=1, \cdots, p$. This is called the Brockett-Clark homomorphism principle. In [1] this was verified to be indeed the case for the case of the Kalman filter of one of the simplest possible linear systems, namely $d x_{t}=d w_{i}, d y_{t}=x_{i} d t+d v_{i}$.

As a matter of fact a filter like (4.6) for $\hat{x}_{z}$ (or for some other statistic) should give rise to an anti-homomorphism from the Lie algebra of differential operators to the Lie algebra of vectorfields generated by the vectorfields in the filter. The reason is that $A \rho$ and $h_{i} \rho$ in (4.3) must be interpreted as vectorfields on $S\left(\mathbb{R}^{n}\right)$ and the mapping which assigns to a linear operator the corresponding linear vectorfield is an injective anti-homomorphism of Lie algebras.
4.8. The Kalman-Bucy filter. Now consider an $n$-dimensional linear system with $m$ inputs and $p$ outputs

$$
\begin{equation*}
(\Sigma) \quad d x_{t}=A x_{t} d t+B d w_{t}, \quad d y_{t}=C x_{t} d t+d v_{t} \tag{4.9}
\end{equation*}
$$

The Kalman-Bucy filter for $\hat{x}_{t}$ is given by the equations

$$
\begin{gather*}
d \hat{x}_{t}=A \hat{x}_{t} d t+P_{t} C^{T}\left(d y_{t}-C \hat{x}_{t} d t\right)  \tag{4.10}\\
d P_{t}=\left(A P_{t}+P_{t} A^{T}+B B^{T}-P_{t} C^{T} C P_{t}\right) d t \tag{4.11}
\end{gather*}
$$

Write $m_{\imath}$ for $\hat{x}_{i}$ and $P_{t}=\left(P_{i j}\right)$. Then the part of the right hand side of (4.10) involving $d y_{k t}$ and contributing to $d m_{i t}$ is equal to

$$
\sum_{j} P_{i j} c_{k j} d y_{k t}
$$

It follows that if we write (4.10), (4.11) in the form (4.6) then the vectorfields $\beta_{1}, \ldots$, $\beta_{p}$ are equal to

$$
\begin{equation*}
\beta_{k}=\sum_{r, r} P_{r s} c_{k s} \frac{\partial}{\partial m_{r}}, \quad k=1, \cdots, p \tag{4.12}
\end{equation*}
$$

Similarly the $\alpha$ vectorfield of (4.10)-(4.11) is equal to

$$
\begin{align*}
\alpha= & \sum_{i, j} a_{i j} m_{j} \frac{\partial}{\partial m_{i}}-\sum_{i, j, r, s} P_{i j} c_{r j} c_{r s} m_{s} \frac{\partial}{\partial m_{i}} \\
& +\sum_{r, i \leqslant j} a_{i r} P_{r i} \frac{\partial}{\partial P_{i j}}+\sum_{r, i \leqslant j} P_{i r} a_{i r} \frac{\partial}{\partial P_{i j}}+\sum_{r, i \leqslant j} b_{i r} b_{i r} \frac{\partial}{\partial P_{i j}} \\
& -\sum_{r, s, t, i \leqslant j} P_{\imath r} c_{s r} c_{s t} P_{t j} \frac{\partial}{\partial P_{i i}} . \tag{4.13}
\end{align*}
$$

4.14. Estimation Lie algebra and Kalman-Bucy filter. Consider again the linear system (4.9). The operators which occur in the DMZ equation for this system are

$$
\begin{gather*}
h_{i}=\sum_{r} c_{i r} x_{r}  \tag{4.15}\\
L=\frac{1}{2} \sum_{i, j, r} b_{i r} b_{i r} \frac{\partial^{2}}{\partial x_{i} \partial x_{j}}-\sum_{r, i} a_{i r} x_{r} \frac{\partial}{\partial x_{i}}-\frac{1}{2} \sum_{i, j, r} c_{r i} c_{r j} x_{i} x_{i}-\sum_{i} a_{i i} \tag{4.16}
\end{gather*}
$$

Let $L(\Sigma)$ be the estimation Lie algebra of the linear system (4.9). This is obviously, cf. (4.16), a sub-Lie algebra of $l s_{n}$, and for varying $\Sigma$ the various $L(\Sigma)$ generate all of $l s$. Whence the name "linear systems Lie algebra" for $l s_{n}$.

