Special Functions Associated with Root Systems:

Recent Progress

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1. Introduction

It is a pleasure for me to contribute to this farewell bundle for Cor Baayen. During his long period of involvement with the Mathematical Centre he has demonstrated from the beginning a wide interest in all kinds of pure mathematics, see his publications and colloquium contributions on many different topics during the sixties. During his years as a head of the Department of Pure Mathematics (ZW) he will have been aware that some pure math was being done within the MC but outside his department. In the late seventies I joined his department, but quite soon Cor was then called for a higher post: director of the whole institute. Next the departments of Pure and of Applied Mathematics joined into the new Department of Analysis, Algebra and Geometry, headed by Michiel Hazewinkel. In his new position, Cor had to deal with a much bigger world than just pure math. It included all kinds of applied math and a very large amount of computer science. Forced by national science policy and by financial constraints he had to take some unfavorable decisions concerning the heritage of his old ZW department. Fortunately, on a countrywide scale, he could give pure (or theoretical) mathematics some important financial injections via the foundation SMC (mathematics research money flow to the universities), of which he was automatically also a director. In my opinion Cor was a good science manager, showing a real interest in the research being done in his institute and in the people performing this research. In this paper I present a brief survey of the active area of Special Functions associated with Root Systems. The article is intended for a general mathematical audience. It will not suppose prerequisites on either special functions or root systems. It will also skip many technical details. Some early developments in this area took place at the Mathematical Centre during the seventies ([17], [32]), and some recent developments ([30], [18]) as well. During the last ten years important break-throughs were made by Heckman and Opdam (Leiden; Heckman later in Nijmegen) [11], [27], [28], [12]. Abroad, I. G. Macdonald [23], [24], C. F. Dunkl [8] and I. Cherednik [3], [4], [5] greatly contributed to the subject. A special period at the Stieltjes Institute (physically at the Universities of Leiden and Amsterdam) was devoted to this subject in the spring of 1994. The subject is also an important theme within the four-year country-wide project "Lie Theory and Special Functions", which just started and which is sponsored by SMC.

A lot of the motivation for the subject of this paper comes from analysis on semisimple Lie groups. Spherical functions on Riemannian symmetric spaces of the compact or non-compact type can be written as special functions depending on parameters which assume only special discrete values. In the one-variable cases these special functions are classical, also for parameter values without group theoretic interpretation, but in the more-variable cases they were new. In the case of group theoretic interpretation, many properties of these special functions, as well as associated harmonic analysis, immediately follow by group theoretic arguments. The case of more general parameter values yields the special functions associated with root systems. Properties derived in the group case can still be formulated in the general case, but now as conjectures rather than theorems. This paper describes some of the progress which has been made in proving these conjectures. For convenience, I will restrict to the polynomial (compact) case, with Bessel functions as a sole exception. (For an introduction to the non-polynomial case see [29], [13].) Neither will I discuss the recent work on commuting operators with elliptic functions as coefficients. An important aspect of the whole theory, which will not be discussed very much in this paper, is the connection with completely integrable systems, for instance the generalized Calogero-Moser system. Special functions associated with root systems have also been developed in the q-case, where q is a deformation parameter giving back the earlier cases when q = 1. Motivation and development of the theory in the q-case has been quite different from the q = 1 case. Except for the case of Hall polynomials [22], theory was developed [23], [24], [18] without interpretation in group theory. But afterwards quantum groups looked very promising as a natural setting for these polynomials. This had turned out to be true in the one-variable case [19], and very recently some interpretations of more-variable cases on quantum groups were found [25], [10]. In any case, a quantum group interpretation for generic values of the parameters cannot be expected. But, by Cherednik's work [3], [4], [5] we know already another algebraic setting for special functions associated with root systems: affine and graded Hecke algebras [20]. As shown by work of Opdam [29], this new algebraic context also allows harmonic analysis.

2. The one-variable case

In this section I will introduce three classical families of special functions, each depending on a real parameter $k \ge 0$, and such that the cases k = 0 and k = 1 are elementary. The three families are connected with each other by limit transitions. Later, for each of the families I will discuss generalizations which are associated with root systems.

2.1. Bessel functions. Consider Bessel functions in a non-standard notation:

$$\mathcal{J}_k(x) := \sum_{j=0}^{\infty} \frac{(-\frac{1}{4}x^2)^j}{(k+\frac{1}{2})_j j!} \quad (x \in \mathbb{R}).$$
(2.1)

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Here we use the notation for shifted factorial:

$$(a)_j := a (a + 1) \dots (a + j - 1) \quad (j = 1, 2, \dots); \qquad (a)_0 := 1.$$

The function \mathcal{J}_k is related to the Bessel function J_{α} in standard notation [9, Ch. VII] by

$$\mathcal{J}_k(x) = \frac{2^{k-\frac{1}{2}} \Gamma(k+\frac{1}{2})}{k-\frac{1}{2}} J_{k-\frac{1}{2}}(x).$$

$$x^{n-2}$$
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Note that

$$\mathcal{J}_k(x) = \mathcal{J}_k(-x), \qquad \mathcal{J}_k(0) = 1. \qquad (2.2)$$

