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State-dependent importance sampling for a Jackson tandem network
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# State-dependent Importance Sampling for a Jackson Tandem Network* 

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This paper considers importance sampling as a tool for rare-event simulation. The focus is on estimating the probability of overflow in the downstream queue of a Jacksonian two-node tandem queue - it is known that in this setting 'traditional' stateindependent importance-sampling distributions perform poorly. We therefore concentrate on developing a state-dependent change of measure, that we prove to be asymptotically efficient.

More specific contributions are the following. (i) We concentrate on the probability of the second queue exceeding a certain predefined threshold before the system empties. Importantly, we identify an asymptotically efficient importance-sampling distribution for any initial state of the system. (ii) The choice of the importance-sampling distribution is backed up by appealing heuristics that are rooted in large-deviations theory. (iii) Our method for proving asymptotic efficiency is substantially more straightforward than some that have been used earlier. The paper is concluded by simulation experiments that show a considerable speed up.


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## 1 Introduction

Rare event analysis of queueing networks has been attracting continuous and growing attention over the past decades. As explicit expressions are hardly available, one usually relies on asymptotic techniques to approximate small overflow probabilities. These asymptotics, however, often lack error bounds, and consequently it is not always clear whether their use is justified for given parameters. This explains why one often opts for simulation methods instead.
The use of simulation for estimating rare event probabilities has an inherent problem: the event under consideration occurs so rarely during the simulation, that it is extremely time consuming to obtain a reliable estimate; a rule of thumb is that the number of occurrences needed to obtain an estimate of a certain predefined accuracy is inversely proportional to the probability of interest. Perhaps the most prominent remedy to this problem is importance sampling (IS), i.e., simulating the system under a new probability measure, and correcting the simulation output by means of likelihood ratios (which essentially capture the likelihood of the realization under the old measure with respect to the new measure) to retain unbiasedness. Evidently, it makes sense to choose an IS distribution which guarantees frequent occurrence of the event of interest. The choice of a 'good' new measure is rather delicate though. It should be chosen such that the above-mentioned likelihood ratio tends to be small on the event of interest; choosing a 'wrong' new measure, one may even end up with an estimator with infinite variance. We refer to, e.g., Heidelberger [9] for more background on IS and its pitfalls.
'Classical' papers on the use of IS in queueing usually rely on a so-called 'state-independent' change of measure, i.e., for any state in the system the probabilistic law is changed in the same manner. Usually, large deviations techniques are used to motivate the choice of the new measure, and to prove that the resulting estimator has specific desirable properties (such as bounds on the likelihood on the event of interest). In this respect we mention the seminal paper by Parekh and Walrand [14], that focuses on the estimation of the probability of overflow in a single queue, but also on the probability of the total queue population in a network reaching some threshold. The new measure then corresponds to an unstable queueing system; for instance in the case of a single $M / M / 1$ system the arrival and service rates should be swapped. A fundamental treatment of this change of measure, in fact even for the multi-server queue GI/GI/ $m$ (where it is tacitly assumed that the service times are light-tailed), was given by Sadowsky [15]. His main result is that the corresponding estimator is asymptotically efficient (or: asymptotically optimal), which effectively means that the variance of the estimator behaves roughly like the square of its first moment. In a setting in which the overflow probability decays exponentially in the buffer size $B$, asymptotic efficiency means that the number of replications needed to obtain an estimator with fixed relative error grows subexponentially fast with the 'rarity parameter' $B$.

Things complicate tremendously when looking at networks rather than one-node systems. For the Jacksonian two-node tandem queue (that is, Poisson arrivals, exponential service times at both queues), aiming at estimating the probability that the network population
exceeds a given threshold, [14] proposed to swap the arrival rate with the rate of the slowest server - this makes, heuristically, sense, as the slowest server corresponds to the bottleneck queue. In this case experimental results were not so encouraging as in the case of a single queue, and the quality of the simulation results was strongly affected by the specific values of the arrival and service rates. Later it was proved that this method is asymptotically efficient for some parameter values, but has unbounded variance for other values, see [8] and [2]. In fact, it was proven that no state-independent change of measure exists that is asymptotically efficient for all parameter values.
It was realized that the main problem of state-independent IS schemes is that the transition rates are changed in a 'uniform manner', i.e., irrespective of whether one of the queues is empty or not. As a result it cannot be guaranteed that the likelihood ratio is bounded on the event of interest, and therefore the IS scheme proposed in [14] performs poorly for some parameter values. Some of the first attempts to solve this problem can be found in [3] and [11], in which state-dependent IS schemes were proposed, i.e., IS distributions that are not uniform over the state space. Dupuis et al. [7] were the first to prove asymptotic efficiency for a state-dependent IS scheme for estimating overflow probabilities in a $d$-node Jackson network.

Several important questions are, however, still open; let us from now on concentrate on the two-node Jackson tandem network. In the first place, the majority of papers on this type of networks deals with the probability that, starting in a situation with both queues empty, the total network population exceeds a certain threshold. One may wonder, though, what the impact of the starting state is on the IS scheme. Also, it is not a priori clear how to change the simulation procedures if one is interested in the event of overflow in a specific queue (rather than the total queue).
The main topic of the present paper concerns the development of an asymptotically efficient IS algorithm for estimating the probability that the content of the downstream queue exceeds a certain threshold $B$ before the system becomes empty, starting in any initial state, say $x \in \mathbb{N} \times\{0, \ldots, B-1\}$.
The search for an appropriate change of measure greatly benefits from powerful largedeviations based heuristics. We express the decay rate of the probability of our interest in terms of so-called 'cost functions', that assign cost to paths; the 'most likely path' is then defined as the 'cheapest' path from state $x$ to the 'overflow set' $\mathbb{N} \times\{B, B+1, \ldots\}$ (that does not visit the origin). The intuition is that, conditional on the event that the second queue indeed reaches $B$ before the system gets empty, the trajectory of the Markov process will be typically close to this most likely path. Then the idea is that knowledge of the most likely path helps in finding a good change of measure. The shape of the most likely path strongly depends on which of the two queues is the bottleneck (i.e., has the lowest service rate). When it comes to proving asymptotic efficiency, the two cases have to be dealt with differently. We remark that the most likely path can have a rather unexpected shape; there are situations that, starting in a state $x$ in which the second queue is nonempty, this path is such that first the second queue becomes empty while the first queue fills (to end up in some state $(y, 0)$ ), and then the first queue drains while the second queue
builds up. Another interesting observation is that the most likely path is not continuous in the starting state $x$ : two nearly identical initial states can reach the 'overflow set' in an entirely different manner. We also mention that a non-trivial technical issue we deal with is the infinite state space, in that the process can attain any value in $\mathbb{N} \times\{0, \ldots, B-1\}$, cf. [11]; this complication does not play a role when analyzing rare-event probabilities related to the total network population.
We expect that the above-mentioned large-deviations heuristic can be rather helpful when analyzing a broad class of networks; see also earlier results in [13] for the model that was introduced in [17], in which the service rate of the first queue depends on the content of the second queue.
The proof technique is essentially based on that of Dupuis et al. [7], but, as in De Boer and Scheinhardt [4], we have managed to simplify the proofs considerably. The change of measure is such that the most likely path is, roughly, followed (that is, with high probability), with corrections for the regions near the axes. The proof of asymptotic efficiency then relies on bounding the likelihood on the event of interest.

We end this section by detailing the structure of the paper. Model and preliminaries, as well as a short overview on the basics of IS, are presented in Section 2. In Section 3 we heuristically construct a state-dependent IS scheme for estimating the probability of our interest; interesting corollary results are (i) the most likely path, and (ii) the corresponding decay rate. In Section 4 we present a number of large-deviations properties of the system, where the main result is the correctness of the decay rate that was heuristically derived in the previous section. Section 5 shows that our IS scheme, after a minor adaptation that deals with visits to the axes, is indeed asymptotically efficient. We conclude this paper with supporting numerical results in Section 6, and conclusions in Section 7.

## 2 Model and preliminaries

We consider a two-node tandem Jackson network with Poisson arrivals at rate $\lambda$ to the first station. Each job receives service at the first station, after which it is routed to the second station. After receiving service at the second station, the job leaves the system. Service times at station $i$ have an exponential distribution with parameter $\mu_{i}, i=1,2$. The waiting rooms at both stations are assumed to be infinitely large.
Let $Q(t)=\left\{\left(Q_{1}(t), Q_{2}(t)\right), t \geq 0\right\}$ be the joint queue-length process, as in [7] and [4], from which we will borrow some more notation. Then it is clear that this is a continuoustime Markov process, with possible jump directions $v_{0}=(1,0), v_{1}=(-1,1)$ and $v_{2}=$ $(0,-1)$ with corresponding transition rates $\lambda, \mu_{1}$ and $\mu_{2}$ respectively. The process $Q(t)$ is regenerative if we impose the stability condition $\lambda<\min \left(\mu_{1}, \mu_{2}\right)$, which we will do from now on.
The queue-length process can also be described by the embedded discrete time Markov chain $Q_{j}=\left(Q_{1, j}, Q_{2, j}\right)$, where $Q_{i, j}$ is the number of jobs in queue $i$ after the $j$-th transition. Without loss of generality we will choose the parameters such that $\lambda+\mu_{1}+\mu_{2}=1$, so that they also represent the transition probabilities of $Q_{j}$ in the interior of the state space. To


Figure 1: State space and transition structure for scaled process $X(t)$.
ensure that the same holds on the boundaries, we shall introduce socalled self-transitions shortly, see below.
Our main interest is to estimate the probability that $Q(t)$ (or equivalently, $Q_{j}$ ) reaches some high level $B$ in the second buffer before it returns to the origin, starting from any state. Thereto, it will be convenient to also consider the scaled processes $X(t)=Q(B t) / B$ (in continuous time) and $X_{j}=Q_{j} / B$ (in discrete time). The advantage of these scalings is the invariance of the state space for any $B$. In particular, our target probability is equivalent to the probability that the second component of either the scaled process $X_{j}$ or the scaled process $X(t)$ reaches 1 before the process returns to the origin.
We introduce the following subsets of the state space

$$
\begin{aligned}
D & :=\left\{\left(x_{1}, x_{2}\right): x_{1}>0,0<x_{2}<1\right\}, \\
\partial_{1} & :=\left\{\left(0, x_{2}\right): 0<x_{2}<1\right\}, \\
\partial_{2} & :=\left\{\left(x_{1}, 0\right): x_{1}>0\right\}, \\
\partial_{e} & :=\left\{\left(x_{1}, 1\right): x_{1}>0\right\},
\end{aligned}
$$

and denote the state space by $\bar{D}=D \cup \partial_{e} \cup \partial_{1} \cup \partial_{2}$ (realize that we can exclude $x_{2}>1$ from the state space). Note that transition $v_{k}$ is impossible when queue $k$ is empty, i.e., when $X_{j} \in \partial_{k}$. We modify the process $X_{j}$ to deal with this by allowing some self-transitions in the following way, see also Figure 1:

$$
\begin{equation*}
\mathbb{P}\left(X_{j+1}=X_{j} \mid X_{j} \in \partial_{k}\right)=\mu_{k}, \text { for } k=1,2 . \tag{1}
\end{equation*}
$$

Next, we introduce the stopping time $\tau_{B}^{x}$, which is the first time that the process $X_{j}$ hits level 1 , starting from state $x=\left(x_{1}, x_{2}\right)$, without visits to the origin:

$$
\begin{equation*}
\tau_{B}^{x}=\inf \left\{k>0: X_{k} \in \partial_{e}, X_{j} \neq 0 \text { for } j=1, \ldots, k-1\right\} \tag{2}
\end{equation*}
$$

and we define $\tau_{B}^{x}=\infty$ if $X_{j}$ hits the origin before $\partial_{e}$. It will also be convenient to let $I_{B}\left(A^{x}\right)$ be the indicator of the event $\tau_{B}^{x}<\infty$ for the path $A^{x}=\left(X_{j}, j=0, \ldots: X_{0}=x\right)$.