As in section 3 above let $N=\frac{1}{2} n(n+1)+n+1$. Consider the projection $\mathbb{R}^{N} \rightarrow$ $\mathbb{R}^{N-1}$ which maps $(m, P, c)$ to ( $m, P$ ). Under this projection the vectorfields occurring in the right hand sides of (3.3)-(3.8) map to vectorfields on $\mathbb{R}^{N-1}$. The vectorfields arising in this way are the same ones except that the $\frac{\partial}{\partial_{c}}$ terms are removed. Let

$$
\begin{equation*}
\kappa^{\prime}: l s_{n} \rightarrow V\left(\underline{\underline{R}}^{N-1}\right) \tag{4.17}
\end{equation*}
$$

be the resulting anti-homomorphism of Lie algebras.
4.18. Theorem. The restriction of $\kappa^{\prime}$ to $L(\Sigma)$ maps the operator $L$ of (4.16) to the vectorfield $\alpha$ of (4.13) and the operators $h_{i}$ of (4.15) to the vectorfields $\beta_{i}$ of (4.12). In other words the restriction of $\kappa^{\prime}$ to $L(\Sigma) \subset l s_{n}$ is the Kalman-Bucy filter for the system ( $\Sigma$ ).

The proof of theorem 4.18 is an entirely straightforward verification, lightly complicated
by the fact that $P_{i j}=P_{i i}$ must be taken into account which is not automatically done by the notation used. Thus the coefficient of $\frac{\partial^{2}}{\partial x_{i} \partial x_{i}}$ in $L$ in (4.16) is equal to $\sum_{r} b_{i r} b_{i r}$ if $i \neq j$ and $\frac{1}{2} \Sigma b_{i r}^{2}$ if $i=j$, and under $\kappa^{\prime}$ which takes

$$
\frac{\partial^{2}}{\partial x_{i} \partial x_{j}} \rightarrow \frac{\partial}{\partial P_{i j}}, \frac{\partial^{2}}{\partial x_{i}^{2}} \rightarrow 2 \frac{\partial}{\partial P_{i i}},
$$

this gives the fifth term of $\alpha$ in (4.13). Similarly the coefficient of $x, \frac{\partial}{\partial x_{i}}$ in $(4.16)$ is $-a_{i r}$. The morphism $\kappa^{\prime}$ takes

$$
x_{r} \frac{\partial}{\partial x_{i}} \rightarrow-m_{r} \frac{\partial}{\partial m_{i}}-P_{r i} \frac{\partial}{\partial P_{i i}}-\sum_{i} P_{r i} \frac{\partial}{\partial P_{i t}}
$$

and these terms account for the first, third and fourth terms in (4.13). Finally the coefficient of $x_{i} x_{i}$ in (4.16) is $-\sum_{r} c_{r i} c_{r i}$ if $i \neq j$ and $-\frac{1}{2} \sum c_{r i}^{2}$ if $i=j$. The morphism $\kappa^{\prime}$ takes $x_{i} x_{j}$ into

$$
\sum_{i}\left(m_{i} P_{i t}+m_{i} P_{i t}\right) \frac{\partial}{\partial m_{t}}+\sum_{t, t} P_{i s} P_{i t} \frac{\partial}{\partial P_{s t}}+\sum_{i} P_{i t} P_{j t} \frac{\partial}{\partial P_{t t}}
$$

and this accounts for the second and sixth terms in (4.13). Similarly (and rather easier) one checks that $\kappa^{\prime}$ takes the $h_{i}$ of(4.15)into the $\beta_{i}$ of (4.12). This proves that $\kappa^{\prime}$ indeed restricts to the Kalman-Bucy filter on $L(\Sigma)$.
4.19. Remarks. Another way to state theorem 4.18 is to say that all possible KalmanBucy filters combine to define an anti-representation of $l s_{n}$ which is faithful modulo the onedimensional centre. The lifted anti-representation $\kappa$ is faithful on $l s_{n}$ itself and permits us to propagate also nongaussian initial densities. Cf. also section 6 below.