The cases k = 0 and k = 1 yield elementary functions:

$$\mathcal{J}_0(x) = \cos x, \qquad \mathcal{J}_1(x) = \frac{\sin x}{x}. \qquad (2.3)$$

The function $x \mapsto \mathcal{J}_k(\lambda x)$ ($\lambda \in \mathbb{R}$) is eigenfunction of a differential operator:

$$\left(\frac{d^2}{dx^2} + \frac{2k}{x}\frac{d}{dx}\right) \mathcal{J}_k(\lambda x) = -\lambda^2 \mathcal{J}_k(\lambda x).$$

It is the unique C^{∞} solution of this differential equation under conditions (2.2).

2.2. Ultraspherical polynomials. Consider ultraspherical or Gegenbauer polynomials [9, §10.9], i.e. polynomials C_n^k of degree n on \mathbb{R} such that

$$\int_0^{\pi} C_n^k(\cos x) C_m^k(\cos x) (\sin x)^{2k} dx = 0 \quad (n, m \in \mathbb{Z}_+, \ n \neq m).$$

Then the C_n^k are determined up to a constant factor (in general, we will not use the standard normalization for Gegenbauer polynomials). For k = 0, 1 we have:

$$C_n^0(\cos x) = \text{const. } \cos(nx), \qquad C_n^1(\cos x) = \text{const. } \frac{\sin((n+1)x)}{\sin x}.$$
 (2.4)

The function $x \mapsto C_n^k(\cos x)$ is eigenfunction of a differential operator:

$$\left(\frac{d^2}{dx^2} + 2k \cot x \, \frac{d}{dx}\right) \, C_n^k(\cos x) = -n(n+2k) \, C_n^k(\cos x).$$

For (n_N) being a sequence of positive integers such that $n_N/N \to \lambda$ for some $\lambda \ge 0$ as $N \to \infty$, we have the limit result

$$\lim_{n \to \infty} \frac{C_{n_N}^k(\cos(x/N))}{C_{n_N}^k(1)} = \mathcal{J}_k(\lambda x).$$

2.3. *q*-Ultraspherical polynomials. Let 0 < q < 1 and define for any complex *a*:

$$(a;q)_{\infty} := \prod_{j=0}^{\infty} (1 - aq^j).$$

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The infinite product converges because of the condition on q. We will consider *q*-ultraspherical polynomials [1] in a non-standard notation. These are polynomials $C_n^{k,q}$ of degree n on \mathbb{R} such that

$$\int_0^{\pi} C_n^{k,q}(\cos x) C_m^{k,q}(\cos x) \left| \frac{(e^{2ix};q)_{\infty}}{(q^k e^{2ix};q)_{\infty}} \right|^2 \, dx = 0 \quad (n,m \in \mathbb{Z}_+, \ n \neq m).$$

Then the $C_n^{k,q}$ are determined up to a constant factor. If we put

$$P_n(e^{ix}) := C_n^{k,q}(\cos x)$$

then P_n is eigenfunction of a q-difference operator:

$$\frac{1-q^k e^{2ix}}{1-e^{2ix}} P_n(q^{\frac{1}{2}}e^{ix}) + \frac{1-q^k e^{-2ix}}{1-e^{-2ix}} P_n(q^{-\frac{1}{2}}e^{ix}) = (q^{-\frac{1}{2}n} + q^{\frac{1}{2}n+k}) P_n(e^{ix}).$$

Note that the P_n on the left hand side have arguments off the unit circle, while orthogonality is on the unit circle. The cases k = 0 and k = 1 are elementary as in (2.4) (not depending on q):

$$C_n^{0,q}(\cos x) = \text{const. } \cos(nx), \qquad C_n^{1,q}(\cos x) = \text{const. } \frac{\sin((n+1)x)}{\sin x}.$$

With suitable normalization there is the limit relation

$$\lim_{q \uparrow 1} C_n^{k,q}(\cos x) = C_n^k(\cos x).$$

The q-ultraspherical polynomials form a subclass of the Askey-Wilson polynomials [2]: a family of orthogonal polynomials depending, apart from q, on four non-trivial parameters.