Thus we can write the probability of our interest as

$$
\begin{equation*}
p_{B}^{x}=\mathbb{E} I_{B}\left(A^{x}\right)=\mathbb{P}\left(\tau_{B}^{x}<\infty\right) . \tag{3}
\end{equation*}
$$

It is clear that it is not efficient to estimate $p_{B}^{x}$ via straightforward simulations when $B$ is large, due to the rarity of the event of interest. In order to reduce the simulation time we will employ Importance Sampling (IS), i.e., we perform simulations under a new measure $\mathbb{Q}$, which replaces the transition rates corresponding to $v_{0}, v_{1}, v_{2}$ by other values. In particular, we will use a state-dependent IS scheme.
This means that the transition rates under the new measure $\mathbb{Q}$ may depend on the current state $x$ of the process; they will be denoted by $\bar{\lambda}(x), \bar{\mu}_{1}(x)$ and $\bar{\mu}_{2}(x)$ respectively.
The probability $p_{B}^{x}$ can now also be expressed as

$$
\begin{equation*}
p_{B}^{x}=\mathbb{E}^{\mathbb{Q}}\left[L\left(A^{x}\right) I_{B}\left(A^{x}\right)\right], \tag{4}
\end{equation*}
$$

where $L\left(A^{x}\right)$ is the likelihood ratio (also known as Radon-Nikodym derivative) of the path $A^{x}$. It is given by

$$
\begin{equation*}
L\left(A^{x}\right)=\prod_{j=0}^{\tau_{B}^{x}-1} \frac{\mathbb{P}\left(Y_{j}\right)}{\mathbb{Q}\left(Y_{j} \mid X_{j}\right)} \tag{5}
\end{equation*}
$$

where $Y_{j}=B\left(X_{j+1}-X_{j}\right)$, unless $X_{j+1}=X_{j}$ in which case $Y_{j}=v_{k}$, if $X_{j} \in \partial_{k}$. Furthermore, $\mathbb{P}\left(Y_{j}\right)$ is the stochastic kernel of the scaled process $X_{j}$ under the old measure, being equal to $\lambda, \mu_{1}$ or $\mu_{2}$ if $j=0,1,2$, respectively, and $\mathbb{Q}\left(Y_{j} \mid X_{j}\right)$ is the kernel under the new measure, given by $\bar{\lambda}(x), \bar{\mu}_{1}(x)$ or $\bar{\mu}_{2}(x)$ when the current state is $X_{j}=x$.

Definition 2.1. The IS scheme for $p_{B}^{x}$ is called asymptotically efficient if

$$
\begin{equation*}
\liminf _{B \rightarrow \infty} \frac{\log \mathbb{E}^{\mathbb{Q}}\left[L^{2}\left(A^{x}\right) I_{B}\left(A^{x}\right)\right]}{\log \mathbb{E}^{\mathbb{Q}}\left[L\left(A^{x}\right) I_{B}\left(A^{x}\right)\right]} \geq 2 \tag{6}
\end{equation*}
$$

In our case it is known that $p_{B}^{x}$ decays exponentially in $B$, so that the exponential decay rate is well defined, i.e.,

$$
\lim _{B \rightarrow \infty} \frac{1}{B} \log p_{B}^{x} \in(0, \infty)
$$

As a result, (6) can be rewritten in the following form:

$$
\limsup _{B \rightarrow \infty} \frac{1}{B} \log \mathbb{E}\left[L\left(A^{x}\right) I_{B}\left(A^{x}\right)\right] \leq 2 \lim _{B \rightarrow \infty} \frac{1}{B} \log p_{B}^{x} .
$$

## 3 Optimal path and related change of measure

In order to find a good change of measure for IS simulations, the first step is usually to find the most probable path to overflow, i.e., the way in which overflow most probably occurs, conditional on its occurrence. In Section 3.1 we explain a method in which minimizing certain 'cost-functions' leads to the most probable path and a good corresponding change
of measure, given by new (state-dependent) transition rates $\tilde{\lambda}(x), \tilde{\mu}_{1}(x)$ and $\tilde{\mu}_{2}(x)$. Then, we split the problem, since the minimization procedure gives different results in different cases. In Section 3.2 we treat the case $\lambda<\mu_{2}<\mu_{1}$, in which the second server is the bottleneck, while Section 3.3 deals with the case $\lambda<\mu_{1} \leq \mu_{2}$, in which the first server is the bottleneck. Beforehand, we would like to point out that the change of measure mentioned above, denoted by tildes, is not the same as the asymptotically efficient change of measure that will be introduced in Section 5 (denoted by bars), although it is closely related.

### 3.1 Cost and structure of path to overflow

The typical path to overflow in the particular case that the origin is the starting point, has already been identified for the $d$-node Jackson tandem network in [1], and hence also for our tandem system. In that paper, the time-reversed process is used to find the shape of the most probable path to overflow. This path to overflow was also obtained as a corollary result in [13], and in this section we present a method similar to the one in [13] to find the optimal path starting from any state $x \in \bar{D}$. The advantage of this method is that it also provides a 'good' change of measure, which ensures that most simulation runs under this new measure will be close to the optimal path. This new measure will be the basis for another change of measure, which is used in our (state-dependent) IS scheme, as presented in Section 5. Another result of our method is the exponential decay rate of $p_{B}^{x}$, which will be determined in Section 4, and which will play a crucial role in the proofs of asymptotic efficiency of Section 5.
Before introducing our method we impose some restrictions on the path structures we consider, leaving the proof that the typical path to overflow indeed satsifies these restrictions to Section 4, see Lemma 4.4. We will only consider the following paths.

## Property 3.1.

- Each path is a concatenation of subpaths, which are straight lines on any of the subsets $D, \delta_{1}$ and $\delta_{2}$, and the new measure stays constant along each subpath, i.e., $\tilde{\lambda}(x)=\tilde{\lambda}, \tilde{\mu}_{1}(x)=\tilde{\mu}_{1}$ and $\tilde{\mu}_{2}(x)=\tilde{\mu}_{2}$, for any state $x$ on the same subpath;
- Each path does not have more than one subpath in each subset if $\mu_{2}<\mu_{1}$;
- Each path does not have more than two subpaths in each subset if $\mu_{2} \geq \mu_{1}$.

With every path that satisfies Property 3.1 we associate a 'cost', the main idea (to be proved in Section 4) being that the minimal cost of the path to overflow in the second buffer, starting from state $x$, can be interpreted as the decay rate of the probability of interest. Our method is based on the family of cost functions $I$, defined by

$$
\begin{equation*}
I(\tilde{\lambda} \mid \lambda):=\lambda-\tilde{\lambda}+\tilde{\lambda} \log \frac{\tilde{\lambda}}{\lambda} \tag{7}
\end{equation*}
$$

see also [16, pages 14 and 20]. Note that the function (7) is convex and equals 0 at $\tilde{\lambda}=\lambda$. Intuitively, we can think of the value $I(\tilde{\lambda} \mid \lambda)$ as the cost we need to pay to let a Poisson process with parameter $\lambda$ behave like a Poisson process with parameter $\lambda$, per time unit.

We will now explain our cost method in more detail in the following two examples. More background can be found in the Appendix of [13].

Example 3.2. As an example, consider a straight path through the interior of the state space, staying away from the boundaries, from some state $x$ to another state $y$, where $x_{1} \geq y_{1}$ and $x_{2}<y_{2}$. We then need to construct a new measure $\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)$, such that $\tilde{\mu}_{1}>\tilde{\mu}_{2}$ and $\tilde{\lambda} \leq \tilde{\mu}_{1}$. This measure ensures that our path has constant north-west drift, or in other words, due to the scaling, our path has a constant slope

$$
\begin{equation*}
\alpha=\frac{\tilde{\mu}_{1}-\tilde{\mu}_{2}}{\tilde{\lambda}-\tilde{\mu}_{1}} \tag{8}
\end{equation*}
$$

The total cost of such a path, per unit time is

$$
\begin{equation*}
\mathbb{I}\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right):=I(\tilde{\lambda} \mid \lambda)+I\left(\tilde{\mu}_{1} \mid \mu_{1}\right)+I\left(\tilde{\mu}_{2} \mid \mu_{2}\right) \tag{9}
\end{equation*}
$$

To find the cost per unit horizontal (vertical) distance, we need to divide this by the horizontal speed $\tilde{\lambda}-\tilde{\mu}_{1}$ (vertical speed $\tilde{\mu}_{1}-\tilde{\mu}_{2}$ ). Thus, minimizing the cost of any straight path from $x$ to $y$ in this case boils down to minimizing

$$
\begin{equation*}
\left(y_{2}-x_{2}\right) \frac{\mathbb{I}\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)}{\tilde{\mu}_{1}-\tilde{\mu}_{2}} \tag{10}
\end{equation*}
$$

over $\tilde{\mu}_{1}$ and $\tilde{\mu}_{2}$, such that $\tilde{\lambda} \leq \tilde{\mu}_{1}$ and $\tilde{\mu}_{1}>\tilde{\mu}_{2}$ hold, as well as

$$
\tilde{\lambda}=\tilde{\mu}_{1}+\frac{y_{1}-x_{1}}{y_{2}-x_{2}}\left(\tilde{\mu}_{1}-\tilde{\mu}_{2}\right)
$$

in addition, we should have that

$$
\frac{y_{2}-x_{2}}{\tilde{\mu}_{1}-\tilde{\mu}_{2}}=\frac{y_{1}-x_{1}}{\tilde{\lambda}-\tilde{\mu}_{1}}
$$

to guarantee that $y$ is indeed the ending state of the path when it starts at $x$.
It is easily checked that the total cost (10) with ending state $y=(0,1)$ attains its minimum when triplet $\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)$ is a solution to

$$
\left\{\begin{array}{l}
\tilde{\lambda}=\tilde{\mu}_{1}-\frac{x_{1}}{1-x_{2}}\left(\tilde{\mu}_{1}-\tilde{\mu}_{2}\right)  \tag{11}\\
\tilde{\lambda}+\tilde{\mu}_{1}+\tilde{\mu}_{2}=\lambda+\mu_{1}+\mu_{2} \\
\tilde{\lambda} \tilde{\mu}_{1} \tilde{\mu}_{2}=\lambda \mu_{1} \mu_{2} \\
\tilde{\lambda} \leq \tilde{\mu}_{1} \text { and } \tilde{\mu}_{1}>\tilde{\mu}_{2} \\
\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}>0
\end{array}\right.
$$

The reason we have chosen the specific ending state $(0,1)$ is that it is the most frequent ending state for our network. Notice also that if $\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)$ is the solution to (11) for some starting state $x$, it also minimizes this system if we replace $x$ by any state that belongs to the straight line between $x$ and $y=(0,1)$.

Example 3.3. Let us now give an example for another type of path with starting state $x \in D$ and ending state $(0,1)$, consisting of two (straight) subpaths. The first subpath belongs to the interior and has north-west drift. The second part belongs to the vertical boundary and has north drift. Thus, it may be denoted as $\left(x_{1}, x_{2}\right) \rightarrow\left(0, x_{2}+\alpha^{-1} x_{1}\right) \rightarrow(0,1)$, for same slope $\alpha$. Property 3.1 tells us that the new measure stays constant along each subpath, so the total cost of such a path is

$$
\alpha^{-1} x_{1} \frac{\mathbb{I}\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)}{\tilde{\mu}_{1}-\tilde{\mu}_{2}}+\left(1-x_{2}-\alpha^{-1} x_{1}\right) \frac{\mathbb{I}\left(\hat{\lambda}, \hat{\mu}_{1}, \hat{\mu}_{2}\right)}{\hat{\lambda}-\hat{\mu}_{2}},
$$

where $\alpha=\left(\tilde{\mu}_{1}-\tilde{\mu}_{2}\right) /\left(\tilde{\lambda}-\tilde{\mu}_{1}\right)$, see (8). The first term in the sum is the cost of the first subpath under some new measure ( $\left.\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)$ and the second term is the cost of the second (vertical) subpath under some measure ( $\hat{\lambda}, \hat{\mu}_{1}, \hat{\mu}_{2}$ ). Optimizing this expression such that $\tilde{\lambda}<\tilde{\mu}_{1}, \tilde{\mu}_{2}<\tilde{\mu}_{1}, \hat{\lambda} \leq \hat{\mu}_{1}$ and $\hat{\mu}_{2}<\hat{\mu}_{1}$, for the case $\mu_{2}<\mu_{1}$, over all parameters marked with tildes and hats, it is readily verified that the minimal cost of this path type is obtained when the new measure is given by

$$
\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)=\left(\hat{\lambda}, \hat{\mu}_{1}, \hat{\mu}_{2}\right)=\left(\mu_{2}, \mu_{1}, \lambda\right),
$$

i.e., by simply interchanging the arrival rate $\lambda$ and the service rate of the second station $\mu_{2}$ for both subpaths.

By considering all possible path types we obtain the overal minimum cost, corresponding to the most probable path, and the corresponding (state-dependent) change of measure $\tilde{\lambda}$, $\tilde{\mu}_{1}$ and $\tilde{\mu}_{2}$. Finally, we also have

$$
\gamma_{x}:=\text { minimal cost over all paths } x \rightarrow \delta_{e},
$$

at our disposal. In Theorem 4.1 we will prove that this is in fact the exponential decay rate of the probability $p_{B}^{x}$ as $B \rightarrow \infty$.
We now present the results of our minimum-cost-path method for both cases of the tandem network.