As a corollary of theorem 4.18 we of course obtain that $L \rightarrow \alpha, h_{i} \rightarrow \beta_{z}$ (with $L, \alpha$, $h_{i}, \beta_{i}$ respectively given by (4.16), (4.13), (4.15), (4.12)) does indeed define an antihomomorphism of Lie algebras,as it should.
4.20. Identification as a nonlinear filtering problem. Consider a linear system

$$
\begin{equation*}
d x=A x d t+B d w, d y=C x d t+d v_{t} \tag{4.21}
\end{equation*}
$$

in which $A, B, C$ are unknown. The problem is to find the best estimates of both $x_{t}$ and the matrices $A, B, C$ given $y_{s}, 0 \leqslant s \leqslant t$. By adding to (4.21) the state equations

$$
\begin{equation*}
d A=0, d B=0, d C=0 \tag{4.22}
\end{equation*}
$$

(so that $A, B, C$ are viewed as random variables (constant in time)), we obtain a nonlinear system and the nonlinear filtering problem of finding the conditional expectation of the extended state $(x, A, B, C)$ is the identification of linear system problem (or at least one version of it).

One potentially interesting statistic is the conditional expectation of the random variable $x_{t} \mid(A, B, C)$. The family of all Kalman filters (4.10)-(4.11) (for varying ( $A, B, C$ )) computes this. According to the Brockett-Clark anti-homomorphism principle, there should be a corresponding anti-homomorphism of Lie algebras. This is the morphism $\kappa^{\prime}$ when the image is
viewed as vectorfields on $\mathbb{R}^{N-1} \times \mathbb{R}^{n^{2}+n p+n m}$ with no $\frac{\partial}{\partial A}, \frac{\partial}{\partial B}$ or $\frac{\partial}{\partial C}$ terms. Thus the BrockettClark principle also holds in this case.
4.23. Example. For special linear systems $L \rightarrow \alpha, h_{i} \rightarrow \beta_{i}$ may accidentally also define a homomorphism of Lie algebras. This happens e.g. for all one-dimensional systems and all systems (4.9) for which the $A$ matrix is zero. In general this is not the case as the following example shows

$$
\begin{aligned}
& \binom{d x_{1}}{d x_{2}}=\left(\begin{array}{ll}
0 & 1 \\
0 & 0
\end{array}\right)\binom{x_{1}}{x_{2}} d t+\binom{0}{1} d w \\
& \binom{d y_{1}}{d y_{2}}=\left(\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right)\binom{x_{1}}{x_{2}} d t+\binom{d v_{1}}{d v_{2}}
\end{aligned}
$$

## 5. The Segal-Shale-Weil Representation

This section simply lists some well-known facts on the basis of [7] with a few elaborations. 5.1. The symplectic group. Let $J$ be as in 2.4 above. Then the symplectic group $S p_{n}$ consists of all real $2 n \times 2 n$ matrices $M$ such that $M J M^{T}=J$. The Lie algebra of $S p_{n}$ is the Lie algebra $s p_{n}$ which we encountered in section 2 above.