2.4. Dunkl operators in one variable. We will now generalize the elementary formulas

$$e^{i\lambda x} = \mathcal{J}_0(\lambda x) + i\lambda x \,\mathcal{J}_1(\lambda x) \quad \text{and} \quad \frac{d}{dx} e^{i\lambda x} = i\lambda \, e^{i\lambda x}$$
(2.5)

(the first formula follows by (2.3)). Dunkl [8] generalized the operator d/dx to a mixture of a differential and a reflection operator:

$$(D^{(k)}f)(x) := f'(x) + k \, \frac{f(x) - f(-x)}{x} \,. \tag{2.6}$$

Note that this Dunkl operator sends smooth functions to smooth functions. Let us define a generalized exponential function in terms of Bessel functions (2.1) by

$$\mathcal{E}_k(\lambda x) := \mathcal{J}_k(\lambda x) + \frac{i\lambda x}{2k+1} \mathcal{J}_{k+1}(\lambda x).$$
(2.7)

Then it follows immediately from well-known differential recurrence formulas for Bessel functions that

$$D^{(k)} \mathcal{E}_k(\lambda x) = i\lambda \mathcal{E}_k(\lambda x).$$
(2.8)

Formulas (2.7) and (2.8) generalize the formulas in (2.5). The function $x \mapsto$ $\mathcal{E}_k(\lambda x)$ is the unique C^{∞} function which equals 1 in 0 and which is eigenfunction with eigenvalue $i\lambda$ of $D^{(k)}$. For $(D^{(k)})^2$ we compute

$$(D^{(k)})^2 f(x) = f''(x) + \frac{2k}{x} f'(x) - k \frac{f(x) - f(-x)}{x^2}.$$

Thus, on even functions f the square of the Dunkl operator acts as the differential operator $(d/dx)^2 + 2kx^{-1} d/dx$. In particular, its action on $\mathcal{J}_k(\lambda x) = \frac{1}{2}(\mathcal{E}_k(\lambda x) + \mathcal{E}_k(-\lambda x))$

yields

$$(D^{(k)})^2 \mathcal{J}_k(\lambda x) = -\lambda^2 \mathcal{J}_k(\lambda x).$$

3. Preliminaries about root systems

3.1. Definition of root system. Let V be a d-dimensional real vector space with inner product $\langle ., . \rangle$. For $\alpha \in V \setminus \{0\}$ let s_{α} denote the orthogonal reflection with respect to the hyperplane orthogonal to α (cf. Fig. 1):

$$S_{\alpha}(\beta) := \beta - \frac{2\langle \beta, \alpha \rangle}{\langle \alpha, \alpha \rangle} \alpha \quad (\beta \in V)$$

A root system in V (see [15]) is a finite subset R of $V \setminus \{0\}$ which spans V and which satisfies for all $\alpha, \beta \in R$ the two properties that

$$s_{\alpha}(\beta) \in R \quad \text{and} \quad rac{2\langle eta, lpha
angle}{\langle lpha, lpha
angle} \in \mathbb{Z}.$$

Clearly, if $\alpha \in R$ then $-\alpha = s_{\alpha}(\alpha) \in R$. For convenience, we will restrict ourselves to the case of a reduced root system, i.e., a root system R such that, if $\alpha, \beta \in R$ and $\alpha = c\beta$ for some $c \in \mathbb{R}$, then $c = \pm 1$. The so-called irreducible root systems can be classified as four infinite families A_n, B_n, C_n, D_n of classical root systems and five exceptional root systems G_2, F_4, E_6, E_7, E_8 . Here the subscript denotes the rank of the root system, i.e. the dimension of V. There is one infinite family of non-reduced irreducible root systems: of type BC_n .





An example for d = 1 is the set $R := \{\pm 2\} \subset \mathbb{R}$ (root system of type A_1 , cf. Fig. 2). An example for d = 2 is the set $R = R_+ \cup (-R_+)$, where $R_+ := \{(1, -1), (2, 0), (1, 1), (0, 2)\} \subset \mathbb{R}^2$ (root system of type C_2 , cf. Fig. 3). In general, when we have a root system R in V then we can write it as a disjoint union $R = R_+ \cup (-R_+)$, where R_+ and $-R_+$ are separated from each other by a hyperplane in V through the origin. The choice for R_+ is not unique. The elements of R are called roots and the elements of R_+ are called positive roots.

Let GL(V) be the group of invertible linear transformations of V. The Weyl group W of the root system R is the subgroup of GL(V) which is generated by the reflections s_{α} ($\alpha \in R$). The group W is finite and it acts on R. It permutes the possible choices of R_+ in a simply transitive way.

3.2. Dunkl operators associated with R. Let R be a root system in V. Let $k: \alpha \mapsto k_{\alpha}: R \to [0, \infty)$ be a function which is W-invariant, i.e., which satisfies $k_{w\alpha} = k_{\alpha}$ for all $w \in W$ and all $\alpha \in R$. If R is an irreducible (reduced) root system then the Weyl group is transitive on all roots of equal length and there are at most two different root lengths. Thus k_{α} then assumes at most two different values. See the above examples: one root length in A_1 and two root lengths in C_2 . The function k is called a multiplicity function. The reason for this name is that root systems naturally arise in the structure theory of real semisimple Lie algebras, where roots have an interpretation as joint eigenvalues of certain operators and the k_{α} then are (integer) multiplicities of such eigenvalues.