### 3.2 Optimal path results for $\lambda<\mu_{2}<\mu_{1}$

When $\mu_{2}<\mu_{1}$, the cost minimization starting in state $x$ as outlined in the previous section (in particular Example 3.3; see also the Appendix in [13] for more examples), yields the following new measure after some calculations:

$$
\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)= \begin{cases}\left(\mu_{2}, \mu_{1}, \lambda\right), & \text { if } x \in A_{1},  \tag{12}\\ \text { solution to }(11), & \text { if } x \in A_{2}, \\ \left(\lambda, \mu_{1}, \mu_{2}\right), & \text { if } x \in A_{3}\end{cases}
$$

Here $A_{i}, i=1,2,3$, is the following partition of the state space $\bar{D}$, see also Figure 2:

$$
\begin{align*}
& A_{1}:=\left\{x \in \bar{D}: x_{2} \leq-x_{1} / \alpha_{1}+1\right\}, \\
& A_{2}:=\left\{x \in \bar{D}:-x_{1} / \alpha_{1}+1<x_{2}<-\alpha_{1} x_{1}+1\right\},  \tag{13}\\
& A_{3}:=\left\{x \in \bar{D}: x_{2} \geq-\alpha_{1} x_{1}+1\right\},
\end{align*}
$$



Figure 2: Partition of $\bar{D}$ and some optimal paths to overflow when $\mu_{2}<\mu_{1}$.
with $\alpha_{1}:=\left(\mu_{1}-\mu_{2}\right) /\left(\mu_{1}-\lambda\right)$. Note that the path considered in Example 3.3 in the previous subsection is optimal for any starting state $x \in A_{1}$, and the corresponding new measure (exchanging $\lambda$ and $\mu_{2}$ ) was earlier found by Parekh and Walrand [14] for the problem of reaching a large total queue population. Also, we point out that the change of measure is continuous in the state $x$, as can be verified by solving system (11) for $x=\left(\alpha_{1}, 0\right)$ and $x=\left(\alpha_{1}^{-1}, 0\right)$, yielding the solutions in the first and third lines of (12), respectively.
The corresponding path from starting state $x=\left(x_{1}, x_{2}\right)$ to some state on $\partial_{e}$ is given by

$$
\begin{array}{ll}
\left(x_{1}, x_{2}\right) \rightarrow\left(0, x_{2}+\alpha_{1}^{-1} x_{1}\right) \rightarrow(0,1), & \text { if } x \in A_{1}, \\
\left(x_{1}, x_{2}\right) \rightarrow(0,1), & \text { if } x \in A_{2},  \tag{14}\\
\left(x_{1}, x_{2}\right) \rightarrow\left(x_{1}-\alpha_{1}^{-1} x_{2}, 1\right), & \text { if } x \in A_{3} .
\end{array}
$$

The resulting cost $\gamma_{x}$ of the optimal path is given by:

$$
\gamma_{x}= \begin{cases}\left(1-x_{1}-x_{2}\right) \gamma, & \text { if } x \in A_{1},  \tag{15}\\ -x_{1} \log \frac{\tilde{\lambda}(x)}{\lambda}-\left(1-x_{2}\right) \log \frac{\tilde{\mu}_{2}(x)}{\mu_{2}}, & \text { if } x \in A_{2}, \\ 0, & \text { if } x \in A_{3},\end{cases}
$$

where

$$
\gamma:=-\log \frac{\lambda}{\mu_{2}},
$$

is the minimal cost of the path $(0,0) \rightarrow(0,1)$.
We end this subsection with some interesting properties of the new measure defined in (12), to be used later. Intuitively, it says that for any state $x$, the new measure 'lies between' the Parekh and Walrand measure where $\lambda$ and $\mu_{2}$ are interchanged, and the 'normal' measure, where the parameters retain their original values. Moreover, the more jobs are present in the system at time zero, either in queue 1 or in queue 2 , the 'less change of measure' we need. This perfectly coincides with the structure of the most probable path, see (14).

Lemma 3.4. When $\mu_{2}<\mu_{1}$, the functions $\tilde{\lambda}(x), \tilde{\mu}_{1}(x)$ and $\tilde{\mu}_{2}(x)$ as defined in (12) are continuous and differentiable, satisfying the following for any $x \in \bar{D}$.
(i) $\frac{\partial \tilde{\lambda}(x)}{\partial x_{1}} \leq 0, \frac{\partial \tilde{\mu}_{2}(x)}{\partial x_{1}} \geq 0, \frac{\partial \tilde{\lambda}(x)}{\partial x_{2}} \leq 0 \quad$ and $\quad \frac{\partial \tilde{\mu}_{2}(x)}{\partial x_{2}} \geq 0$.
(ii) $\tilde{\lambda}(x) \in\left[\lambda, \mu_{2}\right]$ and $\tilde{\mu}_{2}(x) \in\left[\lambda, \mu_{2}\right]$.
(iii) $\gamma=\max _{x \in \bar{D}} \gamma_{x}$.

Proof. (i) We only need to consider $x \in A_{2}$, since otherwise all partial derivatives are zero. Applying implicit differentiation to (11) one finds

$$
\frac{\partial \tilde{\lambda}(x)}{\partial x_{1}}=-\frac{\left(1-x_{2}\right)\left(\tilde{\mu}_{1}(x)-\tilde{\mu}_{2}(x)\right) \tilde{\lambda}(x)}{\left(1-x_{2}\right)^{2} \tilde{\lambda}(x)+\left(1-x_{1}-x_{2}\right)^{2} \tilde{\mu}_{1}(x)+x_{1}^{2} \tilde{\mu}_{2}(x)} \leq 0,
$$

where the last inequality follows from the fact that $\tilde{\mu}_{1}(x)>\tilde{\mu}_{2}(x)$. The other statements follow similarly.
(ii) It follows from (12) that $\tilde{\lambda}(x)=\mu_{2}$ if $x \in \partial_{1}$ and $\tilde{\lambda}(x)=\lambda$ if $x \in A_{3}$, so applying the first statement of this lemma one can find that $\tilde{\lambda}(x) \in\left[\lambda, \mu_{2}\right]$. Using similar arguments one can obtain the same bounds for $\tilde{\mu}_{2}(x)$.
(iii) We show that the partial derivatives with respect to $x_{1}$ and $x_{2}$ of $\gamma_{x}$ as given in (15) are not positive. For $x \in A_{1} \cup A_{3}$ this is obvious, while for $x \in A_{2}$ it can be checked using implicit differentiation, similar to the proof of the first statement.

Lemma 3.4 does not yield results on $\tilde{\mu}_{1}(x)$ since they are not needed in the sequel, but it may be interesting to note that $\tilde{\mu}_{1}(x)$ is not monotone. In fact, $\tilde{\mu}_{1}(x)=\mu$ when $x \in A_{1}$ and when $x \in A_{3}$, but also when $x_{1}+x_{2}=1$, so it is also neither convex nor concave.

### 3.3 Optimal path results for $\lambda<\mu_{1} \leq \mu_{2}$

The new measure under which the path to overflow has minimal cost in terms of (7) is as follows:

$$
\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)= \begin{cases}\left(\mu_{1}, \lambda, \mu_{2}\right), & \text { if } x \in B_{1},  \tag{16}\\ \text { solution to }(11), & \text { if } x \in B_{2}, \\ \left(\lambda, \mu_{2}, \mu_{1}\right), & \text { if } x \in B_{3} .\end{cases}
$$

Again we partitioned the state space into three subspaces $B_{i}, i=1,2,3$ as follows, see also Figure 3.

$$
\begin{align*}
& B_{1}:=\{x \in \bar{D}: f(x) \leq 0\}, \\
& B_{2}:=\left\{x \in \bar{D}: f(x)>0 \text { and } x_{2}<-\alpha_{2} x_{1}+1\right\},  \tag{17}\\
& B_{3}:=\left\{x \in \bar{D}: x_{2} \geq-\alpha_{2} x_{1}+1\right\},
\end{align*}
$$

where $\alpha_{2}:=\left(\mu_{2}-\mu_{1}\right) /\left(\mu_{2}-\lambda\right)$ and

$$
f(x):=\gamma+x_{1} \log \frac{\tilde{\lambda}(x)}{\mu_{1}}+\left(1-x_{2}\right) \log \frac{\tilde{\mu}_{2}(x)}{\mu_{2}}
$$



Figure 3: Partition of $\bar{D}$ and some optimal paths to overflow when $\mu_{1} \leq \mu_{2}$.
with $\tilde{\lambda} \equiv \tilde{\lambda}(x)$ and $\tilde{\mu}_{2} \equiv \tilde{\mu}_{2}(x)$ being the solution to (11). The zero level curve of the function $f(x)$ represents the boundary between subspaces $B_{1}$ and $B_{2}, \beta$ is the unique solution to $f\left(0, x_{2}\right)=0$. Interestingly, for the current case the change of measure is not continuous in states $x$ that lie on this boundary (i.e., $f(x)=0$ ), and the behavior on $B_{1}$ and $B_{2}$ is entirely different. In particular, the change of measure on $B_{2}$ has $\tilde{\lambda}(x)<\tilde{\mu}_{1}(x)$ and $\tilde{\mu}_{2}(x)<\tilde{\mu}_{1}(x)$, as opposed to the first line of (16) where both inequalities are reversed. This is also reflected in a different shape of the typical path from $x=\left(x_{1}, x_{2}\right)$ to $\partial_{e}$ :

$$
\begin{array}{ll}
\left(x_{1}, x_{2}\right) \rightarrow\left(x_{1}+\alpha_{3} x_{2}, 0\right) \rightarrow\left(\alpha_{2}, 0\right) \rightarrow(0,1), & \text { if } x \in B_{1}, \\
\left(x_{1}, x_{2}\right) \rightarrow(0,1), & \text { if } x \in B_{2},  \tag{18}\\
\left(x_{1}, x_{2}\right) \rightarrow\left(x_{1}-\alpha_{2}^{-1} x_{2}, 1\right), & \text { if } x \in B_{3},
\end{array}
$$

where $\alpha_{3}:=\left(\mu_{2}-\lambda\right) /\left(\mu_{1}-\lambda\right)$. Note that the last part of any path with starting state $x \in B_{1}$ is just a special case of a path starting in $B_{2}$ (in this case starting in $\left(\alpha_{2}, 0\right)$ ), but the corresponding new measure on this line (i.e. the solution to system (11) for $x=\left(\alpha_{2}, 0\right)$ ) can be given explicitly as ( $\mu_{1}, \mu_{2}, \lambda$ ). This was already known from [13] for the path starting in the origin.
The next result we give is $\gamma_{x}$, the cost of the optimal path in terms of (7):

$$
\gamma_{x}= \begin{cases}\gamma-x_{1} \log \frac{\mu_{1}}{\lambda}, & \text { if } x \in B_{1}  \tag{19}\\ -x_{1} \log \frac{\tilde{\lambda}(x)}{\lambda}-\left(1-x_{2}\right) \log \frac{\tilde{\mu}_{2}(x)}{\mu_{2}}, & \text { if } x \in B_{2} \\ \left(1-x_{2}\right) \log \frac{\mu_{2}}{\mu_{1}}, & \text { if } x \in B_{3}\end{cases}
$$

We end this subsection with the analogue of Lemma 3.4 in the case when $\mu_{1} \leq \mu_{2}$. For this, we first introduce $z$ as the unique solution in the interval $(0,1)$ of the (essentially cubic) equation

$$
\begin{equation*}
\varphi(z):=\lambda+\mu_{1}+\mu_{2}(1-z)-\sqrt{\frac{\lambda \mu_{1}}{z}}=0 \tag{20}
\end{equation*}
$$

which follows from system (11) by taking $\left(x_{1}, x_{2}\right)=(0,0)$. (The fact that there is a unique solution immediately follows from $\varphi(0)=-\infty, \varphi(1)=\lambda+\mu_{1}-\sqrt{\lambda \mu_{1}}>0$, and the fact that $\varphi^{\prime}(z)=0$ has just a single positive solution, viz. $\sqrt[3]{\lambda \mu_{1} / 4 \mu_{2}^{2}}$.) In fact, $-\log z$ is the cost of the vertical path $(0,0) \rightarrow(0,1)$ in the interior (i.e., in $D$ ), satisfying $\tilde{\lambda}=\tilde{\mu}_{1}$ (as opposed to the vertical path following $\partial_{1}$ in Example 3.3, where $\tilde{\lambda}<\tilde{\mu}_{1}$ ). See also [12, Eqns. (30) and (33)] and [13] for more details.

Lemma 3.5. When $\mu_{1} \leq \mu_{2}$, the functions $\tilde{\lambda}(x), \tilde{\mu}_{1}(x)$ and $\tilde{\mu}_{2}(x)$ as defined in (16) are continuous and differentiable, except in states $x$ with $f(x)=0$. For any such $x \in \bar{D}$, these functions satisfy the following.
(i) $\frac{\partial \tilde{\lambda}(x)}{\partial x_{1}} \leq 0, \frac{\partial \tilde{\mu}_{2}(x)}{\partial x_{1}} \geq 0, \frac{\partial \tilde{\lambda}(x)}{\partial x_{2}} \leq 0, \quad$ and $\quad \frac{\partial \tilde{\mu}_{2}(x)}{\partial x_{2}} \geq 0, \quad$.
(ii) $\tilde{\lambda}(x) \in\left[\lambda, \sqrt{\lambda \mu_{1} / z}\right]$ and $\tilde{\mu}_{2}(x) \in\left[\mu_{2} z, \mu_{1}\right] \cup \mu_{2}$, where $z$ is defined by (20).
(iii) $\gamma=\max _{x \in \bar{D}} \gamma_{x}$.