A certain representation of $S p_{n}$ or more precisely of its two-fold covering $\tilde{S} p_{n}$ on $L^{2}\left(\mathbb{R}^{n}\right)$ which is called the Segal-Shale-Weil representation is of considerable importance in several areas of mathematics, notably number theory [14] and quantum mechanics [12,13]. As we shall see it is also closely related to all Kalman-Bucy filters.
5.2. Definition of the Segal-Shale-Weil representation. One well-known way to obtain this representation is via the Stone-Von Neumann uniqueness theorem. Let $H_{n}$ denote the Heisenberg group, $H_{n}=\mathbb{R}^{n} \times \mathbb{R}^{n} \times S^{1}$, where $S^{1}$ is the circle, with the multiplication $(x, y, z)\left(x^{\prime}, y^{\prime}\right.$, $\left.z^{\prime}\right)=\left(x+x^{\prime}, y+y^{\prime}, e^{-2 \pi i\langle x, y \prime\rangle} z z^{\prime}\right)$. The Lie algebra of $H_{n}$ is $h_{n}$ (which we also encountered in section 2 above). This Lie algebra cän also be described as $h_{n}=\mathbb{R}^{n} \times \mathbb{R}^{n} \times \mathbb{R}$ and then the Lie bracket defines a bilinear form $\mathbb{R}^{2 n} \times \mathbb{R}^{2 n} \rightarrow \mathbb{R}$ which is given by the matrix J. Thus $S p_{n}$ can be seen as a group of automorphisms of $h_{n}$ and $H_{n}$ which moreover is the identity on the centre $S^{1} \subset H_{n}$.

One version of the Stone-Von Neumann theorem says that up to unitary equivalence there is a unique irreducible representation of $H_{n}$ whose character on $S^{1}$ is the identity. Now let $\rho$ be the standard (Schrödinger) representation of $H_{n}$ in $L^{2}\left(\mathbb{R}^{n}\right)$ which is given by

$$
\begin{aligned}
& (x, 0,0) \rightarrow M_{x}, M_{x} f\left(x^{\prime}\right)=e^{2 \pi i\left\langle x, x^{\prime}\right\rangle} f\left(x^{\prime}\right), \\
& (0, y, 0) \rightarrow T_{y}, T_{y} f\left(x^{\prime}\right)=f\left(x^{\prime}-y\right) \\
& (0,0, z) \rightarrow S_{z}, S_{z} f\left(x^{\prime}\right)=z f\left(x^{\prime}\right) .
\end{aligned}
$$

Now let $g \in S p_{n}$ and consider $S p_{n}$ as a group of automorphisms of $H_{n}$. Then $h \rightarrow \rho(g(h))$ is also an irreducible representation of $H_{n}$ with the same central character. By the uniqueness theorem there is an intertwinning operator $\omega(g)$ such that $\omega(g) \rho(h) \omega(g)^{-1}=\rho(g(h))$. These $\omega(g)$ are unique up to a scalar factor. It remains to see whether these scalar factors can be fixed up to yield a representation of $S p_{n}$ on $L^{2}\left(\mathbb{R}^{n}\right)$ (instead of on $P\left(L^{2}\left(\mathbb{R}^{n}\right)\right)$ ). This can almost be done and the result is the Segal-Shale-Weil representation of the two-fold covering
$\tilde{s} p_{n}$ of $S p_{n}$ in $L^{2}\left(\mathbb{R}^{n}\right)$.
5.3. More or less explicit description of the Segal-Shale-Weil representation. Let

$$
M=\left(\begin{array}{ll}
A & B  \tag{5.4}\\
C & D
\end{array}\right) \in S p_{n},
$$

where $A, B, C, D$ are $n \times n$ matrices. Then the $A, B, C, D$ satisfy $A B^{T}=B A^{T}, C D^{T}=$ $D C^{T}, A D^{T}-B C^{T}=1$. Important special elements in $S p_{n}$ are

$$
\left(\begin{array}{rr}
0 & I  \tag{5.5}\\
-I & 0
\end{array}\right),\left(\begin{array}{cc}
A & 0 \\
0 & \left(A^{-1}\right)^{r}
\end{array}\right),\left(\begin{array}{cc}
I & N \\
0 & I
\end{array}\right), N \text { symmetric }
$$

and it is not especially difficult to show that these generate all of $S p_{n}$. Thus in principle to describe the Segal-Shale-Weil representation it suffices to describe the unitary operators corresponding to these matrices. These are as follows