For $\xi \in V$ we will denote by ∂_{ξ} the corresponding directional derivative. The Dunkl operators [8], [16] associated with the root system R and the multi-

plicity function k are defined as the operators $D_{\xi}^{(k)}: C^{\infty}(V) \to C^{\infty}(V)$ $(\xi \in V)$ given by

$$(D_{\xi}^{(k)}f)(x) := (\partial_{\xi}f)(x) + \sum_{\alpha \in R_{+}} k_{\alpha} \langle \alpha, \xi \rangle \, \frac{f(x) - f(s_{\alpha}x)}{\langle \alpha, x \rangle} \,. \tag{3.1}$$

This definition is easily seen to be independent of the choice of R_+ . In case of root system A_1 formula (3.1) reduces for $\xi := 1$ to formula (2.6). Note that the operator (3.1) consists of a term involving a first order derivative and terms involving reflection operators, just as we have seen in (2.6). It is an amazing fact, which can be proved in a rather straightforward way, that the operators $D_{\epsilon}^{(k)}$ commute:

$$[D_{\xi}^{(k)}, D_{\eta}^{(k)}] = 0 \quad (\xi, \eta \in V).$$

Let $\mathbb{D}^{(k)}$ be the algebra generated by the operators $D_{\xi}^{(k)}$. This is a commutative algebra. It can be shown that each W-invariant operator D in $\mathbb{D}^{(k)}$, when restricted in its action to the W-invariant C^{∞} functions on V, coincides with a

partial differential operator (so its reflection terms vanish when acting on a W-invariant function). The joint W-invariant eigenfunctions of the W-invariant operators in $\mathbb{D}^{(k)}$ are called *Bessel functions associated with* R. In the example A_1 things reduce to the one-variable considerations of §2.1 and §2.4. More generally, one may study the joint eigenfunctions of the full algebra $\mathbb{D}^{(k)}$ and one

may try to do harmonic analysis for these eigenfunctions. A lot of satisfactory results have been obtained, see [16] and the references given there.

3.3. Weight lattice associated with R. We still assume a root system R in V. The weight lattice P of R is defined [15] by

$$P := \{ \lambda \in V \mid \frac{2\langle \lambda, \alpha \rangle}{\langle \alpha, \alpha \rangle} \in \mathbb{Z} \text{ for all } \alpha \in R \}.$$

The subset P_+ of dominant weights is then given by

$$P_{+} := \{ \lambda \in P \mid \frac{2\langle \lambda, \alpha \rangle}{\langle \alpha, \alpha \rangle} \ge 0 \quad \text{for all } \alpha \in R_{+} \}.$$

It is easily seen that w(P) = P for $w \in W$, so the Weyl group acts on P. Moreover, it can be shown that each Weyl group orbit in P has a one-point intersection with P_+ :

 $\forall \lambda \in P \quad Card(W\lambda \cap P_+) = 1.$

Thus the dominant weights can be used as a set of representatives for the W-orbits in P.



Fig. 4. Root system C_2 with dominant weights and the set $\{\mu \in P_+ \mid \mu \prec \lambda\}$

We introduce a partial ordering on P which is induced by the root system: for $\lambda, \mu \in P$ we say that $\mu \prec \lambda$ if $\lambda - \mu = \sum_{\alpha \in R_+} m_{\alpha} \alpha$ for certain nonnegative integers m_{α} . For root system C_2 the concepts of this subsection are illustrated in Fig. 4.

3.4. Trigonometric polynomials associated with R. Let P be the weight lattice of a root system R in V. For $\lambda \in P$ define the function e^{λ} on V by

$$e^{\lambda}(x) := e^{i\langle \lambda, x \rangle} \quad (x \in V).$$

Note that $e^{\lambda} e^{\mu} = e^{\lambda + \mu}$. Thus the space

$$\mathcal{A} := \operatorname{Span}\{e^{\lambda} \mid \lambda \in P\}$$

is an algebra: the algebra of trigonometric functions on V (with respect to R). For a function f on V write $(wf)(x) := f(w^{-1}x)$ ($w \in W, x \in V$). Then $we^{\lambda} = e^{w\lambda} \ (w \in W, \ \lambda \in P).$ Put

$$m_{\lambda} := \sum_{\mu \in W\lambda} e^{\mu} \quad (\lambda \in P_+).$$

Then the functions m_{λ} are W-invariant and they form a basis of the space \mathcal{A}^{W} of W-invariant elements in \mathcal{A} .

Let the dual root lattice Q be defined by

 $Q := \{ \lambda \in V \mid \langle \lambda, \mu \rangle \in \mathbb{Z} \text{ for all } \mu \in P \}.$

This lattice gives rise to a torus

 $T := V/(2\pi Q^{\check{}}).$

Let $x \mapsto \dot{x}$ be the natural mapping of V onto T. Then each function f in A actually lives on T: $f(x) = f(\dot{x})$ for a suitable function f on T. In the example A_1 we have $P = \mathbb{Z}, P_+ = \{0, 1, 2, ...\}$, the algebra \mathcal{A} is spanned by the functions $x \mapsto e^{inx}$ $(n \in \mathbb{Z})$ and the subalgebra \mathcal{A}^W by the functions 1 and $x \mapsto 2\cos(nx)$ (n = 1, 2, ...). The torus T equals $\mathbb{R}/(2\pi\mathbb{Z})$.