Proof. To prove these facts we can use similar arguments as in the proof of Lemma 3.4. An exception is the second statement, where the upper bound for $\tilde{\lambda}(x)$ and the lower bound for $\tilde{\mu}_{2}(x)$ is attained when $x \in \partial_{1}$ (see the first statement), i.e., $\tilde{\lambda}(x) \leq \sqrt{\lambda \mu_{1} / z}$ and $\tilde{\mu}_{2}(x) \geq \mu_{2} z$, where $z$ is the unique solution of (20) in the interval $(0,1)$.

Note that part (i) of Lemma 3.5 implies that the functions are monotone on $B_{2} \cup B_{3}$, but not on all $\bar{D}$, due to the discontinuity on the boundary between $B_{1}$ and $B_{2}$.

## 4 Large deviations properties

The goal of this section is to formally prove that the cost of the optimal path to overflow is equal to the exponential decay rate of $p_{B}^{x}$, the probability of interest. We also illuminate some important and interesting large deviations properties of the process $X(t)$.
Consider any absolutely continuous function $\phi:[0, \infty) \rightarrow \bar{D}$, representing a path associated with the scaled process $X(t)$. Our first aim is to define a so-called local rate function $\ell(\phi(t), \dot{\phi}(t))$, which depends both on the position at time $t$ and on the time derivative (or speed vector) $\dot{\phi}(t)$ at time $t$. To do so, we first define three auxiliary functions $L_{i}(y)$, where the argument $y$ should be interpreted as a speed vector:

$$
\begin{equation*}
L_{i}(y):=\sup _{\theta}\left(\langle\theta, y\rangle-g_{i}(\theta)\right), \quad i=0,1,2, \tag{21}
\end{equation*}
$$

where

$$
\begin{aligned}
g_{0}(\theta) & :=\lambda\left(e^{\theta_{1}}-1\right)+\mu_{1}\left(e^{\theta_{2}-\theta_{1}}-1\right)+\mu_{2}\left(e^{-\theta_{2}}-1\right) \\
g_{1}(\theta) & :=\lambda\left(e^{\theta_{1}}-1\right)+\mu_{2}\left(e^{-\theta_{2}}-1\right) \\
g_{2}(\theta) & :=\lambda\left(e^{\theta_{1}}-1\right)+\mu_{1}\left(e^{\theta_{2}-\theta_{1}}-1\right)
\end{aligned}
$$

cf. [16, Eqn. (5.5)]. The second equality applies to $\partial_{1}$ and the third equality applies to $\partial_{2}$. The function $g_{1}(\theta)$ does not have a term with $\mu_{1}$, because jumps of type $v_{1}$ from boundary
$\partial_{1}$ are impossible, and likewise $g_{2}(\theta)$ does not have a term with $\mu_{2}$. Finally let us define the local rate function $\ell$ as:

$$
\ell(\phi(t), \dot{\phi}(t)):= \begin{cases}L_{0}(\dot{\phi}(t)), & \text { if } \phi(t) \in D \cup \partial_{e},  \tag{22}\\ {\left[L_{0} \oplus L_{1}\right](\dot{\phi}(t)),} & \text { if } \phi(t) \in \partial_{1}, \\ {\left[L_{0} \oplus L_{2}\right](\dot{\phi}(t)),} & \text { if } \phi(t) \in \partial_{2},\end{cases}
$$

where

$$
\left[L_{0} \oplus L_{i}\right](y):=\inf \left\{\rho L_{0}\left(y_{0}\right)+(1-\rho) L_{i}\left(y_{i}\right): 0 \leq \rho \leq 1, \rho y_{0}+(1-\rho) y_{i}=y\right\}
$$

for $i=1,2$, is the inf-convolution of the functions $L_{0}$ and $L_{i}$, the infimum being taken over all values $\rho$ and vectors $y_{0}$ and $y_{i}$ that satisfy the given conditions. Let us briefly explain why we use this inf-convolution on the boundaries of the state space. Assume that the scaled process $X(t)$ follows a path $\phi(t) \in \partial_{1}$, such that $\partial \phi_{2} / \partial t>0$ for $t \in[0, T]$. Hence, the first and second component of the vector $y$ should be zero and strictly positive, respectively. It is clear that the original (unscaled) jump process $Q(t)$ can only increase its second component when it is not on $\partial_{1}$, since jumps of type $v_{1}$ are not allowed on $\partial_{1}$. Therefore, the inf-convolution provides a 'mixture' of the functions $L_{0}$ and $L_{1}$, supposing that the process $Q(t)$ spends a fraction of time $\rho$ in the interior $D$ and a fraction $1-\rho$ on the vertical constraint. Note that $\rho$ must be such that $\phi(t)$ has speed $y$ with positive increment in the vertical direction and zero-increment in the horizontal direction, such that the scaled process $X(t)$ remains on $\partial_{1}$.
We are now ready to state the following theorem.
Theorem 4.1. The process $X(t)$ satisfies a large deviations principle with local rate function (22), i.e.,

$$
\lim _{B \rightarrow \infty} \frac{1}{B} \log p_{B}^{x}=-\inf \int_{0}^{\tau} \ell(\phi(t), \dot{\phi}(t)) d t
$$

where $\tau:=\inf \left\{t>0: \phi(t) \in \partial_{e}, \phi(s) \neq 0, s \in(0, t)\right\}$ and the infimum is taken over all absolutely continuous functions $\phi:[0 ; \infty) \rightarrow \bar{D}$ such that $\phi(0)=x$ and $\tau<\infty$.

Proof. The proof of this theorem is based on the results presented in [5]. Let us introduce a process $Z(t)$, which is the unconstrained version of $X(t)$, in other words $Z(t)$ is allowed to have negative values in both components. In addition we will assume that $Z(0)=X(0)=$ $x \in \bar{D}$. One can use [5, Thms. 3.2 and 3.4] to show that the map $\Gamma: Z(t) \rightarrow X(t)$ exists and Theorem 2.2 from the same paper to show that it is Lipschitz continuous. $\Gamma$ is known as the Skorokhod map and the question whether it exists is known as the Skorokhod problem; for more background we refer to [5].
Since the map $\Gamma$ is Lipschitz continuous and the process $Z(t)$ satisfies a large deviation principle, see [16, Thm. 5.1], one can apply the contraction principle (see [16, Thm. 2.13]) and conclude that the process of our interest, $X(t)$, satisfies a large deviations principle with local rate function $\ell(\phi(t), \dot{\phi}(t))$ defined by (22).

Using the local rate function $\ell$, as defined in (22), we can define the rate function of any path $\phi(t)=\left(\phi_{1}(t), \phi_{2}(t)\right)$ with $t \in[0, T]$ for some $T$, as the integral of $\ell$ over time. The following lemma shows that for paths that stay in one of the subsets $D, \partial_{1}, \partial_{2}$, this rate function is minimal when the path is straight, with constant speed vector.

Lemma 4.2. For any $T$, consider an absolutely continuous path $\phi(t)$ that remains in $D$ (or in $\partial_{1}$, or in $\partial_{2}$ ) for all $t \in[0, T]$. Then,

$$
\int_{0}^{T} \ell(\phi(t), \dot{\phi}(t)) d t \geq T \ell\left(\phi(0), \frac{\phi(T)-\phi(0)}{T}\right) .
$$

Equality holds only if $\dot{\phi}(t)$ is a constant, i.e., $\phi(t)$ is a straight line.
Proof. The proof of this lemma can be found in [16, p. 87]; we mention that a related result was established in [13, Lemma 4].

Now assume that $\phi(t) \in D$, for $t \in(0, T)$ is a path between two states $x$ and $y$. Lemma 4.2 tells us that the path $\phi(t)$ has minimal cost if the process $X(t)$ moves along a straight line at constant speed. We can define a corresponding new measure as follows

$$
\begin{align*}
\tilde{\lambda} & =\lambda e^{\theta_{1}} \\
\tilde{\mu}_{1} & =\mu_{1} e^{\theta_{2}-\theta_{1}}  \tag{23}\\
\tilde{\mu}_{2} & =\mu_{2} e^{-\theta_{2}},
\end{align*}
$$

where $\theta=\left(\theta_{1}, \theta_{2}\right)$ is the maximizer of (21) with $i=1$. In fact this is exactly the same change of measure we would find using the cost minimization procedure from Section 3, due to the immediate equality

$$
\begin{equation*}
\ell(\phi(t), \dot{\phi}(t))=\mathbb{I}\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right) . \tag{24}
\end{equation*}
$$

This equality however, does not hold on the boundaries. Instead, when $\phi(t)$ stays on $\partial_{1}$ or $\partial_{2}$ for $t \in[0, T]$, we have

$$
\ell(\phi(t), \dot{\phi}(t)) \leq \mathbb{I}\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right),
$$

where the new measure $\left.\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)$ is again defined as in (23). This is not difficult to see since, e.g. for paths on $\partial_{1}$, we find for some $\rho \in[0,1]$ that

$$
\ell(\phi(t), \dot{\phi}(t))=I(\tilde{\lambda} \mid \lambda)+\rho I\left(\tilde{\mu}_{1} \mid \mu_{1}\right)+I\left(\tilde{\mu}_{2} \mid \mu_{2}\right) \leq \mathbb{I}\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right) .
$$

However, we can still show equality between local rate functions and cost functions on the boundaries, but only for the optimal paths. To state this formally, let $\Phi_{1}\left(\Phi_{2}\right)$ be the set of paths that travels a distance $h>0$ along $\partial_{1}\left(\partial_{2}\right)$ at constant speed during a time $\sigma_{1}\left(\sigma_{2}\right)$, i.e.,

$$
\begin{aligned}
& \Phi_{1}=\left\{\phi(t) \subset \partial_{1}: \phi(0)=\left(0, x_{2}^{*}\right), \phi\left(\sigma_{1}\right)=\left(0, x_{2}^{*}+h\right)\right\}, \\
& \Phi_{2}=\left\{\phi(t) \subset \partial_{2}: \phi(0)=\left(x_{1}^{*}, 0\right), \phi\left(\sigma_{2}\right)=\left(x_{1}^{*}+h, 0\right)\right\},
\end{aligned}
$$

for some $x_{1}^{*}$ and $x_{2}^{*}$. Then we have the following relations between the rate function $\ell$ defined by (22) and the cost function $\mathbb{I}$ from the previous section, defined by (9):

Lemma 4.3. (i) For paths in the interior $D$, we have

$$
\ell(\phi(t), \dot{\phi}(t))=\mathbb{I}\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)
$$

(ii) For paths on the vertical boundary $\partial_{1}$, we have

$$
\inf _{\phi \in \Phi_{1}} \int_{0}^{\sigma_{1}} \ell(\phi(t), \dot{\phi}(t)) d t=h \inf \frac{\mathbb{I}\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)}{\tilde{\mu}_{1}-\tilde{\mu}_{2}},
$$

where the second infimum is taken over all $\tilde{\lambda}, \tilde{\mu}_{1}$ and $\tilde{\mu}_{2}$ such that $\tilde{\lambda}<\tilde{\mu}_{1}$ and $\tilde{\mu}_{1}>\tilde{\mu}_{2}$.
(iii) For paths on the horizontal boundary $\partial_{2}$, we have

$$
\inf _{\phi \in \Phi_{2}} \int_{0}^{\sigma_{2}} \ell(\phi(t), \dot{\phi}(t)) d t=h \inf \frac{\mathbb{I}\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)}{\tilde{\lambda}-\tilde{\mu}_{1}}
$$

where the second infimum is taken over all $\tilde{\lambda}, \tilde{\mu}_{1}$ and $\tilde{\mu}_{2}$ such that $\tilde{\lambda}>\tilde{\mu}_{1}$ and $\tilde{\mu}_{2}>\tilde{\mu}_{1}$.
Proof. Statement (i) is the same as (24). We continue to prove statement (iii), the proof of (ii) being identical. The restriction $\phi(t) \subset \partial_{2}$ implies that $\dot{\phi}(t)=\left(\dot{\phi}_{1}(t), 0\right)$ for any $t \in\left[0, \sigma_{2}\right]$. The definition of the inf-convolution tell us that $\tilde{\lambda} v_{0}+\tilde{\mu}_{1} v_{1}+\rho \tilde{\mu}_{2} v_{2}=\dot{\phi}(t)$. Hence we find that $\rho=\tilde{\mu}_{1} / \tilde{\mu}_{2}$ and $\dot{\phi}_{1}(t)=\tilde{\lambda}-\tilde{\mu}_{1}$, from which we can conclude that

$$
\inf \int_{0}^{\sigma_{2}} \ell(\phi(t), \dot{\phi}(t)) d t=h \inf \frac{I(\tilde{\lambda} \mid \lambda)+I\left(\tilde{\mu}_{1} \mid \mu_{1}\right)+\left(\tilde{\mu}_{1} / \tilde{\mu}_{2}\right) I\left(\tilde{\mu}_{2} \mid \mu_{2}\right)}{\tilde{\lambda}-\tilde{\mu}_{1}}
$$

Straightforward minimization shows that the latter equals $h \inf \left\{\mathbb{I}\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right) /\left(\tilde{\lambda}-\tilde{\mu}_{1}\right)\right\}$, because $\tilde{\mu}_{2}=\mu_{2}$ and hence $I\left(\tilde{\mu}_{2} \mid \mu_{2}\right)=0$.