$$
\begin{gather*}
\left(\begin{array}{cc}
0 & I \\
-I & 0
\end{array}\right) \rightarrow \text { Fourier transform } F: L^{2}\left(\underline{\underline{R}}^{n}\right) \rightarrow L^{2}\left(\underline{\underline{R}}^{n}\right)  \tag{5.6}\\
\left(\begin{array}{cc}
A & 0 \\
0 & \left(A^{-1}\right)^{T}
\end{array}\right) \rightarrow\left(f(x) \rightarrow|\operatorname{det} A|^{\frac{1}{2}} f\left(A^{T} x\right)\right)  \tag{5.7}\\
\left(\begin{array}{cc}
I & N \\
0 & I
\end{array}\right) \rightarrow\left(f(x) \rightarrow e^{x i N(x)} f(x)\right) \tag{5.8}
\end{gather*}
$$

where $N(x)$ is the quadratic form defined by the symmetric matrix $N$.
5.9. The Lie algebra representation defined by the Segal-Shale-Weil representation. First consider a symmetric matrix $N=\left(n_{i j}\right)$. Then

$$
\left.\frac{d}{d t}\left(\left(\begin{array}{cc}
I & 0 \\
t N & I
\end{array}\right) f(x)\right)\right|_{t=0}=(\pi i N(x))(f(x)) .
$$

Next let $B$ be an $n \times n$ matrix, $A=e^{t B}$. Then

$$
\left.\frac{d}{d t}\left(\begin{array}{cc}
e^{t B} & 0 \\
0 & \left(e^{-t B}\right)^{T}
\end{array}\right)(f)\right|_{t=0}=\left(+\frac{1}{2} \operatorname{Tr}(B)+\sum_{i}\left(B^{T} x\right), \frac{\partial}{\partial x_{i}}\right)(f) .
$$

Finally consider the one-parameter subgroup

$$
S_{t}=\left(\begin{array}{rr}
I \cos t & I \sin t \\
-I \sin t & I \cos t
\end{array}\right)
$$

of $S p_{n}$ whose tangent vector at $t=0$ is $J$ (and which also passes through $J$ ). Writing

$$
\begin{aligned}
S_{t}= & \left(\begin{array}{cc}
0 & I \\
-I & 0
\end{array}\right)\left(\begin{array}{cc}
I \cos ^{-1} t & 0 \\
0 & I \cos t
\end{array}\right)\left(\begin{array}{cc}
I & I \sin t \cos t \\
I & 0
\end{array}\right) \\
& \left(\begin{array}{cc}
0 & -I \\
I & 0
\end{array}\right)\left(\begin{array}{cc}
I & I \operatorname{tg} t \\
0 & I
\end{array}\right)
\end{aligned}
$$

it is not difficult to write down $\left(S_{t} f(x)\right.$ and to calculate the derivative at $t=0$. The
result is result is

$$
\pi i\left(x_{1}^{2}+\cdots+x_{n}^{2}\right)-\frac{i}{4 \pi}\left(\frac{\partial^{2}}{\partial x_{1}^{2}}+\cdots+\frac{\partial^{2}}{\partial x_{n}^{2}}\right)
$$

No. 2

It readily follows that the Lie algebra of operators arising from the Segal-Shale-Weil representation is the one with basis

$$
\pi i x_{k} x_{j}, \frac{i}{4 \pi} \frac{\partial^{2}}{\partial x_{k} \partial x_{j}}, x_{k} \frac{\partial}{\partial x_{j}}+\frac{1}{2} \delta_{k j}
$$

which is of course isomorphic to $s p_{n}$, for example to the incarnation of $s p_{n}$ as the subalgebra $\sigma\left(s p_{n}\right) \subset l s_{n}$ via the isomorphism induced by the coordinate change $x_{k} \rightarrow(\sqrt{\pi i}) x_{k}$.