4. Jacobi polynomials associated with R

4.1. Definition of Jacobi polynomials for R. Let R be a root system in V and let $k: R \to [0, \infty)$ be a W-invariant multiplicity function as before. Define a weight function δ_k on T by

$$\delta_k(x) := \prod_{\alpha \in R_+} |2\sin(\langle \alpha, x \rangle)|^{2k_\alpha}.$$
(4.1)

This definition is independent of the choice of R_+ . Define an inner product on the linear space \mathcal{A} by

$$\langle f,g \rangle_k := \int_T f(x) \overline{g(x)} \,\delta_k(x) \,d\dot{x} \quad (f,g \in \mathcal{A}).$$
 (4.2)

Here $d\dot{x}$ denotes Lebesgue measure on T, normalized such that the volume of T is equal to 1.

The Jacobi polynomial $P_{\lambda}^{(k)}$ (cf. [11]) of "degree" $\lambda \in P_{+}$ and of "order" k is an element of \mathcal{A}^W of the form

$$P_{\lambda}^{(k)} = \sum_{\substack{\mu \in P_+\\ \mu \prec \lambda}} c_{\lambda,\mu} m_{\mu}$$

such that $c_{\lambda,\lambda} = 1$ and

$$\langle P_{\lambda}^{(k)}, m_{\mu} \rangle_k = 0 \quad \text{if } \mu \in P_+ \text{ and } \mu \precneqq \lambda.$$
 (4.3)

Instead of (4.3) we can equivalently require that $P_{\lambda}^{(k)}$ satisfies the second order differential equation

$$\left(\Delta + \sum_{\alpha \in R_+} k_\alpha \,\cot(\frac{1}{2}\langle \alpha, x \rangle) \,\partial_\alpha\right) P_\lambda^{(k)}(x) = -\langle \lambda, \lambda + \sum_{\alpha \in R_+} k_\alpha \alpha \rangle \,P_\lambda^{(k)}(x). \tag{4.4}$$

In the example A_1 we obtain that $P_n^{(k)}(x) = \text{const.} C_n^k(\cos x)$, where C_n^k is the ultraspherical polynomial of §2.2. The case of the (non-reduced) root system BC_1 would have given us, more generally, the classical one-variable Jacobi polynomials.

4.2. Three problems and their solutions. As soon as the above definition of Jacobi polynomials associated with R is given, three highly nontrivial questions can naturally be posed:

1. It follows immediately from the definition that the orthogonality

$$\langle P_{\lambda}^{(k)}, P_{\mu}^{(k)} \rangle_k = 0 \tag{4.5}$$

holds if $\mu \not\supseteq \lambda$ or $\lambda \not\supseteq \mu$. What about (4.5) if λ and μ are not related in the partial ordering?

2. Prove the existence of a commutative algebra of differential operators with

d algebraically independent generators, such that the operators in this algebra have the $P_{\lambda}^{(k)}$ ($\lambda \in P_{+}$) as joint eigenfunctions. (Note that the operator in (4.4) can be taken as one of the generators.)

3. Give an explicit expression for $\langle P_{\lambda}^{(k)}, P_{\lambda}^{(k)} \rangle_k$, or rather for its two factors

$$\frac{\langle P_{\lambda}^{(k)}, P_{\lambda}^{(k)} \rangle_{k}}{\langle P_{0}^{(k)}, P_{0}^{(k)} \rangle_{k}} \quad \text{and} \quad \int_{T} \delta_{k}(x) \, d\dot{x}.$$

$$(4.6)$$

In the past few years all these questions have been answered in the positive sense. Let me give some indications.

- If problem 2 can be solved then the answer to 1 follows readily, cf. [11]. Indeed, we need sufficiently many differential operators having the $P_{\lambda}^{(k)}$ as eigenfunctions such that the joint eigenvalues, in their dependence on λ , separate the points of P_+ .
- For certain special choices of k the functions $P_{\lambda}^{(k)}$, renormalized such that $P_{\lambda}^{(k)}(0) = 1$, have an interpretation as spherical functions on compact sym-

metric spaces G/K, cf. [14]. (For instance, in case A_1 the ultraspherical polynomial $C_n^{\frac{1}{2}m-1}$ can be interpreted as spherical function on the (m-1)-dimensional sphere SO(m)/SO(m-1).) Then problems 1, 2 and the first half of problem 3 can be solved by using the group theoretic interpretation.

The orthogonality (4.5) for general λ, μ follows by Schur's orthogonality relations for matrix elements of irreducible unitary representations of G. The first expression in (4.6) was explicitly computed by Vretare [33] in terms of Harish-Chandra's *c*-function related to the spherical functions on the corresponding non-compact symmetric space. The algebra of differential operators in problem 2 can be obtained by taking the radial parts of the G-invariant differential operators on G/K.