The next lemma validates our choice in the previous section to consider only paths that satisfy Property 3.1.

Lemma 4.4. The optimal path from any starting state $x$ to $\partial_{e}$ does not have more than

- one subpath in each subset, if $\mu_{2}<\mu_{1}$,
- two subpaths in each subset, if $\mu_{2} \geq \mu_{1}$.

Proof. Due to the one-to-one correspondence between the rate function of any path in terms of the local rate function $\ell$ and the cost function $\mathbb{I}$, see Lemma 4.3, the proof of this lemma is similar to the proof of Lemma 5 in [13].

Theorem 4.5. The exponential decay rate of $p_{B}^{x}$ equals the minimal cost derived in Section 3, i.e.,

$$
\lim _{B \rightarrow \infty} \frac{1}{B} \log p_{B}^{x}=-\gamma_{x}
$$

Proof. By virtue of Lemma 4.4, the optimal path can be represented as a concatenation of $k$ (at most 6) subpaths $\phi^{(1)}, \ldots, \phi^{(k)}$, which stay on different subsets of the state space (i.e., $D, \partial_{1}$ and $\partial_{2}$ ). If we denote the starting time and position of the $i$ th subpath by $t^{(i)}$
and $x^{(i)}=\phi^{(i)}\left(t^{(i)}\right)=\phi^{(i-1)}\left(t^{(i)}\right)$ respectively, with the convention that $t^{(1)}=0, x^{(1)}=x$, $t^{(k+1)}=\tau, x^{(k+1)} \in \partial_{e}$, then we can write the decay rate identified in Theorem 4.1 as

$$
\inf _{\phi} \int_{0}^{\tau} \ell(\phi(t), \dot{\phi}(t)) d t=\inf _{t^{(1)}, \ldots, t^{(k)}} \sum_{i=1}^{k} \inf _{\phi^{(i)}} \int_{t^{(i)}}^{t^{(i+1)}} \ell\left(\phi^{(i)}(t), \dot{\phi}^{(i)}(t)\right) d t .
$$

Using Lemma 4.3 we can rewrite the last expression as follows:

$$
\inf _{t^{(1)}, \ldots, t^{(k)}} \sum_{i=1}^{k}\left(x_{1}^{(i+1)}-x_{1}^{(i)}\right) \inf _{\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}} \frac{\mathbb{I}\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)}{\tilde{\lambda}-\tilde{\mu}_{1}},
$$

where in some of the terms we may need to replace the denominator $\tilde{\lambda}-\tilde{\mu}_{1}$ by $\tilde{\mu}_{1}-\tilde{\mu}_{2}$, while also changing the prefactor to $x_{2}^{(i+1)}-x_{2}^{(i)}$, see Lemma 4.3. Using the fact that the new measure ( $\left.\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)$ determines the shape of each subpath $\phi^{(i)}$, and the shape of the whole path $\phi$ as a consequence, we may also take the first infimum over $x^{(1)}, \ldots, x^{(k)}$, rather than $t^{(1)}, \ldots, t^{(k)}$. Applying Lemma 4.3 to the last optimization problem we arrive at

$$
\inf _{x^{(1)}, \ldots, x^{(k)}} \sum_{i=1}^{k}\left(\left[\phi_{i+1}\left(t_{i+1}\right)\right]_{1}-\left[\phi_{i}\left(t_{i}\right)\right]_{1}\right) \inf _{\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}} \frac{\mathbb{I}\left(\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)}{\tilde{\lambda}-\tilde{\mu}_{1}}=\gamma_{x},
$$

which completes the proof.

## 5 Asymptotic efficiency

It is known from [13], where the starting state is the origin, that the new measures (12) and (16) are not always asymptotically efficient. For example, when $\mu_{2}<\mu_{1}$, multiple visits of the process $Q(t)$ to the horizontal axis ( $\partial_{2}$ ) under the new measure ( $\mu_{2}, \mu_{1}, \lambda$ ) may cause the likelihood ratio to become very large. We will 'protect' the likelihood ratio by using a specific measure around $\partial_{2}$, under which these visits become harmless. This approach is similar to the one used in [7]. We will also introduce a protection strip along the lower part of the vertical boundary $\partial_{1}$ in the same manner, in the case when $\mu_{1} \leq \mu_{2}$.
We again split the problem into two cases: in Section 5.1 we explain our method in detail for the situation in which the second server is the bottleneck $\left(\lambda<\mu_{2}<\mu_{1}\right)$, and in Section 5.2 we treat the case in which the first server is the bottleneck $\left(\lambda<\mu_{1} \leq \mu_{2}\right)$.

### 5.1 Asymptotically efficient scheme for $\mu_{2}<\mu_{1}$

In order to construct an IS scheme that is provably asymptotically efficient we introduce a function $W(x)$, defined for any point $x=\left(x_{1}, x_{2}\right)$ of the state space. This function will give us an expression for a new measure ( $\left.\tilde{\lambda}, \tilde{\mu}_{1}, \tilde{\mu}_{2}\right)$ in the same manner as it was done in [7].
Let us first introduce three intermediate functions $W_{i}(x), i=1,2,3$ :

$$
\begin{align*}
& W_{1}(x):=2 \gamma_{x}-\delta, \\
& W_{2}(x):=W_{1}\left(x_{1}, \delta / 2 \gamma\right)=2 \gamma_{\left(x_{1}, \delta / 2 \gamma\right)}-\delta,  \tag{25}\\
& W_{3}(x):=2 \gamma-3 \delta,
\end{align*}
$$

where $\delta$ is some small positive number, and $\gamma_{x}$ is given by (15). In the next step we introduce the function which is the minimum of these three functions, see also Figure 4:

$$
\bar{W}(x):=W_{1}(x) \wedge W_{2}(x) \wedge W_{3}(x) .
$$

Note that our particular choice of the functions $W_{i}$ ensures that the shapes of the areas around the origin and $\partial_{2}$ on which $\bar{W}$ coincides with the functions $W_{i}$ are the same as they were in [7]. The last step in the construction is a mollification procedure which makes the


Figure 4: The function $\bar{W}(x)$ and the areas on which $\bar{W}(x)=W_{i}, i=1,2,3\left(\right.$ case $\left.\mu_{2}<\mu_{1}\right)$. resulting function $W(x)$ smooth. We do this by defining:

$$
\begin{equation*}
W(x):=-\epsilon \log \sum_{i=1}^{3} e^{-W_{i}(x) / \epsilon}, \tag{26}
\end{equation*}
$$

where $\epsilon$ is a 'smoothness' parameter; the larger $\epsilon$ is chosen, the smoother the function $W(x)$ is. On the other hand, as $\epsilon \rightarrow 0$ we see that $W(x)$ converges to the (non-smooth) function $\bar{W}(x)$.
The function $W(x)$, and in particular its gradient, will play a main role in the representation of the state-dependent, asymptotically efficient new measure. However, before turning to this, we need some preliminaries, namely a relation between the gradients of the functions $W_{i}$ and the measure from the previous sections, and some assumptions on the parameters $\delta$ and $\epsilon$.

Proposition 5.1. The gradients of the functions $W_{i}(x), i=1,2,3$ can be represented as follows:

$$
\begin{aligned}
D W_{1}(x) & =2\left(\log \frac{\lambda}{\tilde{\lambda}(x)}, \log \frac{\tilde{\mu}_{2}(x)}{\mu_{2}}\right), \\
D W_{2}(x) & =2\left(\log \frac{\lambda}{\tilde{\lambda}\left(x_{1}, \delta / 2 \gamma\right)}, 0\right), \\
D W_{3}(x) & =(0,0) .
\end{aligned}
$$

Proof. It is clear that $D W_{1}(x)=-2 \gamma(1,1)$ if $x \in A_{1}$. When $x \in A_{2}, D W_{1}(x)$ can be represented in the following form:

$$
\begin{aligned}
D W_{1}(x)= & 2\left(\log \frac{\lambda}{\tilde{\lambda}(x)}, \log \frac{\tilde{\mu}_{2}(x)}{\mu_{2}}\right)-2 x_{1}\left(\frac{\partial \tilde{\lambda}(x) / \partial x_{1}}{\tilde{\lambda}(x)}, \frac{\partial \tilde{\lambda}(x) / \partial x_{2}}{\tilde{\lambda}(x)}\right) \\
& -2\left(1-x_{2}\right)\left(\frac{\partial \tilde{\mu}_{2}(x) / \partial x_{1}}{\tilde{\mu}_{2}(x)}, \frac{\partial \tilde{\mu}_{2}(x) / \partial x_{2}}{\tilde{\mu}_{2}(x)}\right) .
\end{aligned}
$$

Although we do not know $\tilde{\lambda}(x), \tilde{\mu}_{1}(x)$ and $\tilde{\mu}_{2}(x)$ explicitly, we can find their partial derivatives with respect to $x_{1}$ and $x_{2}$ by implicit differentiation of (11), see Lemma 3.4 for more insight. After some elementary algebra we find that the sum of the last two vectors in the previous expression equals zero, which proves the first statement. The other two statements follow easily from the definitions of $W_{2}$ and $W_{3}$.

The parameters $\delta$ and $\epsilon$ depend on $B$, and in the sequel we will need the following conditions for their asymptotical behavior as $B$ grows large. Note that these are the same conditions as in [7] and [4].

Assumption 5.2. The parameters $\delta \equiv \delta_{B}$ and $\epsilon \equiv \epsilon_{B}$ are strictly positive and satisfy the following limit conditions: as $B \rightarrow \infty$, (i) $\epsilon_{B} \rightarrow 0$, (ii) $\delta_{B} \rightarrow 0$, (iii) $B \epsilon_{B} \rightarrow \infty$, (iv) $\varepsilon_{B} / \delta_{B} \rightarrow 0$.

We will now show how the new measure is constructed from the function $W$. We inherit the following expression from [7, Prop. 3.2]:

$$
\begin{align*}
\bar{\lambda}(p) & =N(p) \lambda e^{-\left\langle p, v_{0}\right\rangle / 2} \\
\bar{\mu}_{1}(p) & =N(p) \mu_{1} e^{-\left\langle p, v_{1}\right\rangle / 2}  \tag{27}\\
\bar{\mu}_{2}(p) & =N(p) \mu_{2} e^{-\left\langle p, v_{2}\right\rangle / 2}
\end{align*}
$$

where

$$
\begin{equation*}
N(p):=\left[\lambda e^{-\left\langle p, v_{0}\right\rangle / 2}+\mu_{1} e^{-\left\langle p, v_{1}\right\rangle / 2}+\mu_{2} e^{-\left\langle p, v_{2}\right\rangle / 2}\right]^{-1}=e^{\mathbb{H}(p) / 2} \tag{28}
\end{equation*}
$$

Here $\mathbb{H}(p)$ is a function known as the Hamiltonian, which we use to simplify the notation and to enable the comparison with [7] and [4]. The vector $p$ strongly depends on the current state of the process and is in fact taken to be the gradient $D W(x)$. We thus rewrite (27) as

$$
\begin{align*}
& \bar{\lambda}(x)=\lambda e^{-\left\langle D W(x), v_{0}\right\rangle / 2} e^{\mathbb{H}(D W(x)) / 2}  \tag{29}\\
& \bar{\mu}_{i}(x)=\mu_{i} e^{-\left\langle D W(x), v_{i}\right\rangle / 2} e^{\mathbb{H}(D W(x)) / 2}, \quad i=1,2
\end{align*}
$$

We like to mention that we can express the gradient $D W(x)$ as a weighted sum of vectors $D W_{k}(x)$ at point $x$ :

$$
\begin{equation*}
D W(x)=\sum_{k=1}^{3} \rho_{k}(x) D W_{k}(x), \text { where } \rho_{k}(x)=\frac{e^{-W_{k}(x) / \varepsilon}}{\sum_{i=1}^{3} e^{-W_{i}(x) / \varepsilon}} \tag{30}
\end{equation*}
$$

For the Hamiltonian we have the following results.