## 6. Kalman-Bucy Filters and the Segal-Shale-Weil Representation

6.1. Outline of the connection. Given that the Kalman-Bucy filters combine to give an antirepresentation of $s p_{n} \subset l s_{n}$ with $s p_{n}$ realized as a Lie algebra of differential operators and that the differentiated version of the Segal-Shale-Weil representation is also a representation of this same Lie algebra of differential operators, it would be odd if they were not rather closely related. Indeed, as the attentive reader will have seen coming, the filter anti-representation is essentially a real and local version of the Segal-Shale-Weil representation.

The connection is essentially given by assigning to a pair $(P, m), m \in \mathbb{R}^{n}, P$ a symmetric positive definite matrix, the corresponding normal density

$$
\begin{equation*}
-\frac{1}{\sqrt{(2 \pi)^{n}|P|}} e^{-\frac{1}{2} P^{-1}(x-m)} \tag{6.2}
\end{equation*}
$$

where $|P|$ is the absolute value of the determinant of $P$ and $P^{-1}(y)$ is the quadratic form defined by $P^{-1}$. These functions form a total system in $L^{2}\left(\mathbb{R}^{n}\right)$ meaning that the finite linear combinations are dense, so that to define a representation of say $S p_{n}$ of $L^{2}\left(\mathbb{R}^{n}\right)$ it suffices to know what the representation does on these special functions. For the Segal-Shale-Weil representation one uses more generally $n \times n$ matrices $Q$ whose real part is positive definite. And in fact it seems that Weil originally constructed his representation essentially in this way (cf. his comments, also referenced under [14], on the paper in question).

To spell things out in more detail and to avoid equations in $P^{-1}$ (and calculating trouble) it is useful to use the Fourier transform.
6.3. Some Fourier transform facts. Let $F$ denote the Fourier transform. Then we need the following more or less well-known facts

$$
F=\frac{\partial}{\partial x_{k}}=2 \pi i x_{k} F, F x_{k}=-\frac{1}{2 \pi i} \frac{\partial}{\partial x_{k}} F
$$

where e.g. $F x_{k}$ stands for the composition of the operator "multiplication with $x_{k}$ " with the operator $F$. The second fact we need is the formula

$$
\begin{equation*}
F^{-1}\left(\frac{1}{\sqrt{(2 \pi)^{n}|P|}} e^{-\frac{1}{2} P^{-1}(x-m)+c}\right)=e^{c+\langle 2 \pi i m, x\rangle-2 \pi^{2} P(x)} \tag{6.5}
\end{equation*}
$$

Finally it is useful to note that the set of all functions of the form

$$
\begin{equation*}
p(x) e^{-Q(x)} \tag{6.6}
\end{equation*}
$$

where $p(x)$ is a (complex) polynomial and $Q$ a (complex) polynomial of degree 2 whose real homogeneous part of degree 2 is positive definite, is stable under the Fourier transform, multiplication with polynomials and partial differentiation with respect to $x_{k}$.
6.7. Obtaining the filter anti-representation. Consider a function of the type $e^{c+\langle 2 \pi i m, x\rangle}$. ( $m$ and $P$ real). Imagine that $m$ and $P$ vary with time and try to see what this invap an evolution equation of the type

$$
\frac{\partial}{\partial t}\left(e^{c+\langle 2 \pi i m, x\rangle-2 \pi^{2} P(x)}\right)=L e^{c+\langle 2 \pi i m, x\rangle-2 \pi^{2} P(x)}
$$

where $L$ is a differential operator from $l s_{n}$. As is easy to see this yields a system of $a_{k}$ differential equations for $m_{k}$ and $P_{r s}$ provided that $L$ is in $l s_{n}$. This idea is also due ckett. These first order differential equations which have polynomial right hand sides least locally uniquely solvable.