- For the classical root systems question 2 could be answered in a positive way by giving explicit expressions for generators of the algebra, see [17] for BC_2 and A_2 , and Olshanetsky & Perelomov [26], Sekiguchi [31] and Debiard [6] for the higher rank cases.
- Heckman and Opdam [11] have given positive answers to 2, and hence to 1, by use of deep transcendental arguments. This also solved part of Problem 3 (the first expression in (4.6)). In 1982 Macdonald [21] had already given conjectures for the explicit evaluation of the second expression in (4.6), which could be proved in a number of special cases.
- Problem 3 for general λ was solved by Opdam [28] by using so-called shift operators [27]. The most simple example, for case A_1 , of such operators is the following pair of differential recurrence relations for Gegenbauer polynomials:

$$\frac{d}{dx}C_n^k(x) = \text{const.} C_{n-1}^{k+1}(x),$$

$$\left((1-x^2)^{-k+\frac{1}{2}} \frac{d}{dx} \circ (1-x^2)^{k+\frac{1}{2}}\right) C_{n-1}^{k+1}(x) = \text{const.} C_n^k(x).$$

By use of these two formulas we can write $\int_{-1}^{1} (C_n^k(x))^2 (1-x^2)^{k-\frac{1}{2}} dx$ as an explicit constant times $\int_{-1}^{1} (C_{n-1}^{k+1}(x))^2 (1-x^2)^{k+\frac{1}{2}} dx$. Opdam's shift operators in general have a similar structure of lowering λ and raising k, or conversely. The case of root system BC_2 was already considered in [17], [32].

4.3. Dunkl type operators. Some years after Heckman first solved the problems 1 and 2 of the previous subsection he discovered a dramatical simplification [12] for proving these results. For a given root system R in V and a given multiplicity function k he wrote down a trigonometric variant of the Dunkl operators (3.1) for $\xi \in V$:

$$(D_{\xi}^{(k)}f)(x) := (\partial_{\xi}f)(x) + \frac{1}{2} \sum_{\alpha \in R_{+}} k_{\alpha} \langle \alpha, \xi \rangle \cot(\frac{1}{2} \langle \alpha, x \rangle) (f(x) - f(s_{\alpha}x))$$

 $(x \in V, f \in C^{\infty}(V)).$ (4.7)

Now the operators $D_{\mathcal{E}}^{(k)}$ will no longer commute, in general. However, Heckman showed that the operators $\sum_{\eta \in W\xi} (D_{\eta}^{(k)})^j$ ($\xi \in V, j = 0, 1, 2, ...$), when restricted to the W-invariant C^{∞} functions on V, coincide with differental operators which commute with each other and form a commutative algebra. This is the algebra looked for in Problem 2 of the previous subsection. The Jacobi

polynomials $P_{\lambda}^{(k)}$ are the joint eigenfunctions of the operators in this algebra. By this approach, Heckman also obtained a quick existence proof for Opdam's shift operators.

Next Cherednik [3] made a slight but significant variation in Heckman's Dunkl type operators (4.7). He put

$$(\widetilde{D}_{\xi}^{(k)}f)(x) := (\partial_{\xi}f)(x) + \sum_{\alpha \in R_{+}} k_{\alpha} \langle \alpha, \xi \rangle \frac{1}{1 - e^{-\alpha}(x)} (f(x) - f(s_{\alpha}x))$$
$$-\frac{1}{2} \sum_{\alpha \in R_{+}} k_{\alpha} \langle \alpha, x \rangle f(x).$$

(Here I took the part of the right hand side on the second line from Opdam [13, p.86]; Cherednik is not very specific about this part of his formula.) Cherednik's operators have the nice property that they mutually commute, without the need of first restricting to W-invariant functions. On the other hand, they do not share the property $w D_{\xi}^{(k)} w^{-1} = D_{w\xi}^{(k)}$ of Heckman's operators. Anyhow, by means of Cherednik's operators one can draw the same conclusions as by Heckman's operators, and in a similar way. Moreover, a structure of graded Hecke algebra can be associated with Cherednik's operators.

5. Macdonald polynomials associated with R

5.1. Definition of Macdonald polynomials. Let 0 < q < 1. We keep the assumptions of §4.1 except that we replace the weight function δ_k in (4.1) by

$$\delta_{k,q}(x) := \prod_{\alpha \in R_+} \left| \frac{(e^{i\langle \alpha, x \rangle}; q)_{\infty}}{(q^{k_{\alpha}} e^{i\langle \alpha, x \rangle}; q)_{\infty}} \right|^2.$$

Then the Macdonald polynomials $P_{\lambda}^{(k,q)}$ were defined by Macdonald [23], [24] just as the Jacobi polynomials $P_{\lambda}^{(k)}$, but with the inner product in (4.2) replaced by

$$\langle f,g\rangle_{k,q} := \int_T f(x)\overline{g(x)}\,\delta_{k,q}(x)\,dx \quad (f,g\in\mathcal{A}).$$

In the case of root system A_1 the Macdonald polynomials coincide with the q-ultraspherical polynomials $x \mapsto C_n^{k,q}(\cos x)$. For any root system R, in the limit for $q \uparrow 1$, the Macdonald polynomial $P_{\lambda}^{(k,q)}$ tends to the corresponding Jacobi polynomial $P_{\lambda}^{(k)}$.