Lemma 5.3. For any $x, \mathbb{H}\left(D W_{1}(x)\right)=\mathbb{H}\left(D W_{3}(x)\right)=0$ and $\mathbb{H}\left(D W_{2}(x)\right) \geq 0$.
Proof. The first and second claims are due to a direct computation:

$$
\begin{aligned}
& \mathbb{H}\left(D W_{1}(x)\right)=2 \log N\left(D W_{1}(x)\right) \\
& \quad=-2 \log \left[\lambda e^{-\log (\lambda / \tilde{\lambda}(x))}+\mu_{1} e^{-\log \left(\tilde{\mu}_{2}(x) / \mu_{2}\right)+\log (\lambda / \tilde{\lambda}(x))}+\mu_{2} e^{\log \left(\tilde{\mu}_{2}(x) / \mu_{2}\right)}\right] \\
& \quad=-2 \log \left[\tilde{\lambda}(x)+\tilde{\mu}_{1}(x)+\tilde{\mu}_{2}(x)\right]=0
\end{aligned}
$$

and

$$
\mathbb{H}\left(D W_{3}(x)\right)=-2 \log \left[\lambda+\mu_{1}+\mu_{2}\right]=0
$$

Finally, for the third case we have

$$
\begin{aligned}
\mathbb{H}\left(D W_{2}(x)\right) & =2 \log N\left(D W_{2}(x)\right) \\
& =-2 \log \left[\lambda e^{-\log \left(\lambda / \tilde{\lambda}\left(x_{1}, \delta / 2 \gamma\right)\right)}+\mu_{1} e^{\log \left(\lambda / \tilde{\lambda}\left(x_{1}, \delta / 2 \gamma\right)\right)}+\mu_{2} e^{0}\right] \\
& =-2 \log \left[\tilde{\lambda}\left(x_{1}, \delta / 2 \gamma\right)+\mu_{1} \frac{\lambda}{\tilde{\lambda}\left(x_{1}, \delta / 2 \gamma\right)}+\mu_{2}\right]
\end{aligned}
$$

To study the argument of the last logarithm, we consider the function $\psi(x):=x+\lambda \mu_{1} / x+$ $\mu_{2}$, for which we have $\psi(\lambda)=\psi\left(\mu_{1}\right)=1$. Also, $\psi(x)$ is convex, so that for all possible values of $\tilde{\lambda}\left(x_{1}, \delta / 2 \gamma\right)$ in $\left[\lambda, \mu_{2}\right] \subset\left[\lambda, \mu_{1}\right]$, one can conclude that $\psi\left(\tilde{\lambda}\left(x_{1}, \delta / 2 \gamma\right)\right) \leq 1$. This proves the last statement of this lemma.

Clearly there is a difference between the new measures defined in Section 3 (indicated by tildes) and in this section (indicated by bars). In fact it is not difficult to see that the first one also follows from (27) if we replace $W$ by $W_{1}$. However, this change of measure is not asymptotically efficient, while the other one is, due to the protection strips along the boundaries, as we will prove in the remainder of this subsection. We start with some lemmas that are similar to the ones in [4].

Lemma 5.4. The likelihood $L(A)$ of a path $A=\left(X_{j}, j=0, \ldots, \sigma\right)$ under the new measure (29) satisfies

$$
\begin{align*}
\log L(A)= & \frac{B}{2} \sum_{j=0}^{\sigma-1}\left\langle D W\left(X_{j}\right), X_{j+1}-X_{j}\right\rangle \\
& +\sum_{k=1}^{2} \frac{1}{2} \sum_{j=0}^{\sigma-1}\left\langle D W\left(X_{j}\right), v_{k}\right\rangle I\left\{X_{j}=X_{j+1} \in \partial_{k}\right\}  \tag{31}\\
& -\frac{1}{2} \sum_{j=0}^{\sigma-1} \mathbb{H}\left(D W\left(X_{j}\right)\right)
\end{align*}
$$

Proof. The proof is the same as that of Lemma 1 in [4].

Lemma 5.5. Consider the case $\mu_{2}<\mu_{1}$. For any path $A=\left(X_{j}, j=0, \ldots, \sigma\right)$ under the new measure (29), the first term in (31) satisfies

$$
\left|\frac{B}{2} \sum_{j=0}^{\sigma-1}\left\langle D W\left(X_{j}\right), X_{j+1}-X_{j}\right\rangle-\frac{B}{2}\left(W\left(X_{\sigma}\right)-W\left(X_{0}\right)\right)\right| \leq \frac{C}{B \varepsilon} \sigma,
$$

for sufficiently large $B \varepsilon$, where $C$ is some positive constant.
Proof. At first let us introduce the following representation

$$
W(x+y)=W(x)+\langle D W(x), y\rangle+\frac{1}{2} y^{T} H(x) y+|y|^{2} r(y),
$$

where $y=X_{j+1}-X_{j}$ is a one-step increment of the scaled process $X_{j}$, the matrix $H(x)$ is the Hessian matrix of the function $W(x)$ and the function $r(y)$ satisfies $\lim _{|y| \rightarrow 0} r(y)=0$. After transferring two terms to the left hand side and taking the absolute value we find

$$
\begin{aligned}
|(W(x+y)-W(x))-\langle D W(x), y\rangle| & =\left|\frac{1}{2} y^{T} H(x) y+|y|^{2} r(y)\right| \\
& \leq \frac{1}{2}\left|y^{T}\right| \cdot| | H(x) \|_{2} \cdot|y|+|y|^{2}|r(y)| \\
& \leq|y|^{2}\|H(x)\|_{\max }+|y|^{2}|r(y)|
\end{aligned}
$$

where $\|H(x)\|_{\text {max }}$ is the maximum norm of the Hessian matrix, given by

$$
\|H(x)\|_{\max }=\max \left\{h_{11}(x), h_{12}(x), h_{22}(x)\right\} ;
$$

here

$$
h_{11}(x):=\left|\frac{\partial^{2} W(x)}{\partial x_{1}^{2}}\right|, h_{12}(x):=\left|\frac{\partial^{2} W(x)}{\partial x_{1} \partial x_{2}}\right|, h_{22}(x):=\left|\frac{\partial^{2} W(x)}{\partial x_{2}^{2}}\right| .
$$

We now compute an upper bound for $\left|h_{11}(x)\right|$ as an example; the two other terms can be dealt with in the same manner. Using representation (30) one can write

$$
\begin{equation*}
\frac{\partial^{2} W(x)}{\partial x_{1}^{2}}=\sum_{k=1}^{3}\left[\rho_{k}(x) \frac{\partial^{2} W_{k}(x)}{\partial x_{1}^{2}}+\frac{\partial \rho_{k}(x)}{\partial x_{1}} \cdot \frac{\partial W_{k}(x)}{\partial x_{1}}\right], \tag{32}
\end{equation*}
$$

where it follows from the definition of $\rho_{k}(x)$ that

$$
\frac{\partial \rho_{k}(x)}{\partial x_{1}}=-\frac{1}{\varepsilon} \frac{\rho_{k}(x) \sum_{i \neq k} e^{-W_{i}(x) / \varepsilon}\left(\frac{\partial W_{k}(x)}{\partial x_{1}}-\frac{\partial W_{i}(x)}{\partial x_{1}}\right)}{\sum_{i} e^{-W_{i}(x) / \varepsilon}}
$$

Since the second fraction turns out to be bounded as $\varepsilon \rightarrow 0$, and the same holds for the other terms in (32), we find that some positive constant $C_{1}$ exists, such that

$$
\left|\frac{\partial^{2} W(x)}{\partial x_{1}^{2}}\right|<\frac{C_{1}}{\varepsilon} .
$$

Due to similar bounds for the other second-order partial derivatives, and the simple fact that $|y| \leq \sqrt{2} / B$, we have for some positive constant $C_{2}$ that

$$
|y|^{2}\|H(x)\|_{\max } \leq \frac{C_{2}}{B^{2} \varepsilon} .
$$

Finally, if we choose $B$ large enough (and hence $|y|$ small), we have for some positive constant $C_{3}$ that

$$
|y|^{2}|r(y)| \leq \frac{C_{3}}{B^{2}}
$$

The statement of the lemma is a direct consequence of these two bounds.
Lemma 5.6. Consider a two-node tandem Jackson network. For any sequence $\theta_{B}$ such that $\theta_{B} \rightarrow 0$ $(B \rightarrow \infty)$, and $\tau_{B}^{x}$ defined by (2), the following limit holds:

$$
\lim _{B \rightarrow \infty} \frac{1}{B} \log \mathbb{E}\left(\mathrm{e}^{\theta_{B} \tau_{B}^{x}} \mid I_{B}\left(A^{x}\right)=1\right)=0
$$

Proof. Let us first give a sketch of the proof. This proof consist of three steps: (i) we bound the length of any path from $\partial_{e}$ to the origin; (ii) using time reversibility arguments we show that the same bound applies to $\tau_{B}^{0}$; (iii) we show that the path of our interest is shorter (in stochastic sense) than $\left(\tau_{B}^{0} \mid I_{B}\left(A^{0}\right)=1\right)$.
(i) Let $\sigma_{B}$ be the length of the path from any state $(B, \alpha)$ to the origin, for any finite $\alpha$; and let $\omega^{B}$ be the length of the path from any state $\left(x_{1}, x_{2}\right)$, such that $x_{1}+x_{2}=B$. It is clear that

$$
\begin{equation*}
\lim _{B \rightarrow \infty} \frac{1}{B} \log \mathbb{E} e^{\theta_{B} \sigma_{B}}=\lim _{B \rightarrow \infty} \frac{1}{B+\alpha} \log \mathbb{E} e^{\theta_{B} \omega^{B+\alpha}}=\lim _{B \rightarrow \infty} \frac{1}{B} \log \mathbb{E} e^{\theta_{B} \omega^{B}}=0 \tag{33}
\end{equation*}
$$

where the last equality follows from the third statement of [4, Lemma 3].
(ii) Now consider the time-reversed network, see [10, Thm. 1.12]. It is not difficult to check that this is also a tandem queue, but with the first and second queue interchanged. The length of a path in the original system from the origin to level $B$ in the second queue, without visits to the origin in the mean time, equals the length of a path from some state $(B, \alpha)$ to the origin in the reversed system, given that it does not visit any state $(B, \cdot)$ in between, hence

$$
\left(\tau_{B}^{0} \mid I_{B}\left(A^{0}\right)=1\right) \leq_{\text {st }} \sigma_{B}
$$

Combining the last statement with (33) we have

$$
\begin{equation*}
\lim _{B \rightarrow \infty} \frac{1}{B} \log \mathbb{E}\left(e^{\theta_{B} \tau_{B}^{0}} \mid I_{B}\left(A^{0}\right)=1\right) \leq \lim _{B \rightarrow \infty} \frac{1}{B} \log \mathbb{E} e^{\theta_{B} \sigma_{B}}=0 \tag{34}
\end{equation*}
$$

which is similar to the fourth statement of [4, Lemma 3], only there the exit boundary $\partial_{e}$ was different.
(iii) We use stochastic coupling to show that

$$
\begin{equation*}
\left(\tau_{B}^{x} \mid I_{B}\left(A^{0}\right)=1\right) \leq_{s t}\left(\tau_{B}^{0} \mid I_{B}\left(A^{0}\right)=1\right) . \tag{35}
\end{equation*}
$$

To see this, we consider the ('original') process starting in the origin, and couple it to a similar process starting in state $x$. Then the above states that the time to overflow for the original process, given that overflow happens before reaching the origin is stochastically larger than the time to overflow for the coupled process, given that the original process reaches overflow before the origin. Notice that the condition implies that also the coupled process reaches overflow before the origin (since the queue lengths cannot be negative). In other words, for any path with $I_{B}\left(A^{0}\right)=1$, also $I_{B}\left(A^{x}\right)=1$ must hold, but since the opposite does not hold in general, we have $\left\{I_{B}\left(A^{0}\right)=1\right\} \subset\left\{I_{B}\left(A^{x}\right)=1\right\}$. From this we can conclude

$$
\begin{equation*}
\left(\tau_{B}^{x} \mid I_{B}\left(A^{x}\right)=1\right) \leq_{s t}\left(\tau_{B}^{x} \mid I_{B}\left(A^{0}\right)=1\right) \tag{36}
\end{equation*}
$$

Using (34), (35) and (36) we can now write that for any state $x \in \bar{D}$,

$$
\begin{equation*}
\lim _{B \rightarrow \infty} \frac{1}{B} \log \mathbb{E}\left(\mathrm{e}^{\theta_{B} \tau_{B}^{x}} \mid I_{B}\left(A^{x}\right)=1\right) \leq \lim _{B \rightarrow \infty} \frac{1}{B} \log \mathbb{E}\left(\mathrm{e}^{\theta_{B} \tau_{B}^{0}} \mid I_{B}\left(A^{0}\right)=1\right)=0 \tag{37}
\end{equation*}
$$

which completes the proof of the lemma.
Theorem 5.7. When $\mu_{2}<\mu_{1}$ and Assumption 5.2 holds, the new measure in (29), with function $W$ based on (15), is asymptotically efficient.