As one of the most complicated examples for $n=2$ consider $\frac{\partial^{2}}{\partial x_{1} \partial x_{2}} \in l s_{n} . \quad \hbar$ $\exp (-)$ for the $e$-power in (6.8) we find

$$
\begin{gathered}
\frac{\partial}{\partial t} \exp (-)=\exp (-)\left(\dot{c}+2 \pi i \dot{m}_{1} x_{1}+2 \pi i \dot{m}_{2} x_{2}-2 \pi^{2}\left(\dot{P}_{11} x_{1}^{2}+2 \dot{P}_{12} x_{1} x_{2}+\dot{P}_{22} x_{2}^{2}\right)\right)_{\nu} \\
\frac{\partial^{2}}{\partial x_{1} \partial x_{2}} \exp (-)=\exp (1)\left[-4 \pi^{2} m_{1} m_{2}-8 \pi^{3} i m_{2} x_{1} P_{11}-8 \pi^{3} i m_{2} x_{2} P_{12}\right. \\
-8 \pi^{3} i m x_{1} x_{2} P_{22}+16 \pi^{4} x_{1} P_{11} x_{2} P_{22}+16 \pi^{4} x_{2}^{2} P_{12} P_{22}-8 \pi^{3} i m_{1} x_{1} P_{12} \\
\left.+16 \pi^{4} x_{1}^{2} P_{11} P_{12}+16 \pi^{4} x_{1} x_{2} P_{12}^{2}-4 \pi^{2} P_{12}\right] .
\end{gathered}
$$

Comparing these two expressions yields the differential equations

$$
\begin{aligned}
\dot{c} & =-4 \pi^{2} m_{1} m_{2}-4 \pi^{2} P_{12} \\
2 \pi i \dot{m}_{1} & =-8 \pi^{3} i m_{2} P_{11}-8 \pi^{3} i m_{1} P_{12} \\
2 \pi i \dot{m}_{2} & =-8 \pi^{3} i m_{2} P_{12}-8 \pi^{3} i m_{1} P_{22} \\
-2 \pi^{2} \dot{P}_{11} & =16 \pi^{4} P_{11} P_{12} \\
-2 \pi^{2} \dot{P}_{22} & =16 \pi^{4} P_{12} P_{22} \\
-4 \pi^{2} \dot{P}_{12} & =16 \pi^{4} P_{11} P_{22}+16 \pi^{4} P_{12}^{2}
\end{aligned}
$$

Writing down the associated vectorfield and using (6.4) and (6.5) the result is that th evolution of an unnormalized normal probability density $e^{c} N(m, P)$ with mean $m$ and c ance $P$ in an evolution equation

$$
\frac{\partial}{\partial t} e^{c} N(m, P)=x_{1} x_{2} e^{c} N(m, P)
$$

is given by

$$
\frac{\partial}{\partial t}(c, m, P)=\alpha(c, m, P)
$$

where $\alpha$ is the vectorfield

$$
\begin{aligned}
& \left(m_{1} m_{2}+P_{12}\right) \frac{\partial}{\partial c}+\left(m_{2} P_{11}+m_{1} P_{12}\right) \frac{\partial}{\partial m_{1}}+\left(m_{2} P_{12}+m_{1} P_{22}\right) \frac{\partial}{\partial m_{2}} \\
& \quad+2 P_{11} P_{12} \frac{\partial}{\partial P_{11}}+2 P_{12} P_{22} \frac{\partial}{\partial P_{22}}+\left(P_{11} P_{22}+P_{12}^{2}\right) \frac{\partial}{\partial P_{12}}
\end{aligned}
$$

which is of course the special case $n=2$ of formula (3.6).
6.12. Obtaining the Segal-Shale-Weil representation. To obtain the Segal-Shale-Weil representation one can proceed in almost precisely the same way. Now of course one admits complex $m$ and $P$ (with the real part of $P$ positive definite) and one uses
$i \frac{\partial}{\partial x_{j} \partial x_{k}}, i x_{j} x_{k}$ instead of $x_{i} x_{k}, \frac{\partial^{2}}{\partial x_{i} \partial x_{k}}$.
6.13. Finite escape time. The class of functions $e^{c+\langle 2 \pi i m, x\rangle-2 \pi^{2} P(x)}$ is stable under Fourier transform, multiplication with $e^{i Q(x)}, Q$ a real quadratic form and under $x \rightarrow A x, A$ invertible, i.e. they are stable under the transformations corresponding to the special elements (5.5) of $S p_{n}$. As these elements generate $S p_{n}$ it follows that there will be no finite escape time phenomena for the equations of the Segal-Shale-Weil case analogous to (6.10).