Macdonald gives some explicit q-difference operators of which the $P_{\lambda}^{(k,q)}$ are eigenfunctions. Although these operators, except for root system A_n (where they were independently found by Ruijsenaars [30]) do not yet give a full commutative algebra of operators having the $P_{\lambda}^{(k,q)}$ as joint eigenfunctions, the additional parameter q gives enough freedom such that already the eigenvalue of one such operator separates the elements of P_+ for generic q, by which a

positive answer to question 1 in $\S4.2$ can be given for the case of Macdonald polynomials. Taking limits for $q \uparrow 1$ then yields the same positive answer for the case of Jacobi polynomials. This is an alternative to Heckman's approach via Problem 2. Macdonald also gives conjectured explicit expressions for the squared norms $\langle P_{\lambda}^{(k,q)}, P_{\lambda}^{(k,q)} \rangle_{k,q}$.

5.2. Askey-Wilson polynomials for root system BC_n . The author [18] introduced for the non-reduced root system BC_n a class of polynomials having two more parameters than Macdonald's class for BC_n . This extended class reduces for n = 1 to the Askey-Wilson polynomials [2]. In [18] only one explicit q-difference operator was given having the BC_n -polynomials as eigenfunctions, but this was sufficient for establishing orthogonality. Later, van Diejen [7] gave explicit expressions for the generators af a full commutative algebra of operators having the BC_n polynomials as joint eigenfunctions.

5.3. Cherednik's approach to Macdonald polynomials. Cherednik [4], [5] succeeded to give positive answers to questions 2 and 3 in §4.2. In the context of certain representations of affine Hecke algebras he could realize a commutative algebra of operators which have the Macdonald polynomials as joint eigenfunctions. In the same context he could realize q-analogues of Opdam's shift operators and next, by the same technique as in Opdam, prove Macdonald's conjectures in the q-case.

It is beyond the scope of this short survey to explain Cherednik's approach in any detail. In May 1994 I. G. Macdonald delivered some very helpful lectures

in Leiden in order to explain Cherednik's approach. Let me here only give a few indications. Just as a Hecke algebra is a deformation of the group algebra of a Weyl group, an affine Hecke algebra (cf. [20]) deforms the group algebra of an affine Weyl group. If R is an irreducible root system in V with Weyl group W then the (extended) affine Weyl group is the semidirect product $W := W \ltimes P$, where the dual weight lattice P is defined as $P := \{\lambda \in V \mid \langle \lambda, \alpha \rangle \in \mathbb{Z}\}$, an abelian group under addition. Then W acts as a group of motions on V, with P acting as a group of translations. The group W also acts on A, with W acting as before and with the action of P still depending on a parameter q.

The affine Hecke algebra H can be defined in terms of generators and relations which still depend on the values of a W-invariant function $\alpha \mapsto t_{\alpha}$ on R. Corresponding to a choice of R_+ we can define P_+ . The embedding of P_+ in H then generates a commutative subalgebra \mathcal{Y} of H.

For given q we can use the action of W on A in order to define an action of H on A, by specifying the action for a set of generators of H (Demazure operators). This action depends on q and the t_{α} . Put $t_{\alpha} = q^{-k_{\alpha}/2}$, where $k: R \rightarrow [0, \infty)$ is a multiplicity function. Then Cherednik proves that the

Macdonald polynomials $P_{\lambda}^{(k,q)}$ are the joint eigenfunctions of the W-invariant elements in the commutative algebra \mathcal{Y} . This answers question 2 in §4.2.

At the moment it is still an open problem to extend Cherednik's approach to the BC_n polynomials. As van Diejen [7] already answered question 2 in §4.2 in a positive way for this case, it would be nice to complement van Diejen's

constructive approach with the deep conceptual approach via affine Hecke algebras.

References

 R. Askey & M. E. H. Ismail, A generalization of ultraspherical polynomials, in Studies in Pure Mathematics, P. Erdös (ed.), Birkhäuser, 1983, pp. 55-78.