Proof. We will roughly follow the proof of [4, Thm. 1], paying attention to some important issues. First note that Lemma 5.3 provides an upper bound on the last term of the loglikelihood expression in Lemma 5.4:

$$
\begin{equation*}
-\frac{1}{2} \sum_{j=0}^{\tau_{B}^{x}-1} \mathbb{H}\left(D W\left(X_{j}\right)\right) \leq 0 \tag{38}
\end{equation*}
$$

In order to bound the second term in Lemma 5.4 we will prove a result similar to the third statement of [7, Lemma B.1].
From Proposition 5.1 we know that $\left\langle D W_{2}(x),-v_{2}\right\rangle=\left\langle D W_{3}(x),-v_{2}\right\rangle=0$ and also that $\left\langle D W_{1}(x),-v_{2}\right\rangle=2 \log \left(\tilde{\mu}_{2}(x) / \mu_{2}\right)$. Hence, applying (30), we have

$$
\begin{equation*}
\left\langle D W(x),-v_{2}\right\rangle=2 \log \left(\frac{\tilde{\mu}_{2}(x)}{\mu_{2}}\right) \rho_{1}(x) \geq 2 \log \left(\frac{\tilde{\mu}_{2}(x)}{\mu_{2}}\right) e^{-\left(W_{1}(x)-W_{2}(x)\right) / \varepsilon} \tag{39}
\end{equation*}
$$

It is clear that $W_{1}(x)-W_{2}(x)=\delta$ for any $x \in A_{1} \cap \partial_{2}$, see (25). Also, Lemma 3.4 guarantees that $W_{1}(x)-W_{2}(x)$ decreases to 0 as $x$ moves along the horizontal axis from $\left(\alpha_{1}, 0\right)$ to $\left(\alpha_{1}^{-1}, 0\right)$. This immediately leads to $0 \leq W_{1}(x)-W_{2}(x) \leq \delta$ for any $x \in \partial_{2}$. Now keeping in mind that $\tilde{\mu}_{2}(x) \geq \lambda$ (see again Lemma 3.4) and hence $\log \left(\tilde{\mu}_{2}(x) / \mu_{2}\right) \geq-\gamma$, we can write

$$
\left\langle D W(x),-v_{2}\right\rangle \geq-2 \gamma e^{-\frac{\delta}{\varepsilon}}
$$

Using the same technique and keeping Lemma 3.4 in mind one can also show that

$$
\left\langle D W(x),-v_{1}\right\rangle \geq-2 \gamma e^{-\frac{\delta}{\varepsilon}}
$$

for any $x \in \partial_{1}$. Using these two inequalities we obtain the same bound for the second term in Lemma 5.4 as in [4]:

$$
\begin{equation*}
\sum_{k=1}^{2} \frac{1}{2} \sum_{j=0}^{\tau_{B}^{x-1}}\left\langle D W\left(X_{j}\right), v_{k}\right\rangle I\left\{X_{j}=X_{j+1} \in \partial_{k}\right\} \leq \gamma e^{-\delta / \varepsilon} \tau_{B}^{x} \tag{40}
\end{equation*}
$$

To deal with the first term in Lemma 5.4, we first bound $W(x)$ using (26):

$$
W(x) \leq-\varepsilon \log \left(e^{W_{1}(x) / \varepsilon}\right)=-\varepsilon \log \left(e^{\left(-2 \gamma_{x}+\delta\right) /(\varepsilon)}\right)=2 \gamma_{x}-\delta
$$

and, using that $W_{2}(x) \geq W_{1}(x)-\delta$ and the monotonicity of $\gamma_{x}$,

$$
\begin{aligned}
W(x) & \geq-\varepsilon \log \left(e^{-W_{1}(x) / \varepsilon}+e^{\left(-W_{1}(x)+\delta\right) / \varepsilon}+e^{-W_{3}(x) / \varepsilon}\right) \\
& \geq-\varepsilon \log \left(3 e^{\left(-2 \gamma_{x}+3 \delta\right) / \varepsilon}\right)=2 \gamma_{x}-\varepsilon \log (3)-3 \delta
\end{aligned}
$$

Using the same technique we obtain similar bounds for $W\left(X_{\tau_{B}^{x}}\right)$ :

$$
-\varepsilon \log (3)-3 \delta \leq W\left(X_{\tau_{B}^{x}}\right) \leq-\delta
$$

Using the three last inequalities and Lemma 5.5 we can derive an upper bound for the first term in Lemma 5.4,

$$
\begin{equation*}
\frac{B}{2} \sum_{j=0}^{\tau_{B}^{x}-1}\left\langle D W\left(X_{j}\right), X_{j+1}-X_{j}\right\rangle \leq \frac{B}{2}\left(-2 \gamma_{x}+\eta(B)\right)+\frac{C}{B \varepsilon} \tau_{B}^{x} \tag{41}
\end{equation*}
$$

where $\eta(B)$ is such that $\lim _{B \rightarrow \infty} \eta(B)=0$.
Combining (38), (40) and (41) we can rewrite (31) in the following way

$$
\log (L(A)) \leq-B \gamma_{x}+B \eta(B)+\chi(B) \tau_{B}^{x}
$$

where

$$
\chi(B):=\gamma e^{-\delta / \varepsilon}+\frac{C}{B \varepsilon} .
$$

Now for any path $A^{x}$ we have:

$$
\begin{aligned}
& \frac{1}{B} \log \mathbb{E}\left[L\left(A^{x}\right) I_{B}\left(A^{x}\right)\right]=\frac{1}{B} \log \left(\mathbb{E}\left[L(A) \mid I_{B}\left(A^{x}\right)=1\right] \mathbb{P}\left[I_{B}\left(A^{x}\right)=1\right]\right) \\
& \quad \leq \frac{1}{B} \log \left(\mathbb{E}\left[e^{-B \gamma_{x}+B \eta(B)+\chi(B) \tau_{B}^{x}} \mid I_{B}\left(A^{x}\right)=1\right] p_{B}^{x}\right) \\
& \quad=-\gamma_{x}+\eta(B)+\frac{1}{B} \log \mathbb{E}\left[e^{\chi(B) \tau_{B}^{x}} \mid I_{B}\left(A^{x}\right)=1\right]+\frac{1}{B} \log p_{B}^{x} .
\end{aligned}
$$

Using the fact that $\lim _{B \rightarrow \infty} \chi(B)=0$ (see Assumption 5.2), Lemma 5.6 and Theorem 4.5 we conclude that:

$$
\limsup _{B \rightarrow \infty} \frac{1}{B} \log \mathbb{E}\left[L\left(A^{x}\right) I_{B}\left(A^{x}\right)\right] \leq-2 \gamma_{x}=2 \lim _{B \rightarrow \infty} \frac{1}{B} \log p_{B}^{x}
$$

which completes the proof.

### 5.2 Asymptotically efficient scheme for $\mu_{1} \leq \mu_{2}$

We would like to define a function based on the total cost function $\gamma_{x}$ in (19), analogous to the function $W$ in the previous section, see (26). Suppose we define the functions $V_{i}(x)$, $i=1,2,3$, in the same way as (25). In particular we would find $\hat{V}_{1}$ on $B_{1}$ and $B_{2}$ from (19) as,

$$
\begin{align*}
& \hat{V}_{1}(x)=2 \gamma_{\left(x_{1}, 0\right)}-\delta, \quad x \in B_{1}  \tag{42}\\
& \hat{V}_{1}(x)=-2 x_{1} \log \frac{\tilde{\lambda}(x)}{\lambda}-2\left(1-x_{2}\right) \log \frac{\tilde{\mu}_{2}(x)}{\mu_{2}}-\delta, \quad x \in B_{2} . \tag{43}
\end{align*}
$$

where $\left(\tilde{\lambda}(x), \tilde{\mu}_{1}(x), \tilde{\mu}_{2}(x)\right)$ is the solution to (11) if $x \in B_{2}$. As a result, $\hat{V}_{1}(x)$ would not be smooth on the boundary between the sets $B_{1}$ and $B_{2}$ (see also the discussion above (18)), and hence also the resulting mollified function would not be smooth. This would lead to problems when we try to prove the analogue of Lemma 5.5 , where we used continuity and smoothness of $W(x)$. Fortunately, the functions $\hat{V}_{1}(x)$ and $\hat{V}_{2}(x)$ coincide on $B_{1}$ since they are equal on the boundary between $B_{1}$ and $B_{2}$ and both functions do not depend on their second argument (hence their gradients coincide). Hence, instead of using (42)-(43) we prefer to work with a function $V_{1}$ defined as (43) on both $B_{1}$ and $B_{2}$. Mollifying this function with functions $V_{2}$ and $V_{3}$ as in (26), will then provide a smooth function $V$. To be more specific, we first define function $V_{1}(x)$, based on the second line of (19):

$$
V_{1}(x)=-2 x_{1} \log \frac{\tilde{\lambda}(x)}{\lambda}-2\left(1-x_{2}\right) \log \frac{\tilde{\mu}_{2}(x)}{\mu_{2}}-\delta
$$

where

$$
\left(\tilde{\lambda}(x), \tilde{\mu}_{1}(x), \tilde{\mu}_{2}(x)\right)= \begin{cases}\text { solution to (11) } & \text { if } x \in B_{1} \cup B_{2} \\ \left(\lambda, \mu_{2}, \mu_{1}\right) & \text { if } x \in B_{3}\end{cases}
$$

The function $V_{1}(x)$ is an extension of the 'cost' proposed by the solution of (11) to the set $B_{1}$. In other words, for any $x \in B_{1}$ we replace the optimal cost (corresponding to the first type of path in (18)) by the cost that corresponds to the (suboptimal) path that leads straight from $x$ to $(0,1)$. We proceed with the definitions of $V_{2}(x)$ and $V_{3}(x)$ :

$$
\begin{aligned}
& V_{2}(x)=2 \gamma_{\left(x_{1}, \delta / 2 \gamma\right)}-\delta, \\
& V_{3}(x)=2 \gamma-3 \delta,
\end{aligned}
$$

where $\gamma_{x}$ is given in (19). In this way, the minimum of $V_{1}$ and $V_{2}$ is attained by $V_{1}$ for $x \in B_{2}$ (as before), and by $V_{2}$ for $x \in B_{1}$ (rather than by $\hat{V}_{1}$ as before); see also Figure 5, where $\bar{V}(x)=V_{1}(x) \wedge V_{2}(x) \wedge V_{3}(x)$. The mollification procedure now ensures a smooth transition from $B_{1}$ to $B_{2}$ for the function $V(x)$ defined as in (26). Another minor problem is that the function $V_{2}(x)$ is not smooth around $\left(\alpha_{2}, 0\right)$. Specifically for $x_{2}<\delta / \gamma_{2}$, the gradient of $V_{2}(x)$ is not continuous around the vertical line $\left(x_{1}, \cdot\right)$ where the first component satisfies $f\left(x_{1}, \delta / 2 \gamma\right)=0$ with $f(x)$ as defined in (17). Without going into details, we propose to use any suitable mollification procedure to make $V_{2}(x)$ a smooth function, and from now on we will treat $V_{2}(x)$ as such. Thus, mollifying the functions $V_{i}(x), i=1,2,3$, in the


Figure 5: The function $\bar{V}(x)$ and the areas on which $\bar{V}(x)=V_{i}, i=1,2,3$ (case $\mu_{1} \leq \mu_{2}$ ).
same manner as we did in (26), we obtain a smooth and continuous function $V(x)$. Based on this function we define the new change of measure $\left(\bar{\lambda}, \bar{\mu}_{1}, \bar{\mu}_{2}\right)$ as in (29). As for the gradient in these equations, we like to notice that even though the functions $W_{i}(x)$ and $V_{i}(x)$ for $i=1,2,3$ are different, they have similar gradients. In other words, we can replace $D W_{i}(x)$ by $D V_{i}(x)$ in Proposition 5.1 to obtain the shape of the gradients of the functions $V_{i}(x), i=1,2,3$. We will use these gradients in the following proofs.
Turning to the asymptotic efficiency proof of this new change of measure, we first mention that we can prove analogues of Lemmas 5.3, 5.4, 5.5 (using the smoothness of $V$ ) and 5.6 for our current case $\mu_{1} \leq \mu_{2}$, as can be checked easily. Now we can proceed with the main result of this subsection.

Theorem 5.8. When $\mu_{1} \leq \mu_{2}$ and Assumption 5.2 holds, the new measure in (29) with function $V$ based on (19), is asymptotically efficient.

Proof. The proof is similar to that of Theorem 5.7, the main difference being the bound on the second term of the decomposition of $\log L(A)$ in Lemma 5.4. For any $x \in \partial_{2}$, we have $\left\langle D V_{2},-v_{2}\right\rangle=\left\langle D V_{3},-v_{2}\right\rangle=0$ and $\left\langle D V_{1},-v_{2}\right\rangle=2 \log \left(\tilde{\mu}_{2}(x) / \mu_{2}\right)$, so

$$
\left\langle D V(x),-v_{2}\right\rangle \geq 2 \log \left(\frac{\tilde{\mu}_{2}(x)}{\mu_{2}}\right) e^{-\left(V_{1}(x)-V_{2}(x)\right) / \varepsilon},
$$

as in the previous case. For $x=(0,0)$ we have that

$$
V_{1}(0,0)-V_{2}(0,0)=2 \log (1 / z)-2 \gamma .
$$

Using the fact that the optimal cost $\gamma_{x}$ for state $x=(0, \beta)$ is both equal to $\gamma$ (corresonding to the path via state $\left(\alpha_{2}, 0\right)$ ) and to $(1-\beta) \log (1 / z)$ (corresponding to the path along the vertical axis), this can be rewritten as $V_{1}(0,0)-V_{2}(0,0)=2 \beta \log (1 / z)$. Also, the difference $V_{1}(x)-V_{2}(x)$ decreases in $x_{1}$ when $x \in \partial_{2}$, see Lemma 3.5. Combining the above with the fact that $\tilde{\mu}_{2}(x) \geq \mu_{2} z$ (see Lemma 3.5) we obtain the following bound:

$$
\left\langle D V(x),-v_{2}\right\rangle \geq 2 \log (z) \exp \left(-\frac{2 \beta \log (1 / z)}{\varepsilon}\right)
$$

for any $x \in \partial_{2}$.
Now let us consider the situation when $x \in \partial_{1}$. Again $\left\langle D V_{3}(x),-v_{1}\right\rangle=0$ holds, and in addition $\left\langle D V_{2}(x),-v_{1}\right\rangle=-2 \log \frac{\mu_{1}}{\lambda}$ and $\left\langle D V_{1}(x),-v_{1}\right\rangle=2 \log (\lambda / \tilde{\lambda}(x))-2 \log \left(\tilde{\mu}_{2}(x) / \mu_{2}\right)>$ 0 , where the last inequality is due to the simple observation that $\sqrt{\lambda z / \mu_{1}}>z$. Using (30) we have

$$
\left\langle D V(x),-v_{1}\right\rangle \geq\left\langle D V_{2}(x),-v_{1}\right\rangle=-2 \log \left(\mu_{1} / \lambda\right) \rho_{2}(x) \geq-2 \log \left(\mu_{1} / \lambda\right) e^{\left(V_{3}(x)-V_{2}(x)\right) / \varepsilon}
$$

Since for any $x \in \partial_{1}$ we have $V_{3}(x)-V_{2}(x)=-2 \delta$, we conclude that

$$
\left\langle D V(x),-v_{1}\right\rangle \geq-2 \log \left(\mu_{1} / \lambda\right) e^{-2 \delta / \varepsilon}
$$

Analogously to the previous proof we now conclude that

$$
\limsup _{B \rightarrow \infty} \frac{1}{B} \log \mathbb{E}\left[L\left(A^{x}\right) I_{B}\left(A^{x}\right)\right] \leq-2 \gamma_{x}=2 \lim _{B \rightarrow \infty} \frac{1}{B} \log p_{B}^{x}
$$

which completes the proof.