In the real case, i. e. the Kalman-Bucy filter case this can not be guaranteed. Indeed finite escape time does occur (cf. also [9]) and it is easy to see why. In this case $\left(\begin{array}{cc}I & N \\ 0 & I\end{array}\right)$ acts on $f(x)$ by multiplication with $e^{N(x)}$ and depending on $f(x)$ this may or may not result in a function $e^{N(x)} f(x)$ which is not Fourier transformable.

Writing elements of $S p_{n}$ as products of the special elements (5.5) gives more or less explicit solutions of Riccati equations for elements not too far from the identity and this also gives a good deal of information about in what directions (of $s p_{n}$ or $l s_{n}$ ) finite escape time phenomena do not occur. Of course the one parameter subgroups of $L S_{n}$ (the Lie group of $l s_{n}$ ) involve many more directions than those defined by "classically" studied Riccati equations. For complex linear systems the " $S p_{n}$ representation directions" are such that no finite escape time occurs either backwards or forwards. I do not know if this has system theoretic implications.

## References

[1] Brockett, R. W., Clark, J. M. C., The geometry of the conditional density equations, Proc. of the Oxford on Stochastic Systems, Oxford, 1978.
[ 2 ] Brockett, R. W., Geometrical methods in stochastic control and estimation, In: M. Hazewinkel, J. C. Willems (eds), Stochastic Systems: the mathematics of filtering and identification and applications, Reidel Publ. Cy., 1981, 441-478.
[ 3 ] Davis, M. H. A., Marcus, S. I., An introduction to nonlinear filrering, In: Hazewinkel, M., Willems J. C., (eds), Stochastic systems: the mathematics of filtering and identification and applications, Reidel Publ. Cy., 1981, 53-76.
[ 4 ] Honzon, B., Hazewinkel, M., On identification of linear system and the estimation Lie algebra of the associated nonlinear filtering problem. In: Proc. 6-th IFAC Conf. on identifications, Washington, June 1982, Pergamon Press, 1983, 63-68.
[5] Hazewinkel, M., Willems, J. C. (eds), Stochastic systems: the mathematics of filtering and identification and applications, Reidel Publ. Cy., 1981.
[6] Hazewinkel, M., Marcus, S. I., On Lie algebras and nonlinear filtering, Report 8019, Econometric Inst., Erasmus Univ. Rotterdam, 1980. To appear in Stochastics.
[6] Howe, R. E., On the role of the Heisenberg group in harmonic analysis, Aull. Amer. Math. Soc., 3 (1980), 821-844.
[ 8 ] Jazwinsky, A. H., Stochastic processes and filtering theory, Acad. Pr., 1970.
[ 9 ] Kalman, R. E., Contributions to the theory of optimal control, Bol. Soc. Mat. Mexicana 1960, 102-119.
[10] Krishnaprasad, P. S., Marcus, S. I., Identification and tracking: a class of nonlinear filtering probIems, In: Proc. JACC, June 1981.
[11] Mitter, S. K., On the analogy between mathematical problems of nonlinear filtering and quantum
physics，Ricerche di Automatica，10：2（1979），163－216．
［12］Segal，I．E．，Transforms for operators and symplectic automorphisms over a locally compact abelian group，Math．Scand．，13（1963），31－43．
［13］Shale，D．，Linear symmetrics of free boson fields，Trans．Amer．Math．，Soc．， 103 （1962），149－167．
［14］Weil，A．，Sur certains groupes d＇operateurs unitaires，In：Collected Papers Vol．III，1－71，Sprin－ ger 1980；Comments on this paper，ibid，443－447．