- [2] R. Askey & J. Wilson, Some basic hypergeometric orthogonal polynomials that generalize Jacobi polynomials, Memoirs Amer. Math. Soc. (1985), no. 319.
- [3] I. Cherednik, A unification of Knizhnik-Zamolodchikov and Dunkl operators via affine Hecke algebras, Invent. Math. 106 (1991), 411–431.
- [4] I. Cherednik, Double affine Hecke algebras, Knizhnik-Zamolodchikov equations, and Macdonald's operators, Duke Math. J., Internat. Math. Res. Notices (1992), no. 9, 171–180.
- [5] I. Cherednik, The Macdonald constant-term conjecture, Duke Math. J., Internat. Math. Res. Notices (1993), no. 6, 165–177.
- [6] A. Debiard, Système différentiel hypergéométrique et parties radiales des opérateurs invariants des espaces symétriques de type BC_p, in Séminaire d'Algèbre Paul Dubreil et Marie-Paule Malliavin, M.-P. Malliavin (ed.), Lecture Notes in Math. 1296, Springer, 1988, pp. 42–124.
- [7] J. F. van Diejen, Commuting difference operators with polynomial eigenfunctions, Compositio Math., to appear.
- [8] C. F. Dunkl, Differential-difference operators associated to reflection groups, Trans. Amer. Math. Soc. 311 (1989), 167–183.
- [9] A. Erdélyi, W. Magnus, F. Oberhettinger & F. G. Tricomi, Higher transcendental functions, Vol. II, McGraw-Hill, 1953.
- [10] P. I. Etingof & A. A. Kirillov. Jr., Macdonald's polynomials and representations of quantum groups, Math. Res. Lett. 1 (1994), no. 3, 279–296.
- [11] G. J. Heckman (part I with E. M. Opdam), Root systems and hypergeometric functions I, II, Compositio Math. 64 (1987), 329-352, 353-373.
- [12] G. J. Heckman, An elementary approach to the hypergeometric shift operators of Opdam, Invent. Math. 103 (1991), 341-350.
- [13] G. J. Heckman (part I, II) and E. M. Opdam (part III), Hypergeometric functions and representation theory, in "Seminar Hypergeometric Functions", Communications of the Mathematical Institute 18, Utrecht University, 1994, pp. 72–93.
- [14] S. Helgason, Groups and geometric analysis, Academic Press, 1984.
- [15] J. E. Humphreys, Introduction to Lie algebras and representation theory, Springer, 1972.
- [16] M. F. E. de Jeu, Dunkl operators, Dissertation, University of Leiden, 1994.

- [17] T. H. Koornwinder, Orthogonal polynomials in two variables which are eigenfunctions of two algebraically independent partial differential operators I, II, III, IV, Nederl. Akad. Wetensch. Proc. Ser. A 77 (1974), 48–58, 59–66, 357–369, 370–381.
- [18] T. H. Koornwinder, Askey-Wilson polynomials for root systems of type BC, in Hypergeometric functions on domains of positivity, Jack polynomials, and applications, D. St. P. Richards (ed.), Contemp. Math. 138, Amer. Math. Soc., 1992, pp. 189–204.
- [19] T. H. Koornwinder, Askey-Wilson polynomials as zonal spherical functions on the SU(2) quantum group, SIAM J. Math. Anal. 24 (1993), 795–813.
- [20] G. Lusztig, Affine Hecke algebras and their graded version, J. Amer. Math. Soc. 2 (1989), 599–635.
- [21] I. G. Macdonald, Some conjectures for root systems, SIAM J. Math. Anal.
 13 (1982), 988–1007.
- [22] I. G. Macdonald, Symmetric functions and Hall polynomials, Oxford University Press, 1979; second ed. 1994.
- [23] I. G. Macdonald, A new class of symmetric functions, in Séminaire Lotharingien de Combinatoire, 20e Session, L. Cerlienco & D. Foata (eds.), Publication de l'Institut de Recherche Mathématique Avancée 372/S-20, Strasbourg, 1988, pp. 131–171.
- [24] I. G. Macdonald, Orthogonal polynomials associated with root systems, preprint, 1988.
- [25] M. Noumi, Macdonald's symmetric polynomials as zonal spherical functions on some quantum homogeneous spaces, Adv. in Math., to appear.
- [26] M. A. Olshanetsky & A. M. Perelomov, Quantum integrable systems related to Lie algebras, Phys. Rep. 94 (1983), 313-404.
- [27] E. M. Opdam, Root systems and hypergeometric functions III, IV, Compositio Math. 67 (1988), 21-49, 191-209.
- [28] E. M. Opdam, Some applications of hypergeometric shift operators, Invent. Math. 98 (1989), 1–18.
- [29] E. M. Opdam, Harmonic analysis for certain representations of graded Hecke algebras, Report W 93-18, Dept. of Math., Univ. of Leiden, 1993.
- [30] S. N. M. Ruijsenaars, Complete integrability of relativistic Calogero-Moser systems and elliptic function identities, Comm. Math. Phys. 110 (1987), 191–213.
- [31] J. Sekiguchi, Zonal spherical functions on some symmetric spaces, Publ. Res. Inst. Math. Sci. 12 Suppl. (1977), 455-459.
- [32] I. G. Sprinkhuizen, Orthogonal polynomials in two variables. A further analysis of the polynomials orthogonal on a region bounded by two lines and a parabola, SIAM J. Math. Anal. 7 (1976), 501-518.
- [33] L. Vretare Elementary spherical functions on symmetric spaces, Math. Scand. 39 (1976), 343-358.