## 6 Numerical results

We provide some supporting simulation results in this section. To simplify the implementation of the IS scheme, we introduce a slightly different new measure $\mathbb{Q}$, defined as follows:

$$
\begin{align*}
\breve{\lambda}(x) & =\lambda \sum_{k=1}^{3} \rho_{k}(x) e^{-\left\langle D W_{k}(x), v_{0}\right\rangle / 2} e^{\mathbb{H}\left(D W_{k}(x)\right) / 2}  \tag{44}\\
\breve{\mu}_{i}(x) & =\mu_{i} \sum_{k=1}^{3} \rho_{k}(x) e^{-\left\langle D W_{k}(x), v_{i}\right\rangle / 2} e^{\mathbb{H}\left(D W_{k}(x)\right) / 2}, \quad i=1,2 \tag{45}
\end{align*}
$$

where $\rho_{k}(x)$ is defined by (30).
Theorem 6.1. Under Assumption 5.2,
(i) when $\mu_{2}<\mu_{1}$, the measure $\mathbb{Q}$ in (44)-(45), with $W$ defined by (26), is asymptotically efficient.
(ii) when $\mu_{1} \leq \mu_{2}$, the measure $\mathbb{Q}$ in (44)-(45), with $W$ replaced by the function $V$ as defined in Section 5.2 , is asymptotically efficient.

Proof. (i) It is clear that the log-likelihood ratio for a transition of type $v_{0}$ from any state $x$ under $\mathbb{Q}$ satisfies

$$
\begin{aligned}
\log \frac{\lambda}{\breve{\lambda}(x)} & =-\log \sum_{k=1}^{3} \rho_{k}(x) e^{-\left\langle D W_{k}(x), v_{0}\right\rangle / 2} e^{\mathbb{H}\left(D W_{k}(x)\right) / 2} \\
& \leq-\sum \rho_{k}(x) \log e^{-\left\langle D W_{k}(x), v_{0}\right\rangle / 2}=\left\langle D W(x), v_{0}\right\rangle / 2
\end{aligned}
$$

| $(0.6 B, 0)$ |  |  | $(B, 0)$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $B$ | $\psi_{B}$ | $p_{B}$ | $R E$ | $B$ | $\psi_{B}$ | $p_{B}$ | $R E$ |
| 20 | 1.96 | $2.00 \cdot 10^{-5} \pm 2.74 \cdot 10^{-8}$ | $7.01 \cdot 10^{-4}$ | 20 | 1.92 | $1.29 \cdot 10^{-2} \pm 1.61 \cdot 10^{-5}$ | $6.38 \cdot 10^{-4}$ |
| 50 | 1.98 | $3.12 \cdot 10^{-12} \pm 4.69 \cdot 10^{-15}$ | $7.67 \cdot 10^{-4}$ | 50 | 1.96 | $3.95 \cdot 10^{-5} \pm 5.37 \cdot 10^{-8}$ | $7.20 \cdot 10^{-4}$ |

Table 1: Simulation results; original state-dependent scheme
where the last inequality holds due to the fact that $\mathbb{H}\left(D W_{k}(x)\right) \geq 0$, thanks to Lemma 5.3 , and concavity of the logarithm (note that $\sum_{k=1}^{3} \rho_{k}(x)=1$ ). It is obvious that we have similar bounds for transitions $v_{1}$ and $v_{2}$. Summing these expressions over all steps of sample path $A=\left(X_{j}, j=0 \ldots \sigma\right)$ we will get the righthand side of expression (31), but without the last term. Since the function $W(x)$ stays the same, we may use the proof of Theorem 5.7 to verify the statement of the current problem.
(ii) The second claim is proved analogously to the first claim.

All simulations were performed under new measure $\mathbb{Q}$ defined by (44)-(45) and the joint queue-length process around the boundaries was modified according to (1).
Here we present results of dynamic IS simulations for the two-node Jackson tandem network with initial parameters $\left(\lambda, \mu_{1}, \mu_{2}\right)=(0.1,0.55,0.35)$. We restrict ourselves to the case when the second buffer is the bottleneck (i.e., $\mu_{2}<\mu_{1}$ ), since the optimal path and new measure $\mathbb{Q}$ are very similar for both cases when the initial state $x$ lies in $A_{2} \cup A_{3}$ (respectively, $B_{2} \cup B_{3}$ ); only when $x \in B_{1}$ (for the case $\mu_{1} \leq \mu_{2}$ ), there is an interesting difference with the other case ( $\mu_{1} \geq \mu_{2}$ ), but the shape of the optimal path and corresponding new measure are then very similar to the ones studied before in [13].
We have two different starting states: $(0.6 B, 0)$ and $(B, 0)$. Both belong to the most interesting subspace $A_{2}$, see Figure 2, where the new measure is found as the solution to system (11).
We performed three different types of IS simulations. At first we simulate the system based on the asymptotic efficient state-dependent scheme obtained in Section 5, i.e., scheme (12) with protection along the horizontal boundary. In these (and all further) simulations we chose $\varepsilon=0.05$ and $\delta=-\varepsilon \log \varepsilon$, as motivated in Remark 3.7 in [7]. Moreover, we always performed $10^{6}$ simulation runs, leading to comparable computation times in the order of a few minutes; in fact these were approximately linear in the value of $B$, as could be expected.
The results are presented in Table 1. The value $\psi_{B}$ in the second column of each table (for both panels) is the estimator of the right hand side of (6), without the limit, and is used as an indication for the efficiency of the scheme. The obtained results are indeed good (that is, close to 2), but it is mentioned that the computation time needed is considerable when $B$ grows large, due to the precalculation of the new measure $\mathbb{Q}$ for all states in $A_{2}$.
We tried to simplify the scheme to reduce the computation time, in the following way. We divide the set $A_{2}$ into a number of triangles $D_{i}$ of equal area, each having state $(0,1)$ as one of the corners, the other corners being given by points on the horizontal axis between $\left(\alpha_{2}, 0\right)$ and $\left(\alpha_{2}^{-1}, 0\right)$, at equal distances. In each of these subsets $D_{i}$ we use a separate,
fixed, new measure, based on the solution of system (11) where $x$ lies in the middle of the two corners on the horizontal axis. In this way we only need to precalculate a few new measures, rather than dozens for $B=50$ and hundreds for $B=100$. In Table 2 the simulation results are given, when $A_{2}$ is divided into six subsets. Due to this precalculation reduction it became possible to add the extra line $(B=100)$ in Table 2 (and Table 3).

| $(0.6 B, 0)$ |  |  |  | $(B, 0)$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $B$ | $\psi_{B}$ | $p_{B}$ | $R E$ | $B$ | $\psi_{B}$ | $p_{B}$ | $R E$ |
| 20 | 1.95 | $2.00 \cdot 10^{-5} \pm 3.04 \cdot 10^{-8}$ | $7.77 \cdot 10^{-4}$ | 20 | 1.86 | $1.29 \cdot 10^{-2} \pm 2.28 \cdot 10^{-5}$ | $8.94 \cdot 10^{-4}$ |
| 50 | 1.97 | $3.12 \cdot 10^{-12} \pm 6.11 \cdot 10^{-15}$ | $9.98 \cdot 10^{-4}$ | 50 | 1.91 | $3.94 \cdot 10^{-5} \pm 9.36 \cdot 10^{-8}$ | $1.20 \cdot 10^{-3}$ |
| 100 | 1.98 | $1.82 \cdot 10^{-23} \pm 5.06 \cdot 10^{-26}$ | $1.41 \cdot 10^{-3}$ | 100 | 1.93 | $3.66 \cdot 10^{-9} \pm 1.13 \cdot 10^{-11}$ | $1.57 \cdot 10^{-3}$ |

Table 2: Simulation results; simplified scheme with six domains

As we can see, the variance of the estimator increases as a result of the simplification of the original scheme (although the effect is relatively modest). This may be explained by the following. Under the simplified scheme a path that starts to, say, the left of the middle point of $D_{i}$ will at some point hit the boundary between $D_{i}$ and $D_{i-1}$, after which it follows this boundary. Hence, when $B$ grows large and therefore the sizes of (the unscaled counterparts of) the $D_{i}$ also increase, the sample path will move back and forth between two different changes of measure for a substantial period of time.
We also simulated the system using an even simpler scheme, getting rid of the set $A_{2}$ altogether, expanding $A_{1}$ and $A_{3}$ so that they meet at the line $x_{1}+x_{2}=1$. That is, we simply used the measure $\mathbb{Q}=\left(\mu_{2}, \mu_{1}, \lambda\right)$ when the total population of the system is less than $B$ and used no change of measure otherwise. Clearly, this method provides worse

| $(0.6 B, 0)$ |  |  | $(B, 0)$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $B$ | $\psi_{B}$ | $p_{B}$ | $R E$ | $B$ | $\psi_{B}$ | $p_{B}$ | $R E$ |
| 20 | 1.86 | $2.00 \cdot 10^{-5} \pm 7.30 \cdot 10^{-8}$ | $1.87 \cdot 10^{-3}$ | 20 | 1.10 | $1.29 \cdot 10^{-2} \pm 1.78 \cdot 10^{-4}$ | $7.02 \cdot 10^{-3}$ |
| 50 | 1.91 | $3.12 \cdot 10^{-12} \pm 1.67 \cdot 10^{-14}$ | $2.73 \cdot 10^{-3}$ | 50 | 1.18 | $3.67 \cdot 10^{-5} \pm 4.34 \cdot 10^{-6}$ | $6.00 \cdot 10^{-2}$ |
| 100 | 1.94 | $1.81 \cdot 10^{-23} \pm 1.59 \cdot 10^{-25}$ | $4.48 \cdot 10^{-3}$ | 100 | 1.30 | $5.62 \cdot 10^{-9} \pm 7.09 \cdot 10^{-9}$ | $6.40 \cdot 10^{-1}$ |

Table 3: Simulation results; simplified scheme with two domains
results, as can be concluded from Table 3. Looking at the relative error or at the confidence intervals, it is clear that this method is inferior to the ones presented above.
Finally, we tried a change of measure that replaces the original parameters on $A_{2}$ by a simple linear interpolation between the values on $A_{1}$ and $A_{3}$ (e.g., when $\mu_{2}<\mu_{1}$ and $x \in A_{2}$ we let $\breve{\lambda}(x)=\alpha(x) \mu_{2}+(1-\alpha(x)) \lambda$, where $\alpha(x) \in[0,1]$ depends on the location of $x$ relative to $A_{1}$ and $A_{3}$ ). Unfortunately, this approach gives only slightly better results than those in Table 3, but much worse than those in Table 2.

## 7 Conclusions

In this paper we focused on the event that, starting from an arbitrary state, the second queue in a two-node Jackson tandem network reaches overflow before the system becomes
empty. The main focus is on the development of efficient simulation techniques for estimating this probability. We have proposed a particular change of measure, motivated by large-deviations arguments, and we have proved asymptotic efficiency of a subtly modified version (that differs close to the axes, and thus nicely controls the likelihood).
We strongly feel that the methods used in the current paper are applicable to other, more complex queueing networks. For example, we expect that it can be applied to a so-called 'slow-down network', i.e., a tandem network with Poisson arrivals and exponential service times, in which the first server decreases its speed as soon as the second buffer reaches some prescribed utilization, see [17]. Such an analysis has recently been published in [6] for a specific parameter setting, with the origin as starting state, but several issues remain open (general parameter settings, general initial point, simplification of the asymptotic efficiency proof).

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